

UNCLASSIFIED

AD **246 226**

*Reproduced
by the*

ARMED SERVICES TECHNICAL INFORMATION AGENCY
ARLINGTON HALL STATION
ARLINGTON 12, VIRGINIA



UNCLASSIFIED

NOTICE: When government or other drawings, specifications or other data are used for any purpose other than in connection with a definitely related government procurement operation, the U. S. Government thereby incurs no responsibility, nor any obligation whatsoever; and the fact that the Government may have formulated, furnished, or in any way supplied the said drawings, specifications, or other data is not to be regarded by implication or otherwise as in any manner licensing the holder or any other person or corporation, or conveying any rights or permission to manufacture, use or sell any patented invention that may in any way be related thereto.

AEDC-TN-60-181

CORNELL AERONAUTICAL LABORATORY, INC.

CATALOGED BY ASTIA
AS AD NO 246 226

REPORT NO. AD-1345-W-3

APPROXIMATE SOLUTIONS FOR NONEQUILIBRIUM AIRFLOW IN HYPERSONIC NOZZLES

BY

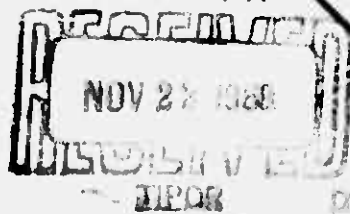
D. W. BOYER

A. Q. ESCHENROEDER

A. L. RUSSO

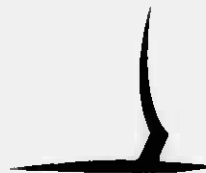
AUGUST 1960

ASTIA



XEROX

B U F F A L O , N E W Y O R K



CORNELL AERONAUTICAL LABORATORY, INC.
Buffalo 21, New York

REPORT NO. AD-1345-W-3

APPROXIMATE SOLUTIONS FOR NONEQUILIBRIUM
AIRFLOW IN HYPERSONIC NOZZLES

BY

D. W. Boyer, A. Q. Eschenroeder, and A. L. Russo

AUGUST 1960

Program Area 75D A

ARDC Project - 4773

Contract No. AF40(6DD)-804

AEDC TN-6D-181

Arnold Engineering Development Center
Arnold Air Force Station, Tennessee

AIR RESEARCH AND DEVELOPMENT COMMAND
UNITED STATES AIR FORCE

FOREWORD

This technical note was prepared as part of the Wave Superheater Hypersonic Tunnel project being constructed at Cornell Aeronautical Laboratory, Inc., Buffalo, New York.

This research is being supported by the Advance Research Project Agency and monitored by the Arnold Engineering Development Center of the Air Research and Development Command under Contract AF40(600)-804.

The authors are pleased to acknowledge the helpful discussions held with Dr. J. G. Hall of CAL throughout the course of this work.

ABSTRACT

A method has been described to estimate the effects of finite chemical reaction rates on the one-dimensional expansion of air in hypersonic nozzles. The calculations have considered stagnation temperatures from 4000 to 6000°K, stagnation pressures from 100 to 1000 atmospheres, and a range of wedge and axisymmetric nozzle shapes. Vibrational equilibrium is assumed.

A relaxation length criterion has been applied to a simplified kinetic model of air to determine the approximate location of freezing for flow with finite reaction rates in each nozzle configuration. Equilibrium composition profiles obtained by machine calculation were used for the calculation of the relaxation lengths. The resultant chemically frozen expansions have been calculated and are presented in tabular and graphical form. In all cases freezing occurs fairly early in the nozzle. Further, freezing in the nozzle is delayed by an increase in stagnation pressure, an increase in stagnation temperature, and by the use of gradually expanding nozzles. The effect of freezing is to reduce the pressure, velocity and temperature at a particular area ratio from the corresponding equilibrium values, and to increase the density and Mach number. The change in temperature and Mach number may be considerable whereas the density and flow velocity are relatively unaffected by freezing.

TABLE OF CONTENTS

	<u>Page</u>
FOREWORD	i
ABSTRACT	iii
NOTATION	vii
I. INTRODUCTION	1
II. SOLUTION OF THE EQUILIBRIUM FLOW	3
A. Governing Equations	3
B. Equilibrium Analysis	5
III. APPROXIMATE SOLUTIONS FOR FINITE-RATE AIR EXPANSIONS	11
A. Relaxation Length Criterion for Freezing	11
B. Nozzle Geometry Considered	17
C. Results of Approximate Finite-Rate Solutions	18
IV. SOLUTIONS FOR EXPANSION AFTER FREEZING	20
V. FURTHER KINETIC CONSIDERATIONS	23
VI. CONCLUSIONS	25
APPENDIX A - NEWTON-RAPHSON METHOD FOR SOLVING THE EQUILIBRIUM PROBLEM	27
REFERENCES	30
TABLES 1 - 26	32
FIGURES 1 - 31	68

NOTATION

Dimensional

- A' cross sectional area of nozzle
- D dissociation energy per mole
- ΔF_ℓ° change in standard free energy for ℓ th reaction
- h_i' enthalpy of i th species including energy of formation, cal/mole
- h Planck's constant
- H' enthalpy of mixture, cal/gm
- k_f dissociation rate coefficient, $\text{cm}^3/\text{mole sec}$
- k_r recombination rate coefficient, $\text{cm}^6/\text{mole}^2 \text{ sec}$
- K_e equilibrium constant = k_f / k_r , mole/cm^3
- $K_{p\ell}, K_{x\ell}$ equilibrium constant referred to partial pressures and mole fractions respectively
- K Boltzmann constant
- ℓ characteristic nozzle dimension L/a
- L characteristic throat dimension, cm
- m' mass flow $\rho' u' A'$
- p' fluid pressure
- r' local relaxation length, Equation (35)
- r'_∞ local relaxation length based on infinite-rate equilibrium, Equation (36)
- R_o universal gas constant
- s'_o entropy (per unit mass) at reservoir conditions
- $S_i^{\circ'}$ molar entropy of i th species at standard pressure
- T' local absolute temperature of gas

u'	local gas velocity
x'	distance from nozzle throat
γ_i	concentration of species i, moles/gm
θ'_0	characteristic dissociation temperature D/R_0
θ'_v	characteristic vibrational temperature $h\nu/K$
μ	molecular weight of undissociated mixture
$\bar{\mu}$	molecular weight of mixture
ν	characteristic frequency of molecule
ρ'	fluid density
ψ	parameter $\frac{2 \ell \rho_0'^2}{\sqrt{\frac{R_0 T_0'}{\mu}}} k_{R_0}^{(2)}$

Dimensionless

a	tangent of nozzle semi-angle
A_{jk}	member of correction coefficient matrix in Newton-Raphson method
A	A'/A_*
C	number of linearly independent species
F_j	function used in Newton-Raphson method defined by Equations (A-2) and (A-3)
h_i	$h'_i / R_0 T_0'$
h_k	fractional correction to the k th independent variable in Newton-Raphson iteration
H	$H' / \frac{R_0 T_0'}{\mu}$
m	mass flow $\rho u A$

m	number of chemical elements present
M	Mach number
M_i	chemical formula of i th species
n	exponent in area relation $A = 1 + x^n$
p	p'/p_0'
q_j	mole fraction of j th component in a hypothetical system containing components only
Q_k	number of gram atoms of k th chemical element present
r	r'/ℓ
r_∞	r'_∞/ℓ
s	number of species
s_0	s_0'/R_0
S_i^0	$S_i^{00'}/R_0$
T	T'/T_0'
u	$u'/\sqrt{\frac{R_0 T_0'}{\mu}}$
x	x'/ℓ
X_i	mole fraction of i th species
y_i	formula vector of i th species
α_{jk}	atoms of k th chemical element per molecule of j th chemical species
β_i	$k_f^{(i)}/k_f^{(2)}$, Equation (32)
δ_{jk}	Kronecker delta
ϵ_h	maximum permissible fractional correction to mole fractions
θ_0	θ'_0/T_0'
θ_v	θ'_v/T_0'

ν_{ij} stoichiometric coefficients
 ρ ρ'/ρ_0'

Subscripts

α, i, j, k, l sum, product or matrix indices
o denotes condition at reservoir
e denotes local equilibrium value
 ∞ denotes infinite-rate equilibrium value
* denotes value at throat, $A = 1$
f denotes value at freezing

I INTRODUCTION

The stagnation enthalpies encountered in the flow of gases through nozzles at low supersonic velocities are usually sufficiently low so that chemical kinetics are unimportant. Such flows may be accurately described by simple equations which consider the participation of only the translational and rotational degrees of freedom of the molecules (constant γ). However, in the high temperature flows occurring in hypersonic nozzles, such simplified considerations are no longer applicable because of the interchange of energy between the internal and translational degrees of freedom and the coupling of the chemistry with the gas dynamics through dissociation. Figure 1 shows that at stagnation conditions pertinent to the present study, a significant portion of the total energy resides in the internal energy modes.

The gas dynamic and chemical behavior of a high temperature gas undergoing expansion in complete equilibrium may be calculated with reasonable precision (Sec. II). Complete chemical equilibrium, however, implying infinite reaction rates, is a condition which may only be approached, to a greater or lesser degree, in the expansion of high temperature real gases. Studies of the effects of chemical nonequilibrium, in particular nozzle flows^{1, 2, 3, 4} involving only a single dissociating molecular species, have shown that the flow properties may depart considerably from their local equilibrium values. The effect of finite chemical reaction rates, as well as nozzle geometry, may result in actual freezing of the gas composition at some point in the nozzle.

The purpose of this report is to present quantitative estimates for the effect of flow freezing on the departure of test section properties from their equilibrium values in the expansion of air through hypersonic wind tunnel nozzles. The analysis has considered stagnation temperatures from 4000°K to 6000°K and stagnation pressures from 100 to 1000 atmospheres, representing a range of stagnation enthalpies of either current or future interest. The expansions are considered for two nozzle configurations; a two-dimensional wedge-type nozzle with a sharp throat and a hyperbolic axisymmetric nozzle

which is conical for large area ratios. The solutions cover a comprehensive range of geometries for both types of nozzle.

The outline of the report is as follows. The equilibrium chemical composition of the gas at the reservoir condition is first computed for all stagnation states by the method outlined in Sec. II. The equilibrium (infinite-rate) expansions are then calculated, the solutions providing both the gas dynamic behavior and gas composition as a function of temperature. Because of their usefulness, the results of the equilibrium calculations are included in this report in tabular and graphical form.

The concept of a local chemical relaxation length⁵ is then introduced (Sec. III) and used to develop an approximate criterion for freezing. This concept was first used to develop a freezing criterion in Ref. 3, for both a pure diatomic gas or diatomic gas plus any diluent. Some air calculations were also performed in Ref. 3. The present kinetic model for air assumes finite reaction rates for oxygen, with nitrogen and nitric oxide assumed to remain in chemical (infinite-rate) equilibrium. A relaxation length is calculated, based upon the equilibrium flow properties, and used to determine the approximate nozzle location at which appreciable freezing has occurred, i.e., to predict, approximately, the final chemical state of the expanding flow for finite reaction rates.

Finally (Sec. IV), calculation of the nozzle expansions are completed with all chemical species fixed at their "freezing" mass fractions. The equilibrium solutions serve as a basis of comparison for the frozen (zero-rate) solutions. Both the equilibrium and frozen expansion processes are isentropic.¹⁴ The comparisons have been referred to as providing quantitative "estimates" since although the calculations of the frozen expansions are in themselves exact, the criterion for the selection of the equilibrium state to serve as input to the frozen calculations is approximate.

II SOLUTION OF THE EQUILIBRIUM FLOW

A. Governing Equations

It is assumed that the gas is a mixture of ideal gases, that the flow is quasi-one-dimensional, and that diffusion, heat conduction and viscous processes are of negligible significance. Steady, adiabatic, critical flow is assumed to exist in the nozzle.

The equations which determine the flow are,

Conservation of Momentum:

$$u' du' + \frac{dp'}{\rho'} = 0 \quad (1)$$

where u' = velocity
 p' = pressure
 ρ' = density

Global Conservation of Energy:

$$\frac{u'^2}{2} + H' = H_0' \quad (2)$$

where H' = mixture enthalpy (cal/gm)

Global Conservation of Mass:

$$m' = \rho' u' A' \quad (3)$$

where m' = mass flow
 A' = cross-sectional area

The state equations are,

Thermodynamic State:

$$p' = \rho' \frac{R_0}{\bar{\mu}} T' \quad (4)$$

where $\bar{\mu}$ = molecular weight of mixture and is given by

$$\bar{\mu} = \frac{1}{\sum_{i=1}^s x_i} \quad (5)$$

where x_i = mass concentration of species i in moles per unit mass.

Caloric State:

$$H' = \sum_{i=1}^s x_i h_i' \quad (6)$$

where h_i' = molar enthalpy of i th species including thermal and formation contributions. h_i' is only a function of temperature.

In dimensionless form, the equations are written

$$u du + \frac{dp}{\rho} = 0 \quad (7)$$

$$\frac{u^2}{2} + H = H_0 \quad (8)$$

$$m = \rho u A \quad (9)$$

$$p = \rho T \frac{\sum_{i=1}^s x_i}{\sum_{i=1}^s x_{i_0}} \quad (10)$$

$$H = \mu \sum_{i=1}^s x_i h_i \quad (11)$$

where

$$u = u' / \sqrt{\frac{R_o T_o'}{\mu}}$$

$$H = H' / \frac{R_o T_o'}{\mu}$$

$$p = p' / p_o'$$

$$\rho = \rho' / \rho_o'$$

$$T = T' / T_o'$$

$$A = A' / A'_*$$

μ = molecular weight of undissociated (cold) gas

R_o = Universal gas constant

A'_* = throat area

Except for the molecular weight, μ , the reference conditions for the non-dimensional forms are those at the equilibrium reservoir state.

B. Equilibrium Analysis

The methods and notation of Brinkley⁶ were used for the equilibrium analysis wherever applicable. At the end of this section, the Brinkley technique is extended to include the computation of states along an isentropic path.

The chemical formula of the i th species may be expressed in vectorial form as

$$y_i = (\alpha_{i1}, \alpha_{i2}, \dots, \alpha_{ik}, \dots, \alpha_{im}) \quad (i=1, 2, \dots, s) \quad (12)$$

where α_{ik} is the number of atoms of the k th chemical element per molecule of the i th species and m is the number of elements present. If the rank of the matrix of α_{ik} is c , there exist c linearly independent y_i , and $(s-c)$ dependent y_i which are given by linear combinations of the c independent ones.

$$\sum_{j=1}^c \nu_{lj} y_j = y_l \quad (l = c+1, c+2, \dots, s) \quad (13)$$

As expressed by Eq. (13), ν_{lj} are coefficients in the linear dependencies of vectors y_l upon y_j .

In the equilibrium analysis the independent species are designated ($j = 1, 2, \dots, c$) and are called "components." The number of components c usually equals the number of chemical elements present m . Equations (13) correspond to $(s-c)$ chemical reactions postulated to form the dependent species if the formula vectors are replaced with chemical symbols.

It is useful to consider a hypothetical system. This system has the same atomic constitution as the actual system, but is chemically combined such that only components are present. If Q_k is the number of gram-atoms of the k th element present, and q'_j is the number of moles of the j th component contained in the hypothetical system, then

$$\sum_{j=1}^c \alpha_{jk} q'_j = Q_k \quad (k=1, 2, \dots, m) \quad (14)$$

The q'_j are normalized by

$$q_j = \frac{q'_j}{\sum_{j=1}^c q'_j} \quad (15)$$

thereby scaling the total size of the system if q_j are considered to be numbers of moles.

Global mass conservation may be expressed by

$$\sum_{j=1}^c n_j \mu_j + \sum_{l=c+1}^s n_l \mu_l = \sum_{j=1}^c q_j \mu_j \quad (16)$$

where μ_j is the molecular weight and n_j is the number of moles of j th species. Mass is conserved in each formation reaction of Eqs. (13) so that

$$\mu_l = \sum_{j=1}^c \nu_{lj} \mu_j \quad (l = c+1, c+2, \dots, s) \quad (17)$$

Substituting μ_l from Eq. (17) and equating coefficients of μ_j in Eq. (16) gives

$$n_j + \sum_{l=c+1}^s \nu_{lj} n_l = q_j \quad (j = 1, 2, \dots, c) \quad (18)$$

Division of Eq. (18) by the total number of moles n gives

$$X_j + \sum_{l=c+1}^s \nu_{lj} X_l = \frac{q_j}{n} \quad (j = 1, 2, \dots, c) \quad (19)$$

where X are mole fractions. By definition

$$\sum_{j=1}^c X_j + \sum_{l=c+1}^s X_l = 1 \quad (20)$$

Equation (19) is summed over j to give

$$\sum_{j=1}^c X_j + \sum_{l=c+1}^s \nu_l X_l = \frac{1}{n} \quad (21)$$

where

$$\nu_l = \sum_{j=1}^c \nu_{lj}$$

Substituting for X_j from Eq. (20) into Eq. (21) results in

$$1 + \sum_{\ell=c+1}^S (\nu_{\ell} - 1) X_{\ell} = \frac{1}{n} \quad (22)$$

Therefore, Eq. (19) may be written

$$X_j = q_j - \sum_{\ell=c+1}^S [\nu_{\ell j} - q_j (\nu_{\ell} - 1)] X_{\ell} \quad (j = 1, 2, \dots, c) \quad (23)$$

giving c of the S equations necessary for determining the equilibrium composition at a specified temperature and pressure. The remaining equations are

$$X_{\ell} = K_{x\ell} \prod_{j=1}^c X_j^{\nu_{\ell j}} \quad (\ell = c+1, c+2, \dots, S) \quad (24)$$

where $K_{x\ell}$ are equilibrium constants based on mole fractions. These constants are related to $K_{p\ell}$ by

$$K_{x\ell} = K_{p\ell} p^{-(\nu_{\ell}-1)} \quad (\ell = c+1, c+2, \dots, S) \quad (25)$$

$K_{p\ell}$ are calculated from chemical potentials by

$$K_{p\ell} = \exp \left(- \frac{\Delta F_{\ell}^{\circ}}{R_0 T'} \right) \quad (26)$$

where ΔF_{ℓ}° is the change in free energy at standard pressure for the ℓ th reaction. The ΔF_{ℓ}° were calculated by means of polynomial fits to various tabulated and derived thermodynamic functions provided by R. E. Duff of Los Alamos Scientific Laboratory. Equations (23) and (24) determine the equilibrium composition when temperature and pressure are specified; therefore, they can be employed for the reservoir state calculations.

For computation along an isentropic path, the Brinkley method must be modified. The composition and pressure will be computed at various temperatures for the entropy at the prescribed reservoir state. Since pressure

is an additional unknown quantity, one more equation is required. A statement that entropy is a constant is the desired equation.

$$s_0 = \frac{1}{\bar{\mu}} \sum_{i=1}^s x_i s_i^0 - \sum_{i=1}^s x_i \ln x_i - \ln p' \quad (27)$$

s_0 is the entropy (per unit mass) calculated at the reservoir state, and $\bar{\mu}$ is the mixture molecular weight which is calculated from

$$\bar{\mu} = \sum_{i=1}^s x_i \mu_i$$

s_i^0 is the molar entropy of the i th species at standard pressure. Simultaneous solution of the nonlinear algebraic equations (23), (24), and (27) gives the composition and pressure at various temperatures along an isentrope characterized by s_0 . The Newton-Raphson method as employed in the solution is outlined in Appendix A.

Additional quantities are computed from the results of the solution. Concentrations in moles per unit mass γ_i are found from

$$\gamma_i = \frac{x_i}{\bar{\mu}} \quad (i = 1, 2, \dots, s) \quad (28)$$

Density is found from Eq. (4), enthalpy from Eq. (6), velocity from Eq. (2), and mass flow per unit area m'/A' from Eq. (3). The state at which mass flow per unit area goes through a maximum defines the conditions at the throat of the equilibrium nozzle. The local to throat area ratios are then found by dividing the maximum m'/A' by the local m'/A' values. A Mach number is calculated at each step from the expression

$$M = u / \sqrt{\left(\frac{p}{\rho}\right)_{\text{isentropic equilibrium}}} \quad (29)$$

The thermodynamic state is independent of the form of $A(x)$ because the satisfaction of equilibrium conditions at any area is independent of the history of the flow.

The equilibrium air expansions were calculated on the CAL IBM-704 computer* for reservoir temperatures of 4000, 5000, 6000, 7000, and 8000°K and reservoir pressures of 100, 300, and 1000 atmospheres. The calculations assumed a 13 chemical specie system for air. The solutions for the gas dynamic properties and gas composition are presented for each equilibrium expansion in Tables 1-16. In addition, the air composition is shown as a function of area ratio for each case in Figs. 2-17.

*The Fortran programming for IBM machine computation was done by D. B. Larson, Computer Applications Branch, Systems Research Department, Cornell Aeronautical Laboratory. The IBM-704 program used in the solution of the equilibrium flow is available at CAL. In the event further equilibrium flow calculations are required, details will be provided upon request.

III APPROXIMATE SOLUTIONS FOR FINITE-RATE AIR EXPANSIONS

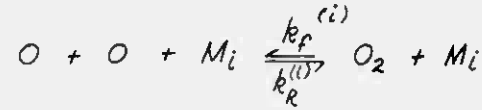
A. Relaxation Length Criterion for Freezing

The solution of nonequilibrium flows, which involve finite reaction rates, is rendered difficult by the need for simultaneous solution of inter-related chemical rate equations and the gas dynamic equations. Also, at the present time, the magnitudes of the required rate constants, and their temperature dependence, are still not known with certainty. It is worthwhile, therefore, to investigate a simple method for the approximate calculation of the final state of a gas undergoing expansion with finite reaction rates.

The present approach is that of Ref. 3 in which specific non-equilibrium airflows in hypersonic nozzles were calculated on the basis of a relaxation length. The results of the previous studies^{2,3} have shown that for flows involving a single relaxing species, the use of a local relaxation length, evaluated for the corresponding equilibrium flow, provides an excellent estimate of the final frozen state of the gas obtained from exact finite-rate solutions. Accordingly, a relaxation length is introduced in the air calculations on the basis of an assumption concerning the relative importance of the various chemical reactions.

The energy associated with the normally inert degrees of freedom for the air species, as a fraction of the total enthalpy, is shown in Fig. 1. It is seen that over the range of reservoir conditions pertinent to the present frozen expansion calculations, i.e., from 4000 to 6000°K, oxygen dissociation is the most important internal mode as it represents the source of greatest energy loss due to freezing. The percentage of the gas enthalpy associated with nitrogen dissociation is small compared with that of oxygen. Furthermore, since the energy associated with the formation of nitric oxide is small and essentially constant over a wide range of temperatures and pressures, the potential energy loss due to freezing of the nitric oxide composition is also quite small. Consequently, a relaxation length is derived assuming that only oxygen dissociation-recombination reactions occur at finite rates.

The reaction may be written



where M_1 = oxygen atom
 M_2 = oxygen molecule
 M_3 = nitrogen atom
 M_4 = nitrogen molecule
 M_5 = nitric oxide molecule

The equilibrium constant in terms of molar concentrations per unit mass (γ_i) has the assumed form

$$K_e(T') = \frac{k_f^{(i)}}{k_R^{(i)}} = \rho' \frac{\gamma_{O_e}^2}{\gamma_{O_2e}} \quad (30)$$

where $k^{(i)}$ are reaction rate coefficients. The subscript e denotes the local equilibrium value. The rate equation is then

$$\begin{aligned} \frac{d\gamma_0}{dt} = u' \frac{d\gamma_0}{dx'} &= 2\rho' k_f^{(2)} \gamma_{O_2}^2 + 2\rho' k_f^{(1)} \gamma_0 \gamma_{O_2} - 2\rho'^2 k_R^{(2)} \gamma_0^2 \gamma_{O_2} \\ &\quad - 2\rho'^2 k_R^{(1)} \gamma_0^3 + 2\rho' \gamma_{O_2} \sum_{i=3}^5 k_f^{(i)} \gamma_{M_i} - 2\rho'^2 \gamma_0^2 \sum_{i=3}^5 k_R^{(i)} \gamma_{M_i} \end{aligned} \quad (31)$$

On substituting Eq. (30) and rearranging, Eq. (31) becomes

$$u' \frac{d\gamma_0}{dx'} = 2\rho' K_e k_R^{(2)} \gamma_{O_2}^2 \left\{ 1 - \frac{\rho'}{K_e} \frac{\gamma_0^2}{\gamma_{O_2}} \right\} \left\{ 1 + \beta_1 \frac{\gamma_0}{\gamma_{O_2}} + \sum_{i=3}^5 \beta_i \frac{\gamma_{M_i}}{\gamma_{O_2}} \right\} \quad (32)$$

where

$$\beta_i = \frac{k_f^{(i)}}{k_f^{(2)}}$$

Now the conservation equation for oxygen atoms is

$$\gamma_o + \gamma_{NO} + 2\gamma_{O_2} = \text{CONST.}$$

$$\text{therefore } \gamma_{O_2} = B - \frac{\gamma_o}{2} \quad \text{where } B = \text{CONST} - \frac{\gamma_{NO}}{2}$$

Defining

$$k_R^{(2)} = k_{R_o}^{(2)} T^{-2}$$

$$x = x'/\ell$$

ℓ = length characterizing nozzle geometry

$$= \frac{L}{a} = \frac{\text{half throat height (cm)}}{\text{tangent of nozzle semi-angle}}$$

$$\psi = \frac{2\ell s_o'^2}{\sqrt{R_o T_o'}} k_{R_o}^{(2)}$$

and substituting in Eq. (32), noting that K_e (Eq. 30) is defined for equilibrium values of the γ_i , there results

$$\frac{d\gamma_o}{dx} = -\frac{\psi s_o'^2}{u T^2} \left\{ 1 + \beta_1 \frac{\gamma_o}{\gamma_{O_2}} + \sum_{i=3}^5 \beta_i \frac{\gamma_{M_i}}{\gamma_{O_2}} \right\} \left\{ B - \frac{\gamma_o}{2} \right\} \quad (33)$$

$$\left\{ \gamma_o + \frac{B\gamma_{oe}}{B - \frac{\gamma_{oe}}{2}} \right\} \left\{ \gamma_o - \gamma_{oe} \right\}$$

where γ_{oe} = local finite-rate equilibrium value of γ_o , i.e., the equilibrium value of γ_o for the actual temperature and density distribution.

Equation (33) is then of the form introduced by Heims⁵

i. e. ,

$$\frac{d\gamma_0}{dx} = -\frac{\gamma_0 - \gamma_{0e}}{r} \quad (34)$$

where

$$r = \frac{r'}{l} = \frac{u T^2}{\psi \rho^2 \left[1 + \beta_1 \frac{\gamma_0}{\gamma_{02}} + \sum_{i=3}^5 \beta_i \frac{\gamma_{Mi}}{\gamma_{02}} \right] \left[\beta - \frac{\gamma_0}{2} \right] \left[\gamma_0 + \frac{B \gamma_{0e}}{\beta - \frac{\gamma_{0e}}{2}} \right]} \quad (35)$$

The rate equation, Eq. (34), which incorporates a local relaxation length r , is used to derive an approximate freezing criterion for the finite-rate flow. The freezing criterion is given in a general form in Ref. 3 and its development is repeated here for completeness.

For flow near equilibrium, $\gamma_0 - \gamma_{0e} \ll \gamma_{0e}$ and it follows from Eq. (34) that $\left| \frac{d\gamma_0}{dx} \right| \ll \frac{\gamma_{0e}}{r}$. For near frozen flow, $\gamma_{0e} \ll \gamma_0$ so that $\left| \frac{d\gamma_0}{dx} \right| \approx \frac{\gamma_0}{r} \gg \frac{\gamma_{0e}}{r}$. Thus an indication of where significant freezing has occurred may be obtained by setting $\left| \frac{d\gamma_0}{dx} \right| = \frac{\gamma_{0e}}{r}$. This equality may be evaluated on the basis of the infinite rate equilibrium flow by using the average values of $\frac{d\gamma_0}{dx}$ and γ_0 through the freezing region. That is

$$\frac{d\gamma_0}{dx} = \frac{1}{2} \left(\frac{d\gamma_0}{dx} \right)_\infty$$

$$\gamma_{0e} = \frac{1}{2} (\gamma_0)_\infty$$

to give

$$\left(\frac{d\gamma_0}{dx} \right)_\infty = \frac{\gamma_{0\infty}}{r_\infty} \quad \text{or} \quad \left(\frac{d\gamma_0}{dA} \right)_\infty \left(\frac{dA}{dx} \right) = \frac{\gamma_{0\infty}}{r_\infty} \quad (36)$$

where the subscript ∞ refers to infinite-rate equilibrium values.

Some remarks should be made at this point on the selection of "equilibrium" values. As mentioned earlier in this section, the concept of

a local chemical relaxation length has been found useful for flows in which only a single mode of relaxation is involved, such as the dissociation relaxation of a diatomic gas, with or without inert diluent. For the more complex case of air, the relaxation length criterion has been applied to two simplified kinetic models of air in which the only finite-rate reactions considered are the oxygen dissociation and recombination processes. All species except O and O_2 were assumed to act only as second or third body colliders for the oxygen dissociation-recombination kinetics. The appropriate reaction rate constants were obtained from simple collision theory. Vibrational equilibrium was assumed for all molecular species.

Two basic models were considered in the calculations. In the first, the nitrogen atom and molecule and nitric oxide mass concentrations were frozen at their respective reservoir values during the expansion. In the second, these mass concentrations were taken as those for infinite-rate equilibrium (Sec. II). The flows with oxygen atom-molecule equilibration corresponding to these two models were used in the solution of Eq. (36).

Both models gave values for the frozen degree of oxygen dissociation (α_f) which agreed quite closely up to stagnation temperatures of 6000°K. For example, in a typical hyperbolic nozzle ($L/a = 1.0$) for a stagnation temperature of 5000°K and a stagnation pressure of 100 atm., $\alpha_f = .125$ with N_2 , N , and NO frozen at their reservoir concentrations and $\alpha_f = .149$ with N_2 , N , and NO in infinite-rate equilibrium. Since these are the limiting conditions for the N_2 , N , and NO reactions, this comparison demonstrates that oxygen dissociation-recombination reactions play the major role in air kinetics. The second model has been used in the calculations reported herein, i. e., the equilibrium solutions of Sec. II were used to evaluate relaxation lengths from Eq. (35).

In the above calculations the magnitude of $k_R^{(2)}$ at 3000°K has been taken as $1.16 \times 10^{15} \text{ cm}^6 \text{ mole}^{-2} \text{ sec}^{-1}$, from the results of Byron's⁸ interferometric studies of oxygen dissociation rates. Then since the temperature dependence of the recombination rate coefficient has been assumed as

$$k_R^{(2)} \sim T^{-2} \quad \text{(following Ref. 3)}$$

then

$$k_{R_0}^{(2)} = 1.16 \times 10^{15} \left(\frac{3000}{T_0'} \right)^2 \frac{\text{cm}^6}{\text{mole}^2 \text{sec}}$$

The remaining rate constants for the second or third bodies (β_i) were determined from ordinary collision theory. The form of the rate coefficient is, from collision theory⁹

$$k_f \sim \frac{KD'^2}{\sigma \sqrt{\mu_{ij}}} \left(\frac{\Theta'_{di}}{T'} \right)^\eta e^{-\frac{\Theta'_{di}}{T'}}$$

where K = steric factor

D' = average diameter of colliding particles

μ_{ij} = reduced mass of colliding particles (i and j)

$$= \frac{\mu_i + \mu_j}{\mu_i \mu_j}$$

η = constant

σ = symmetry number = 1 unlike particles
= 2 like particles

Θ'_{di} = characteristic dissociation temperature for species i

For oxygen (from Ref. 8),

$$\frac{k_f^{(1)}}{k_f^{(2)}} = \beta_1 = 35 \frac{T'}{\Theta'_0}$$

For N_2 and NO the steric factors and molecular diameters were assumed the same as for O_2 , whence

$$\frac{k_f^{(i)}}{k_f^{(2)}} = \frac{\sigma_2 \sqrt{\mu_{O_2, O_2}}}{\sigma_i \sqrt{\mu_{i, O_2}}}$$

then for N_2 ,

$$\frac{k_f^{(4)}}{k_f^{(2)}} = \beta_4 = 2.07$$

and NO ,

$$\frac{k_f^{(5)}}{k_f^{(2)}} = \beta_5 = 2.033$$

For N , the steric factor and atomic diameter were assumed to be the same as for O , so

$$\frac{k_f^{(3)}}{k_f^{(1)}} = 1.047$$

whence

$$\frac{k_f^{(3)}}{k_f^{(2)}} = \frac{k_f^{(3)}}{k_f^{(1)}} \frac{k_f^{(1)}}{k_f^{(2)}} = 1.047 \left(35 \frac{T'}{\theta_D'} \right) = 36.6 \frac{T'}{\theta_D'}$$

Relaxation lengths may then be calculated from Eq. (35).

Equation (36) is solved graphically to locate the point of freezing in the nozzle. The slopes $\left(\frac{d\gamma_o}{dA} \right)_\infty$ are obtained graphically from tangents to the curve of oxygen atom mass concentration as a function of area ratio. Calculated values of $(\gamma_o/r)_\infty$ and the value of $\left(\frac{d\gamma_o}{dA} \right)_\infty \left(\frac{dA}{dx} \right)$ are plotted over a range of area ratios. The area ratio at which these two curves cross is the location at which significant freezing has occurred. The frozen value of γ_o that can be expected for finite-rate flow is given by the value of $\gamma_{o\infty}$ which corresponds to this area ratio. Figure 18 shows the solution of Eq. (36) for the 5000°K, 100 atm. case.

B. Nozzle Geometry Considered

The nozzle geometries considered are given by the general expression $A = 1 + x^n$. These nozzles are symmetric about the throat for even values of the exponent n . In particular, two nozzle geometries have been considered which correspond to values of the exponent equal to 1 and 2. For $n = 1$, the

nozzle is a simple two-dimensional wedge-type nozzle with a sharp throat. The nozzle parameter, $\ell = L/a$, which enters the equation for the area ratio through the relationship $x = \frac{x'}{L/a}$, is given by one-half the throat height in centimeters divided by the tangent of half the total expansion angle. For $n = 2$, the nozzle is an axisymmetric hyperbolic nozzle with a smooth throat and which becomes a conical nozzle for large values of the area ratio, i.e., for nozzle areas not close to the throat. The nozzle geometry parameter is given by the throat radius in centimeters divided by the tangent of the asymptotic half angle.

The nozzle geometry parameter, L/a , represents the rapidity of the nozzle expansion. For a given throat size, small values of L/a correspond to large expansion angles and therefore to rapidly expanding nozzles. On the other hand, large values of the parameter L/a represent small expansion angles and accordingly correspond to gradually expanding nozzles. Many of the nozzles currently used for hypersonic flow research are designed with a value of L/a of about unity. In the present analyses, values of L/a of 0.1, 1, and 10 have been used to bracket the values of current design practice.

C. Results of Approximate Finite-Rate Solutions

The results of the relaxation length criterion calculations for flow in a hyperbolic axisymmetric and wedge-type nozzle are shown in Figs. 19, 20, and 21. The calculations were performed for stagnation temperatures of 4000, 5000, and 6000°K and stagnation pressures from 100 to 1000 atmospheres, for three values of the nozzle geometry parameter L/a (0.1, 1.0, 10).

Figure 19 shows the area ratio for freezing (A_f) plotted against the stagnation pressure for each stagnation temperature. In all cases freezing occurs fairly early in the nozzle ($A_f < 25$). Further, freezing in the nozzle is delayed by an increase in reservoir pressure, an increase in reservoir temperature, and by the use of gradually expanding nozzles (larger L/a).

The frozen degree of oxygen dissociation, α_f (the mass fraction of oxygen atoms to total oxygen) is given by

$$\alpha_f = \frac{\gamma_{O_f}}{\gamma_{O_f} + \gamma_{NO_f} + 2\gamma_{O_2f}}$$

where the subscript f denotes the value at oxygen freezing. The fraction of the stagnation enthalpy represented by frozen chemical energy is given by

$$\frac{H_f}{H_o} = \frac{\mu}{H_o} \left\{ \gamma_{O_f} \frac{\Theta_{DO_2}}{2} + \gamma_{N_f} \frac{\Theta_{DN_2}}{2} + \gamma_{NO_f} \left(\frac{\Theta_{DO_2} + \Theta_{DN_2}}{2} - \Theta_{DNO} \right) \right\}$$

where

$$\Theta_{Di} = \frac{\Theta'_{Di}}{T_o'}$$

The characteristic dissociation temperatures (Θ'_{Di}) assumed for oxygen, nitrogen and nitric oxide are 59,400°K, 113,200°K, and 75,500°K respectively^{7,10}.

The frozen degree of dissociation and H_f/H_o are plotted versus the stagnation pressure for each stagnation temperature in Figs. 20 and 21 respectively. Figure 20 shows that the frozen level of oxygen dissociation is appreciably reduced by an increase in stagnation pressure. This indicates that high stagnation pressures are required in order to reduce the static enthalpy loss through freezing at high levels of dissociation. Figure 21 shows, for example, that at a stagnation temperature of 6000°K and a stagnation pressure of 100 atm., as much as 20% of the total enthalpy may be frozen chemical energy. A wedge nozzle is superior to a hyperbolic nozzle of the same L/a because of a lower frozen degree of dissociation.

IV SOLUTIONS FOR EXPANSION AFTER FREEZING

For the simplified kinetic model of air, the exact finite-rate solution is approximated by assuming equilibrium flow up to the freezing point followed by a chemically-frozen flow downstream. This approximation was introduced by Bray¹ for the case of a pure diatomic gas. The area ratio for freezing as determined by the methods outlined in the last section, represents a unique point in the equilibrium solution. The equilibrium gasdynamic properties and composition at this area ratio are the initial values for the subsequent chemically-frozen expansion.

In the frozen expansion calculations, vibrational equilibrium is assumed with the energy taken as that for a quantized linear oscillator. The total gas enthalpy for translational, rotational and vibrational equilibrium, may then be written as

$$\begin{aligned}
 H = 3T + \mu \left\{ \gamma_{O_2} \left(\frac{T}{2} + \frac{\Theta_{VO_2}}{e^{\Theta_{VO_2}/T} - 1} \right) + \gamma_{N_2} \left(\frac{T}{2} + \frac{\Theta_{VN_2}}{e^{\Theta_{VN_2}/T} - 1} \right) \right. \\
 + \gamma_{NO} \left(\frac{T}{2} + \frac{\Theta_{VNO}}{e^{\Theta_{VNO}/T} - 1} \right) + \gamma_O \left(T + \frac{\Theta_{DO_2}}{2} \right) \\
 \left. + \gamma_N \left(T + \frac{\Theta_{DN_2}}{2} \right) + \gamma_{NO} \left(\frac{\Theta_{DO_2} + \Theta_{DN_2}}{2} - \Theta_{DNO} \right) \right\}
 \end{aligned} \quad (37)$$

where μ = molecular weight undissociated gas (28.85)

$$\Theta_{Vi} = \frac{\Theta'_{Vi}}{T_0'}$$

The characteristic vibrational temperatures (Θ'_{Vi}) used for oxygen, nitrogen and nitric oxide are 2300°K, 3390°K and 2740°K, respectively.

The method of solution is as follows. Since the gas composition is now frozen, all γ_i in Eq. (37) are known constants and the gas enthalpy may be simply determined for any specified temperature ratio T . The velocity is then calculated from Eq. (8).

The frozen expansion is isentropic (by virtue of zero reaction rates ¹⁴), hence

$$dH - \frac{dp}{\rho} = 0 \quad (38)$$

Then from Eq. (10), in which both $\sum \gamma_{i_0}$ and $\sum \gamma_i$ are now constant, we have

$$\int \frac{dH}{T} - \frac{\sum \gamma}{\sum \gamma_{i_0}} \ln T - \frac{\sum \gamma}{\sum \gamma_{i_0}} \ln \rho = \text{CONST.} \quad (39)$$

where $\sum \gamma = \sum \gamma_i$ at freezing

$$= (\gamma_{O_2} + \gamma_{N_2} + \gamma_{NO} + \gamma_O + \gamma_N)_f$$

Substituting Eq. (37) for H, the left side of Eq. (39) is evaluated to obtain

$$\begin{aligned} & \frac{1}{\sum \gamma} \left(\frac{3}{\mu} + \frac{\sum \gamma_2}{2} + \sum \gamma_1 - \frac{\sum \gamma''}{\mu \sum \gamma_{i_0}} \right) \ln T \\ & + \frac{\gamma_{O_2}}{\sum \gamma} \left(\frac{1}{T} \frac{\theta_{VO_2}}{e^{\theta_{VO_2}/T-1}} - \ln \frac{e^{\theta_{VO_2}/T-1}}{e^{\theta_{VO_2}/T}} \right) \\ & + \frac{\gamma_{N_2}}{\sum \gamma} \left(\frac{1}{T} \frac{\theta_{VN_2}}{e^{\theta_{VN_2}/T-1}} - \ln \frac{e^{\theta_{VN_2}/T-1}}{e^{\theta_{VN_2}/T}} \right) \\ & + \frac{\gamma_{NO}}{\sum \gamma} \left(\frac{1}{T} \frac{\theta_{VNO}}{e^{\theta_{VNO}/T-1}} - \ln \frac{e^{\theta_{VNO}/T-1}}{e^{\theta_{VNO}/T}} \right) - \frac{1}{\mu \sum \gamma_{i_0}} \ln \rho = \text{CONST.} \end{aligned} \quad (40)$$

where

$$\begin{aligned}\sum \gamma_2 &= (\gamma_{O_2} + \gamma'_{N_2} + \gamma_{NO})_f \\ \sum \gamma_1 &= (\gamma_O + \gamma'_N)_f\end{aligned}$$

Equation (40) relates ρ and τ explicitly for flow with frozen dissociation and equilibrium vibration. The constant is evaluated initially for equilibrium values at the point of freezing. The density is then calculated from Eq. (40) for all subsequent τ . The pressure is next obtained from Eq. (10). Also the area ratio from

$$A = \frac{(\rho u)_*}{\rho_{1L}}$$

and a Mach number is calculated at each step from

$$M = u / \sqrt{\left(\frac{dp}{d\rho} \right)_{\text{isentropic frozen}}}$$

The frozen air expansions were calculated on the CAL Datatron computer* and the results for all initial reservoir states and all nozzle configurations are tabulated in Tables 17 to 26. The parameters in each table commence with their values at freezing, i.e., the equilibrium values at the appropriate A_f . The behavior of the flow temperature, pressure, density, velocity and Mach number during expansion through the hyperbolic-type nozzles from each initial reservoir state is shown in Figs. 22-31. The corresponding equilibrium solutions are included to show the calculated location of freezing in the nozzle and to provide a basis of comparison for the frozen expansions. Only the results for the extreme nozzle shapes ($\ell = 0.1, 10$) are shown; data for the $\ell = 1.0$ hyperbolic nozzle flows, as well as for all the wedge-type nozzle flows may be extracted from the tables.

*The programming for Datatron machine computation was done by V.L. Widler, Head, Data Handling Section, Hypersonic Tunnel Department, Cornell Aeronautical Laboratory.

V FURTHER KINETIC CONSIDERATIONS

In the application of a relaxation length freezing criterion for air it has been assumed that the freezing of the oxygen kinetics, taken here as the governing chemical mechanism, implies the simultaneous freezing of all kinetics. Actually this assumption has little significance in the present cases since at the predicted freezing temperatures, little dissociation energy exists in the N and NO, and consequently, only a small amount of enthalpy is involved whether these species are assumed to freeze or not.

However, an analysis similar to that outlined in Sec. III may be performed for a different chemical species. The appropriate equations would all be of the same form as those given previously for oxygen. Now if relaxation length considerations were to indicate freezing of this species at area ratios earlier than those predicted in Sec. III on the basis of oxygen kinetics, the previous assumption of simultaneous freezing of all species would result in a greater loss of static enthalpy because of the freezing of oxygen at higher temperatures. This in turn could result in significant changes in some of the gas dynamic properties calculated in Sec. IV. It was considered expedient, therefore, to indicate the approximate solution to a specific case as dictated by the kinetics of some species other than oxygen.

The 6000°K, 100 atmosphere reservoir case was chosen as it involved the greatest amount of initial oxygen dissociation and is therefore the case most sensitive to earlier oxygen freezing. Furthermore, since the equilibrium solution exhibited a more rapid decrease in χ_N compared to χ_O (Fig. 9), the analysis was based on the nitrogen recombination kinetics. Citing only the results, the application of a relaxation length criterion indicated the freezing of nitrogen approximately at the throat for all nozzle configurations. Then, consistent with the previous assumption, the nitric oxide and oxygen were frozen simultaneously with nitrogen and the resultant frozen expansion is indicated by the dashed profiles in Fig. 29.

It is seen that the calculated frozen expansions determined by the nitrogen kinetics may differ appreciably from those determined by oxygen kinetics. The

difference is mainly dependent on the location of oxygen freezing. However, an argument against the freezing of nitrogen is provided by the nitric oxide "shuffle" reactions. The reactions



provide two fast two-body processes for equilibration of $N - N_2$ and $N - NO$ which have not been considered in the nitrogen three-body processes discussed above.

VI CONCLUSIONS

A method has been described to estimate the effects of finite chemical reaction rates on the one-dimensional expansion of air in hypersonic nozzles. Numerical solutions have been computed for reservoir temperatures from 4000°K to 6000°K and reservoir pressures from 100 to 1000 atmospheres. The calculations have considered a range of wedge and axisymmetric nozzle shapes. The air was assumed to be at rest and in chemical equilibrium at the reservoir state. The flow parameters and gas composition in the infinite-rate equilibrium expansion from each reservoir state were calculated, and a relaxation length criterion, based on the equilibrium solution, was used to indicate the approximate location of freezing for finite-rate airflow in each nozzle configuration. The resultant frozen air expansions have been presented in tabular and graphical form.

The results of the analysis have shown that in all cases, freezing occurred fairly early in the nozzle downstream of the throat, at area ratios less than about 25 (Fig. 19). As a result of such early freezing the frozen degree of dissociation may be quite large (Fig. 20) representing an appreciable loss in the static enthalpy. It is seen, for example (Fig. 21), that in the high temperature, low pressure cases, the unavailable enthalpy owing to freezing may be as much as 19% of the stagnation enthalpy.

The following features are evident in the results. The effect of the static enthalpy loss through freezing is to reduce the pressure, velocity and temperature at a particular nozzle location (area ratio) from the corresponding equilibrium values, and to increase the density and Mach number. While the change in temperature and Mach number may be considerable, the density and velocity are relatively unaffected by freezing.

As expected, the effects of freezing are minimized for larger L/a values, i.e., by the use of gradually expanding nozzles. In practice, however, limitations will be imposed by boundary layer growth in the nozzle and a compromise will be required between maximizing L/a and maintaining a reasonable nozzle size. In this respect the wedge nozzle is superior to the hyperbolic nozzle for

a given L/a by virtue of a lower frozen degree of dissociation of the flow (later freezing). It should be noted that the effect of a change in the nozzle L/a parameter also reflects the effect of a change in the kinetic rates, in particular $k_{R_0}^{(2)}$. Since L/a and the recombination rate constant $k_{R_0}^{(2)}$ occur in the rate equation as a product (in ψ), a change in L/a is comparable to a similar change in $k_{R_0}^{(2)}$, the product remaining the same. This factor permits a simple compensation for discrepancies in existing rate constant data. Also noticeable in the results is the reduction in the frozen degree of dissociation with increase in reservoir pressure. This indicates that high reservoir pressures are required in order to reduce the static enthalpy loss through freezing in the operation of hypersonic wind tunnels at high stagnation temperatures. This requirement is compatible with aerodynamic requirements for testing at very high Mach numbers.

The present nonequilibrium airflow calculations are meant as approximate estimates of the gross features of particular finite-rate nozzle flows. In the interests of simplicity several effects have been omitted from the approximations used, notably the neglect of the nitrogen and nitric oxide reactions at finite rates in the simplified kinetic model of air. However, since the assumption of both zero and infinite rates for reactions involving these species resulted in relatively small changes in the frozen levels of dissociation (Sec. 3.1), the calculated expansions from each reservoir state should represent realistic approximations to actual gas behavior in the nozzle. Ultimate verification must come from comparisons with exact solutions to more complex kinetic models. Such a program is presently underway at CAL.

APPENDIX A NEWTON-RAPHSON METHOD^{11,12,13} FOR SOLVING THE EQUILIBRIUM PROBLEM

The use of the Newton-Raphson method for the iterative solution of these simultaneous nonlinear algebraic equations was originally suggested by Brinkley⁶ for the determination of equilibrium composition at a given temperature and pressure. Brinkley's method is directly applicable to the computation of the stagnation state; however, certain modifications are introduced for computing composition and pressure at a given entropy for various temperatures along the expansion. Only the modified form is presented here because the essential features of the Brinkley method are contained therein. Equations (23) and (27) are expressed as

$$F_j = 0 \quad (j = 1, 2, \dots, c, c+1) \quad (\text{A-1})$$

where

$$F_j = q_j + \sum_{l=c+1}^S \left[q_j (\nu_l - 1) - \nu_{lj} \right] X_l - X_j \quad (\text{A-2})$$

$(j = 1, 2, \dots, c)$

$$F_{c+1} = \sum_{i=1}^S X_i (S_i^\circ - \mu_i s_0 - \ln X_i) - \ln p' \quad (\text{A-3})$$

For given S_0 and T' , values of X_j and p' (the $(c+1)$ st independent variable) must be guessed such that Equations (A-1) are satisfied. X_l are related to X_j by Equation (24).

$$X_l = K_{x_l} \prod_{j=1}^c X_j^{\nu_{lj}} \quad (l = c+1, c+2, \dots, S) \quad \text{Repeat (24)}$$

The Newton-Raphson method provides a means of improving successive guesses by utilizing a Taylor expansion of F_j about X_j and p' to first order terms. This leads to a set of fractional corrections $h_k^{(r)}$ ($k=1, 2, \dots, c, c+1$) applicable to the X_j and p' values obtained at the r th iteration. The improved $(r+1)$ st values of the independent variables are found from

$$X_k^{(r+1)} = X_k^{(r)} (1 + h_k^{(r)}) \quad (k=1, 2, \dots, c) \quad (\text{A-4})$$

and

$$p'^{(r+1)} = p'^{(r)} (1 + h_{c+1}^{(r)}) \quad (\text{A-5})$$

The set of h_k is computed from a set of linear equations expressed in matrix notation by

$$\begin{bmatrix} A_{jk}^{(r)} \end{bmatrix} \begin{bmatrix} h_k^{(r)} \end{bmatrix} = \begin{bmatrix} F_j^{(r)} \end{bmatrix} \quad \begin{matrix} (j=1, 2, \dots, c, c+1) \\ (k=1, 2, \dots, c, c+1) \end{matrix} \quad (\text{A-6})$$

$h_k^{(r)}$ and $F_j^{(r)}$ are column vectors

$$A_{jk} = X_j \delta_{jk} + \sum_{\ell=c+1}^s \nu_{\ell k} [\nu_{\ell j} - q_j (\nu_{\ell} - 1)] X_{\ell} \quad (\text{A-7})$$

$$\delta_{jk} = 0 \quad \text{for } j \neq k \quad (j=1, 2, \dots, c)$$

$$\delta_{jk} = 1 \quad \text{for } j = k \quad (k=1, 2, \dots, c)$$

$$A_{j,c+1} = \sum_{\ell=c+1}^s (\nu_{\ell} - 1) [\nu_{\ell j} - q_j (\nu_{\ell} - 1)] X_{\ell} \quad (j=1, 2, \dots, c) \quad (\text{A-8})$$

$$A_{c+1,k} = \left\{ \begin{aligned} &[1 - (S_k^0 - \mu_k s_0 - \ln X_k)] X_k + \\ &\sum_{\ell=c+1}^s \nu_{\ell k} [1 - (S_{\ell}^0 - \mu_{\ell} s_0 - \ln X_{\ell})] X_{\ell} \end{aligned} \right\} \quad (k=1, 2, \dots, c) \quad (\text{A-9})$$

$$A_{c+1,c+1} = 1 + \sum_{\ell=c+1}^s (\nu_{\ell} - 1) [1 - (S_{\ell}^0 - \mu_{\ell} s_0 - \ln X_{\ell})] X_{\ell} \quad (\text{A-10})$$

The $(r+1)$ st iteration is started by calculating $h_k^{(r)}$ from Equation (A-6) using the formulas in Equations (24), (A-2), (A-3), (A-7), (A-8), (A-9), and (A-10). Improved values of λ_j and p' are found from Equations (A-4) and (A-5), and improved values of λ_ℓ are found from Equation (24) using $\lambda_j^{(r+1)}$ and $p'^{(r+1)}$. The cycle begins again with the computation of $[h_k^{(r+1)}]$ from Equation (A-6). When all h_k become equal to or smaller than a prescribed maximum fractional error ϵ_h , the computation is terminated.

REFERENCES

1. Bray, K. N. C. Atomic Recombination in a Hypersonic Wind Tunnel Nozzle, J. Fluid Mechs. Vol. 6, Pt. 1, July 1959
2. Hall, J. G. Dissociation Nonequilibrium in Hypersonic Nozzle Flow. Paper presented at Symposium on Thermodynamics of Jet and Rocket Propulsion, A.I.Ch.E. National Meeting, 1959
3. Hall, J. G. and Russo, A. L. Studies of Chemical Nonequilibrium in Hypersonic Nozzle Flows. C.A.L. Report No. AD-1118-A-6, AFOSR TN-59-1090, November 1959
4. Wegener, P. P. Experiments on the Departure from Chemical Equilibrium in a Supersonic Flow. ARS Journal, 30, 4, April 1960
5. Heims, S. P. Effect of Oxygen Recombination on One-Dimensional Flow at High Mach Numbers. NACA TN 4144, January 1958
6. Brinkley, S. R. Computational Methods in Combustion Calculations. Article 1. C. in Combustion Processes, Vol. II of Princeton Series, pp 64-98 (1956)
7. Logan, J. G. and Treanor, C. E. Tables of Thermodynamic Properties of Air from 3000°K to 10,000°K at Intervals of 100°K. C.A.L. Report No. BE-1007-A-3, January 1957
8. Byron, S. B. Interferometric Measurement of the Rate of Dissociation of Oxygen Heated by Strong Shock Waves. Cornell University Report, 1958
9. Fowler and Guggenheim Statistical Thermodynamics. Cambridge University Press, 1952
10. Stupochenko, et al Thermodynamic Properties of Air in the Temperature Interval from 1000 to 12000°K and the Pressure Intervals from 0.001 to 1000 Atmosphere. ARS Journal Supplement, 30, 1, January 1960
11. Newton, I. Method of Fluxions (1669)

T	ρ	P	\dot{u} (M/sec)	M	A	T_H	T_0	T_{H_2}	T_{O_2}	T_A	T_N	T_O	T_{NO}	T_{NO^+}	T_{N^+}	T_{O^+}	T_{O^-}	T_{e^-}
1.00	1.00 01	1.00 01	0.000 00	0.00 00	1.501 01	3219-03	1640-05	6596-03	3393-02	8772-14	2018-10	8092-08	1078-16	1250-13	1078-16	1250-13	4257-08	3655-08
.98	910 00	.691 01	1.889 04	.437 00	1.501 01	3219-03	1291-05	5747-03	3250-02	4737-14	1378-10	6105-08	4866-17	6948-14	4866-17	6948-14	3101-06	3018-08
.96	826 00	.722 00	2.356 04	.624 00	1.170 01	3219-03	1005-05	5217-03	3104-02	2493-14	9245-11	4350-08	2141-17	3762-14	2141-17	3762-14	2221-06	2335-08
.94	749 00	.793 00	2.920 04	.772 00	1.054 01	3219-03	7735-06	4706-03	2950-02	1279-14	8056-11	3346-06	9049-18	1961-14	9049-18	1961-14	1571-06	1783-08
.92	678 00	.842 00	3.598 04	.900 00	1.010 01	3219-03	5875-06	4218-03	2807-02	6369-15	3941-11	2429-08	3679-18	1012-14	3679-18	1012-14	1090-08	1343-06
.90	613 00	.890 00	4.381 04	1.02 00	1.000 01	3219-03	4401-06	3753-03	2657-02	3071-15	2496-11	1736-08	1635-18	5011-15	1635-18	5011-15	7421-09	9968-09
.88	553 00	.935 00	5.271 04	1.12 01	1.013 01	3219-03	3249-06	3315-03	2508-02	1438-15	1546-11	1222-08	5355-19	2397-15	5355-19	2397-15	4933-09	7279-09
.86	496 00	.978 00	6.266 04	1.23 01	1.043 01	3219-03	2360-06	2904-03	2358-02	8500-16	9357-12	8452-09	1908-19	1106-15	1908-19	1106-15	3238-09	5226-09
.84	440 00	1.020 01	7.373 04	1.34 01	1.066 01	3219-03	1686-06	2521-03	2208-02	2828-16	5521-12	5743-09	6448-20	4904-16	6448-20	4904-16	2068-09	3683-07
.82	384 00	1.061 01	8.597 04	1.44 01	1.114 01	3219-03	1182-06	2168-03	2061-02	1182-16	3171-12	3827-09	2068-20	4066-16	2068-20	4066-16	1286-09	2544-09
.80	328 00	1.099 01	9.944 04	1.54 01	1.171 01	3219-03	8128-07	1845-03	1911-02	4733-17	1765-12	2498-09	6240-21	4476-17	6240-21	4476-17	7797-10	1720-09
.78	272 00	1.135 01	11.411 04	1.64 01	1.231 01	3219-03	5469-07	1553-03	1771-02	1809-17	9567-13	1594-09	1770-21	3262-17	1770-21	3262-17	4592-10	1136-09
.76	216 00	1.169 01	12.904 04	1.74 01	1.291 01	3219-03	3586-07	1391-03	1629-02	1678-18	5003-13	9927-10	4666-22	1205-17	4666-22	1205-17	4592-10	7310-10
.74	160 00	1.203 01	14.411 04	1.84 01	1.351 01	3219-03	2305-07	1059-03	1491-02	2266-18	2523-13	6023-10	1153-22	1460-16	1153-22	1460-16	1441-10	4579-10
.72	104 00	1.237 01	15.922 04	1.94 01	1.411 01	3219-03	1438-07	870-04	1358-02	2243-19	1233-13	3552-10	2816-23	1363-18	2816-23	1363-18	7698-11	2764-10
.70	48 00	1.271 01	17.433 04	2.04 01	1.471 01	3219-03	8710-08	6826-04	1228-02	2243-19	5682-14	2051-10	5443-24	1157-19	5443-24	1157-19	5938-11	1636-10
.68	12 00	1.305 01	18.944 04	2.14 01	1.531 01	3219-03	5107-08	5346-04	1101-02	6378-20	2502-14	1123-10	1030-24	1178-19	1030-24	1178-19	1928-11	9368-11
.66	160 00	1.339 01	20.455 04	2.24 01	1.591 01	3219-03	2890-08	4311-04	1101-02	1680-20	1063-14	5987-11	1757-25	3082-20	1757-25	3082-20	5997-12	5069-11
.64	104 00	1.373 01	21.966 04	2.34 01	1.651 01	3219-03	1574-08	3097-04	8877-03	4077-21	4239-15	3066-11	2675-26	2390-21	2675-26	2390-21	5997-12	4667-11
.62	48 00	1.407 01	23.477 04	2.44 01	1.711 01	3219-03	8210-09	2283-04	7603-03	9028-22	1591-15	1502-11	3597-27	1610-21	3597-27	1610-21	1673-12	1355-11
.60	12 00	1.441 01	24.988 04	2.54 01	1.771 01	3219-03	4088-09	1642-04	6598-03	1609-22	5586-16	7012-12	4415-26	3154-22	4415-26	3154-22	6593-13	6534-12
.58	160 00	1.475 01	26.499 04	2.64 01	1.831 01	3219-03	1933-09	1149-04	5665-03	1609-22	5586-16	7012-12	4415-26	3154-22	4415-26	3154-22	6593-13	6534-12
.56	104 00	1.509 01	28.010 04	2.74 01	1.891 01	3219-03	8659-10	7812-05	4807-03	1609-22	5586-16	7012-12	4415-26	3154-22	4415-26	3154-22	6593-13	6534-12
.54	48 00	1.543 01	29.521 04	2.84 01	1.951 01	3219-03	8659-10	7812-05	4807-03	1609-22	5586-16	7012-12	4415-26	3154-22	4415-26	3154-22	6593-13	6534-12
.52	160 00	1.577 01	31.032 04	2.94 01	2.011 01	3219-03	1410-10	3255-05	3324-03	1609-22	5586-16	7012-12	4415-26	3154-22	4415-26	3154-22	6593-13	6534-12
.50	104 00	1.611 01	32.543 04	3.04 01	2.071 01	3219-03	5082-11	1979-05	2700-03	1609-22	5586-16	7012-12	4415-26	3154-22	4415-26	3154-22	6593-13	6534-12
.48	48 00	1.645 01	34.054 04	3.14 01	2.131 01	3219-03	1674-05	1149-05	2154-03	1609-22	5586-16	7012-12	4415-26	3154-22	4415-26	3154-22	6593-13	6534-12
.46	160 00	1.679 01	35.565 04	3.24 01	2.191 01	3219-03	4982-12	6326-06	1683-03	1609-22	5586-16	7012-12	4415-26	3154-22	4415-26	3154-22	6593-13	6534-12
.44	104 00	1.713 01	37.076 04	3.34 01	2.251 01	3219-03	1321-12	3282-06	1288-03	1609-22	5586-16	7012-12	4415-26	3154-22	4415-26	3154-22	6593-13	6534-12
.42	48 00	1.747 01	38.587 04	3.44 01	2.311 01	3219-03	3071-13	1590-06	9562-04	1609-22	5586-16	7012-12	4415-26	3154-22	4415-26	3154-22	6593-13	6534-12
.40	160 00	1.781 01	40.098 04	3.54 01	2.371 01	3219-03	1613-14	6133-07	6901-04	1609-22	5586-16	7012-12	4415-26	3154-22	4415-26	3154-22	6593-13	6534-12
.38	104 00	1.815 01	41.609 04	3.64 01	2.431 01	3219-03	1028-14	7214-07	4810-04	1609-22	5586-16	7012-12	4415-26	3154-22	4415-26	3154-22	6593-13	6534-12
.36	48 00	1.849 01	43.120 04	3.74 01	2.491 01	3219-03	1402-15	1072-07	3219-04	1609-22	5586-16	7012-12	4415-26	3154-22	4415-26	3154-22	6593-13	6534-12
.34	160 00	1.883 01	44.631 04	3.84 01	2.551 01	3219-03	1503-16	3478-08	2034-04	1609-22	5586-16	7012-12	4415-26	3154-22	4415-26	3154-22	6593-13	6534-12
.32	104 00	1.917 01	46.142 04	3.94 01	2.611 01	3219-03	1209-17	9729-09	1238-04	1609-22	5586-16	7012-12	4415-26	3154-22	4415-26	3154-22	6593-13	6534-12
.30	48 00	1.951 01	47.653 04	4.04 01	2.671 01	3219-03	6893-19	2377-09	4976-05	1609-22	5586-16	7012-12	4415-26	3154-22	4415-26	3154-22	6593-13	6534-12
.28	160 00	1.985 01	49.164 04	4.14 01	2.731 01	3219-03	2386-20	4288-10	3621-05	1609-22	5586-16	7012-12	4415-26	3154-22	4415-26	3154-22	6593-13	6534-12
.26	104 00	2.019 01	50.675 04	4.24 01	2.791 01	3219-03	5793-22	6183-11	1699-05	1609-22	5586-16	7012-12	4415-26	3154-22	4415-26	3154-22	6593-13	6534-12
.24	48 00	2.053 01	52.186 04	4.34 01	2.851 01	3219-03	6812-24	6379-12	7026-06	1609-22	5586-16	7012-12	4415-26	3154-22	4415-26	3154-22	6593-13	6534-12
.22	160 00	2.087 01	53.697 04	4.44 01	2.911 01	3219-03	3524-26	4297-13	2476-06	1609-22	5586-16	7012-12	4415-26	3154-22	4415-26	3154-22	6593-13	6534-12
.20	104 00	2.121 01	55.208 04	4.54 01	2.971 01	3219-03	1661-29	1661-14	7085-07	1609-22	5586-16	7012-12	4415-26	3154-22	4415-26	3154-22	6593-13	6534-12
.18	48 00	2.155 01	56.719 04	4.64 01	3.031 01	3219-03	2681-32	3060-16	1537-07	1609-22	5586-16	7012-12	4415-26	3154-22	4415-26	3154-22	6593-13	6534-12
.16	160 00	2.189 01	58.230 04	4.74 01	3.091 01	3219-03	1615-36	2031-18	2278-08	1609-22	5586-16	7012-12	4415-26	3154-22	4415-26	3154-22	6593-13	6534-12
.14	104 00	2.223 01	59.741 04	4.84 01	3.151 01	3219-03	3332-21	3332-21	1963-09	1609-22	5586-16	7012-12	4415-26	3154-22	4415-26	3154-22	6593-13	6534-12
.12	48 00	2.257 01	61.252 04	4.94 01	3.211 01	3219-03	5592-25	5592-25	1499-11	1609-22	5586-16	7012-12	4415-26	3154-22	4415-26	3154-22	6593-13	6534-12
.10	160 00	2.291 01	62.763 04	5.04 01	3.271 01	3219-03	2768-30	2768-30	7815-13	1609-22	5586-16	7012-12	4415-26	3154-22	4415-26	3154-22	6593-13	6534-12

Table 3.

EQUILIBRIUM AIRFLOW EXPANSION

$T_0' = 4000^\circ\text{K}$, $p_0' = 1000 \text{ atm.}$, $\rho_0' = 8.744 \times 10^{-2} \text{ gm/cm}^3$,
 $H_0 = 4.599, (\rho u)^* = .6547 \text{ at } T = .903, \bar{\mu}_0 = 28.70 \text{ gm/mole}$

[illegible]

EQUILIBRIUM AIRFLOW EXPANSION

Table 5.

EQUILIBRIUM AIRFLOW EXPANSION

$$T_o' = 5000^\circ\text{K}, p_o' = 200 \text{ atm.}, \rho_o' = 1.317 \times 10^{-2} \text{ gm/cm}^3, \\ H_o = 5.436, (\mu)_* = .6729 \text{ at } T = .915, \bar{\mu}_o = 27.01 \text{ gm/mole}$$

[illegible]

$T_o' = 5000^\circ\text{K}$, $p_o' = 300 \text{ atm.}$, $\rho_o' = 1.995 \times 10^{-2} \text{ gm/cm}^3$,
 $H_o = 5.313$, $(\rho u)_* = .6697 \text{ at } T = .913$, $\bar{\mu}_o = 27.28 \text{ gm/mole}$

21

[illegible]

Table 7.

$$T_O' = 5000^\circ\text{K}, p_O' = 1000 \text{ atm}, \rho_O' = 6.810 \times 10^{-2} \text{ gm/cm}^3, \\ H_O' = 5.026, (\rho u)_* = .6628 \text{ at } T = .909, \bar{\mu}_O = 27.94 \text{ gm/mole}$$

[illegible]

Table 8.

$$T_o' = 6000^\circ\text{K}, p_o' = 100 \text{ atm.}, \rho_o' = 5.078 \times 10^{-3} \text{ gm/cm}^3$$

T	ρ	P	\dot{u} (ft/sec.)	M	A	Th_2	To_2	Ta	Tn	To	Tno	Tn_2	To_2	Tno'	Tn'	To'	Tn'	To'	Tn'
1.00	1.00	1.00	0.000	0.00	0.00	2.475-01	1.278-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.98	.906	.866	.2212	.00	.446	2.475-01	1.357-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.96	.823	.783	.4130	.00	.892	2.475-01	1.437-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.94	.743	.693	.6037	.00	1.337	2.475-01	1.517-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.92	.669	.609	.7937	.00	1.782	2.475-01	1.597-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.90	.600	.530	.9836	.00	2.227	2.475-01	1.677-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.88	.539	.469	1.1735	.00	2.672	2.475-01	1.757-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.86	.477	.407	1.3634	.00	3.117	2.475-01	1.837-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.84	.423	.353	1.5533	.00	3.562	2.475-01	1.917-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.82	.373	.303	1.7432	.00	3.999	2.475-01	2.000-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.80	.327	.257	1.9331	.00	4.436	2.475-01	2.082-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.78	.285	.215	2.1230	.00	4.873	2.475-01	2.165-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.76	.246	.180	2.3129	.00	5.310	2.475-01	2.248-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.74	.212	.150	2.5028	.00	5.747	2.475-01	2.331-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.72	.181	.130	2.6927	.00	6.184	2.475-01	2.414-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.70	.153	.110	2.8826	.00	6.621	2.475-01	2.497-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.68	.129	.090	3.0725	.00	7.058	2.475-01	2.580-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.66	.107	.070	3.2624	.00	7.495	2.475-01	2.663-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.64	.089	.051	3.4523	.00	7.932	2.475-01	2.746-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.62	.073	.040	3.6422	.00	8.369	2.475-01	2.829-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.60	.058	.030	3.8321	.00	8.806	2.475-01	2.912-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.58	.043	.020	4.0220	.00	9.243	2.475-01	2.995-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.56	.030	.011	4.2119	.00	9.680	2.475-01	3.078-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.54	.020	.006	4.4018	.00	1.012	2.475-01	3.161-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.52	.014	.003	4.5917	.00	1.057	2.475-01	3.244-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.50	.009	.002	4.7816	.00	1.102	2.475-01	3.327-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.48	.006	.001	4.9715	.00	1.147	2.475-01	3.410-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.46	.004	.000	5.1614	.00	1.192	2.475-01	3.493-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.44	.003	.000	5.3513	.00	1.237	2.475-01	3.576-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.42	.002	.000	5.5412	.00	1.282	2.475-01	3.659-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.40	.001	.000	5.7311	.00	1.327	2.475-01	3.742-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.38	.000	.000	5.9210	.00	1.372	2.475-01	3.825-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.36	.000	.000	6.1109	.00	1.417	2.475-01	3.908-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.34	.000	.000	6.3008	.00	1.462	2.475-01	3.991-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.32	.000	.000	6.4907	.00	1.507	2.475-01	4.074-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.30	.000	.000	6.6806	.00	1.552	2.475-01	4.157-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.28	.000	.000	6.8705	.00	1.597	2.475-01	4.240-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.26	.000	.000	7.0604	.00	1.642	2.475-01	4.323-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.24	.000	.000	7.2503	.00	1.687	2.475-01	4.406-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.22	.000	.000	7.4402	.00	1.732	2.475-01	4.489-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.20	.000	.000	7.6301	.00	1.777	2.475-01	4.572-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.18	.000	.000	7.8200	.00	1.822	2.475-01	4.655-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.16	.000	.000	8.0099	.00	1.867	2.475-01	4.738-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.14	.000	.000	8.1998	.00	1.912	2.475-01	4.821-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.12	.000	.000	8.3897	.00	1.957	2.475-01	4.904-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.10	.000	.000	8.5796	.00	2.002	2.475-01	4.987-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.08	.000	.000	8.7695	.00	2.047	2.475-01	5.070-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.06	.000	.000	8.9594	.00	2.092	2.475-01	5.153-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.04	.000	.000	9.1493	.00	2.137	2.475-01	5.236-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05
.02	.000	.000	9.3392	.00	2.182	2.475-01	5.319-02	3.219-03	1.191-03	7.793-02	3.993-02	9.340-09	1.359-07	1.985-05	9.463-10	2.141-06	9.463-10	2.141-06	1.482-05

Table 9.

EQUILIBRIUM AIRFLOW EXPANSION

$T_0' = 6000^\circ\text{K}$, $p_0' = 300 \text{ atm.}$, $\rho_0' = 1.574 \times 10^{-2} \text{ gm/cm}^3$,
 $H_0 = 5.850$, $(\rho u)_* = .692$

Table 11. EQUILIBRIUM AIRFLOW EXPANSION

$T_o' = 7000^\circ\text{K}$, $p_o' = 100 \text{ atm.}$, $\rho_o' = 4.133 \times 10^{-3} \text{ gm/cm}^3$,
 $H_o = 6.805$, $(\rho u)_\infty = .7156 \text{ at } T = .920$, $\bar{\mu}_o = 23.74 \text{ gm/mole}$

δ	P	μ' (H/sec)	M	A	TH_2	TO_2	TA	TH	TO	TH_0	TH_3	TO_3	TH_0'	TH'	TO'	TH''	TO''	TH'''	TO'''
1.00	1.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00
9.9	9.90	0.376	0.4	1.500	1.0	9.907	3.119	9.914	8.058	9.914	9.924	5.138	5.485	9.924	1.701	5.485	9.924	1.701	5.485
9.8	9.80	0.376	0.4	1.500	1.0	10.39	3.119	9.914	7.980	9.924	9.924	4.367	4.790	9.924	1.386	4.790	9.924	1.386	4.790
9.7	9.70	0.376	0.4	1.500	1.0	10.87	3.119	9.914	7.945	9.924	9.924	3.684	4.159	9.924	0.964	4.159	9.924	0.964	4.159
9.6	9.60	0.376	0.4	1.500	1.0	11.35	3.119	9.914	7.910	9.924	9.924	2.998	3.580	9.924	0.543	3.580	9.924	0.543	3.580
9.5	9.50	0.376	0.4	1.500	1.0	11.83	3.119	9.914	7.875	9.924	9.924	2.312	2.969	9.924	0.122	2.969	9.924	0.122	2.969
9.4	9.40	0.376	0.4	1.500	1.0	12.31	3.119	9.914	7.840	9.924	9.924	1.626	2.556	9.924	0.305	2.556	9.924	0.305	2.556
9.3	9.30	0.376	0.4	1.500	1.0	12.79	3.119	9.914	7.805	9.924	9.924	0.940	2.143	9.924	0.488	2.143	9.924	0.488	2.143
9.2	9.20	0.376	0.4	1.500	1.0	13.27	3.119	9.914	7.770	9.924	9.924	0.254	1.730	9.924	0.671	1.730	9.924	0.671	1.730
9.1	9.10	0.376	0.4	1.500	1.0	13.75	3.119	9.914	7.735	9.924	9.924	0.568	1.317	9.924	0.854	1.317	9.924	0.854	1.317
9.0	9.00	0.376	0.4	1.500	1.0	14.23	3.119	9.914	7.700	9.924	9.924	0.882	0.904	9.924	1.037	0.904	9.924	1.037	0.904
8.9	8.90	0.376	0.4	1.500	1.0	14.71	3.119	9.914	7.665	9.924	9.924	1.196	0.491	9.924	1.220	0.491	9.924	1.220	0.491
8.8	8.80	0.376	0.4	1.500	1.0	15.19	3.119	9.914	7.630	9.924	9.924	1.510	0.078	9.924	1.403	0.078	9.924	1.403	0.078
8.7	8.70	0.376	0.4	1.500	1.0	15.67	3.119	9.914	7.595	9.924	9.924	1.824	0.665	9.924	1.586	0.665	9.924	1.586	0.665
8.6	8.60	0.376	0.4	1.500	1.0	16.15	3.119	9.914	7.560	9.924	9.924	2.138	1.252	9.924	1.769	1.252	9.924	1.769	1.252
8.5	8.50	0.376	0.4	1.500	1.0	16.63	3.119	9.914	7.525	9.924	9.924	2.452	1.839	9.924	1.952	1.839	9.924	1.952	1.839
8.4	8.40	0.376	0.4	1.500	1.0	17.11	3.119	9.914	7.490	9.924	9.924	2.766	2.426	9.924	2.135	2.426	9.924	2.135	2.426
8.3	8.30	0.376	0.4	1.500	1.0	17.59	3.119	9.914	7.455	9.924	9.924	3.080	3.013	9.924	2.318	3.013	9.924	2.318	3.013
8.2	8.20	0.376	0.4	1.500	1.0	18.07	3.119	9.914	7.420	9.924	9.924	3.394	3.600	9.924	2.501	3.600	9.924	2.501	3.600
8.1	8.10	0.376	0.4	1.500	1.0	18.55	3.119	9.914	7.385	9.									

Table 13.

EQUILIBRIUM AIRFLOW EXPANSION

$$T_0 = 7000^\circ\text{K}, p_0 = 1000 \text{ atm.}, \rho_0 = 4.460 \times 10^{-2} \text{ gm/cm}^3,$$

$$H_0 = 5.349, (pu)_* = .6962 \text{ at } T = .904, \bar{\mu}_0 = 25.62 \text{ gm/mole}$$

Table 15.

$$T_O' = 8000^\circ\text{K}, p_O' = 300 \text{ atm.}, \rho_O' = 1.060 \times 10^{-2} \text{ gm/cm}^3, \\ H_O = 6.999, (\rho u)_* = .7234 \text{ at } T = .924, \bar{\mu}_O = 23.20 \text{ gm/mole}$$
[illegible]

[illegible]
$$T_o' = 8000^\circ\text{K}, p_o' = 1000 \text{ atm.}, \rho_o' = 3.722 \times 10^{-2} \text{ gm/cm}^3$$

$$H_o = 6.289, (\rho u)_* = .7105 \text{ at } T = .910, \bar{\mu}_o = 24.43 \text{ gm/mole}$$

$l = 0.1$ $l = 1.0$ $l = 10$

T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A
(ft/sec.)														
84.4	400 00	5101 04	141 01	1134 01	73	196 00	140 00	196 01	1795 01	626	985-01	7545 04	243 01	3113 01
84	394 00	5143 04	141 01	1143 01	72	187 00	132 00	198 01	1858 01	62	953-01	7587 04	244 01	3100 01
82	392 00	5293 04	149 01	1193 01	70	170 00	116 00	205 01	1998 01	58	852-01	7729 04	252 01	3117 01
80	393 00	5553 04	156 01	1251 01	68	154 00	102 00	213 01	2156 01	56	760-01	7867 04	261 01	3071 01
78	396 00	5747 04	164 01	1317 01	66	139 00	898-01	221 01	2335 01	54	675-01	8003 04	270 01	3086 01
76	390 00	5934 04	171 01	1394 01	64	125 00	784-01	229 01	2540 01	52	597-01	8136 04	280 01	3078 01
74	385 00	6116 04	179 01	1481 01	62	113 00	683-01	238 01	2773 01	50	525-01	8266 04	289 01	3085 01
72	380 00	6292 04	186 01	1580 01	60	101 00	591-01	246 01	3040 01	48	461-01	8395 04	300 01	3085 01
70	375 00	6463 04	194 01	1692 01	58	898-01	503-01	255 01	3347 01	46	402-01	8521 04	310 01	3078 01
68	370 00	6629 04	202 01	1820 01	56	797-01	433-01	264 01	3701 01	44	349-01	8664 04	321 01	3070 01
66	365 00	6791 04	210 01	1965 01	54	706-01	375-01	274 01	4111 01	42	301-01	8816 04	333 01	3070 01
64	360 00	6949 04	218 01	2131 01	52	622-01	318-01	283 01	4589 01	40	258-01	8985 04	345 01	3062 01
62	355 00	7104 04	226 01	2320 01	50	545-01	267-01	293 01	5149 01	38	220-01	9003 04	358 01	3052 01
60	350 00	7254 04	235 01	2536 01	48	476-01	223-01	304 01	5809 01	36	186-01	9118 04	372 01	3032 01
58	345 00	7402 04	244 01	2785 01	46	413-01	188-01	315 01	6590 01	34	156-01	9231 04	387 01	3012 01
56	340 00	7546 04	253 01	3072 01	44	357-01	158-01	327 01	7521 01	32	130-01	9342 04	402 01	2992 01
54	335 00	7687 04	262 01	3405 01	42	306-01	128-01	339 01	8640 01	30	107-01	9451 04	419 01	2972 01
52	330 00	7826 04	272 01	3791 01	40	261-01	102-01	352 01	9994 01	28	878-02	9567 04	437 01	2952 01
50	325 00	7962 04	282 01	4244 01	38	221-01	82-02	365 01	1165 02	26	709-02	9684 04	456 01	2932 01
48	320 00	8095 04	293 01	4776 01	36	186-01	653-02	380 01	1368 02	24	566-02	9864 04	478 01	2912 01
46	315 00	8225 04	304 01	5406 01	34	155-01	513-02	395 01	1622 02	22	445-02	9962 04	501 01	2892 01
44	310 00	8353 04	315 01	6156 01	32	128-01	400-02	412 01	1941 02	20	344-02	1006 05	528 01	2872 01
42	305 00	8478 04	327 01	7056 01	30	105-01	307-02	430 01	2347 02	18	261-02	1025 05	557 01	2852 01
40	300 00	8601 04	340 01	8144 01	28	846-02	233-02	449 01	2871 02	16	194-02	1045 05	591 01	2832 01
38	295 00	8722 04	353 01	9469 01	26	675-02	172-02	470 01	3557 02	14	140-02	1064 05	631 01	2812 01
36	290 00	8840 04	367 01	1110 02	24	531-02	123-02	494 01	4472 02	12	978-03	1083 05	678 01	2792 01
34	285 00	8957 04	383 01	1312 02	22	412-02	885-03	520 01	5717 02	10	651-03	1102 05	736 01	2772 01
32	280 00	9070 04	399 01	1567 02	20	313-02	613-03	549 01	7448 02	08	406-03	1121 05	811 01	2752 01
30	275 00	9182 04	417 01	1890 02	18	233-02	403-03	583 01	9927 02	06	310-03	1140 05	857 01	2732 01
28	270 00	9291 04	436 01	2307 02	16	168-02	263-03	622 01	1360 03	04	229-03	1159 05	912 01	2712 01
26	265 00	9397 04	456 01	2851 02	14	117-02	161-03	668 01	1930 03	02	163-03	1178 05	978 01	2692 01
24	260 00	9502 04	479 01	3575 02	12	783-03	913-04	726 01	2870 03	00	110-03	1197 05	1066 05	2672 01
22	255 00	9603 04	505 01	4558 02	10	489-03	478-04	800 01	4558 03	00	695-04	1216 05	117 02	2652 01
20	250 00	9702 04	534 01	5922 02	08	374-03	320-04	845 01	5942 03	00	394-04	1235 05	131 02	2632 01
18	245 00	9799 04	567 01	7871 02	06	277-03	217-04	900 01	7989 03	00	190-04	1254 05	152 02	2612 01
16	240 00	9893 04	605 01	1075 03	04	197-03	133-05	965 01	1117 04	00	679-05	1273 05	187 02	2592 01
14	235 00	9985 04	650 01	1521 03	02	134-03	783-05	105 02	1643 04	00	117-05	1292 05	266 02	2572 01
12	230 00	1007 05	707 01	2254 03	00	842-04	411-05	105 05	2597 04	00	111-05	1311 05	1087 05	2552 01
10	225 00	1016 05	779 01	3565 03	00	478-04	187-05	129 02	4551 04	00	103 05	1330 05	1087 05	2532 01
08	220 00	1026 05	824 01	4639 03	00	231-04	677-06	150 02	9388 04	00	103 05	1349 05	1087 05	2512 01
06	215 00	1035 05	877 01	6223 03	00	128-05	162-06	184 02	1601 05	00	103 05	1368 05	1087 05	2492 01
04	210 00	1044 05	941 01	8675 03	00	143-05	142-07	260 02	1502 06	00	103 05	1387 05	1087 05	2472 01
02	205 00	1053 05	102 02	1273 04	00	01	01	01	01	00	01	01	01	01
00	200 00	1062 05	112 02	2005 04	00	01	01	01	01	00	01	01	01	01
	195 00	1071 05	126 02	3500 04	00	01	01	01	01	00	01	01	01	01
	190 00	1080 05	146 02	7184 04	00	01	01	01	01	00	01	01	01	01
	185 00	1089 05	160 02	1983 05	00	01	01	01	01	00	01	01	01	01
	180 00	1098 05	180 02	1983 05	00	01	01	01	01	00	01	01	01	01
	175 00	1107 05	252 02	1128 06	00	01	01	01	01	00	01	01	01	01

Table 17a. EXPANSION AFTER FREEZING - WEDGE NOZZLE

 $T_0' = 4000^\circ K, p_0' = 100 \text{ atm.}$

$l = 1.0$

$l = 0.1$

$l = 0.1$

$l = 1.0$

T	P	P	u'	M	A	T	P	P	u'	M	A	T	P	P	u'	M	A
.88	.502 00	.437 00	.4517 04	.122 01	.1020 01	.723	.227 00	.167 00	.6304 04	.186 01	.1617 01	.665	.127 00	.827-01	.7202 04	.225 01	.2520 01
.86	.484 00	.394 00	.4754 04	.129 01	.1049 01	.74	.214 00	.155 00	.6415 04	.188 01	.1687 01	.66	.124 00	.800-01	.7239 04	.225 01	.2572 01
.84	.468 00	.356 00	.4980 04	.136 01	.1085 01	.76	.195 00	.137 00	.6583 04	.196 01	.1805 01	.64	.112 00	.698-01	.7368 04	.233 01	.2800 01
.82	.452 00	.320 00	.5195 04	.144 01	.1129 01	.78	.177 00	.121 00	.6766 04	.203 01	.1940 01	.62	.100 00	.607-01	.7533 04	.242 01	.3080 01
.80	.436 00	.287 00	.5402 04	.152 01	.1181 01	.80	.160 00	.107 00	.6960 04	.211 01	.2092 01	.60	.898-01	.525-01	.7675 04	.250 01	.3358 01
.78	.420 00	.257 00	.5601 04	.159 01	.1241 01	.82	.145 00	.935-01	.7001 04	.219 01	.2265 01	.58	.800-01	.453-01	.7814 04	.259 01	.3700 01
.76	.404 00	.229 00	.5793 04	.167 01	.1310 01	.84	.130 00	.816-01	.7213 04	.227 01	.2462 01	.56	.711-01	.388-01	.7951 04	.268 01	.4055 01
.74	.388 00	.204 00	.5979 04	.174 01	.1390 01	.86	.117 00	.710-01	.7382 04	.236 01	.2687 01	.54	.620-01	.331-01	.8085 04	.278 01	.4553 01
.72	.372 00	.181 00	.6159 04	.182 01	.1481 01	.88	.105 00	.615-01	.7508 04	.244 01	.2944 01	.52	.554-01	.281-01	.8216 04	.287 01	.5086 01
.70	.356 00	.160 00	.6333 04	.190 01	.1584 01	.90	.934-01	.530-01	.7650 04	.253 01	.3240 01	.50	.485-01	.237-01	.8345 04	.298 01	.5711 01
.68	.340 00	.141 00	.6503 04	.198 01	.1701 01	.92	.829-01	.455-01	.7790 04	.262 01	.3581 01	.48	.424-01	.198-01	.8472 04	.308 01	.6448 01
.66	.324 00	.123 00	.6668 04	.206 01	.1834 01	.94	.734-01	.388-01	.7926 04	.272 01	.3977 01	.46	.368-01	.165-01	.8597 04	.319 01	.7320 01
.64	.308 00	.108 00	.6829 04	.214 01	.1987 01	.96	.647-01	.329-01	.8061 04	.281 01	.4438 01	.44	.317-01	.136-01	.8719 04	.331 01	.8301 01
.62	.292 00	.939-01	.6986 04	.222 01	.2160 01	.98	.567-01	.278-01	.8192 04	.292 01	.4977 01	.42	.272-01	.112-01	.8839 04	.343 01	.9612 01
.60	.276 00	.914-01	.7140 04	.231 01	.2359 01	.48	.495-01	.233-01	.8321 04	.302 01	.5613 01	.40	.232-01	.906-02	.8957 04	.356 01	.1113 02
.58	.260 00	.703-01	.7290 04	.240 01	.2588 01	.46	.430-01	.194-01	.8448 04	.313 01	.6365 01	.38	.197-01	.728-02	.9073 04	.370 01	.1298 02
.56	.244 00	.604-01	.7436 04	.249 01	.2852 01	.44	.372-01	.160-01	.8573 04	.325 01	.7262 01	.36	.165-01	.580-02	.9186 04	.384 01	.1526 02
.54	.228 00	.516-01	.7580 04	.258 01	.3158 01	.42	.319-01	.131-01	.8695 04	.337 01	.8340 01	.34	.138-01	.456-02	.9298 04	.400 01	.1810 02
.52	.212 00	.438-01	.7720 04	.268 01	.3514 01	.40	.272-01	.107-01	.8815 04	.350 01	.9643 01	.32	.113-01	.354-02	.9407 04	.417 01	.2167 02
.50	.196 00	.370-01	.7857 04	.278 01	.3929 01	.38	.231-01	.858-02	.8932 04	.363 01	.1123 02	.30	.927-02	.271-02	.9514 04	.435 01	.2623 02
.48	.180 00	.311-01	.7992 04	.288 01	.4419 01	.36	.194-01	.683-02	.9048 04	.378 01	.1319 02	.28	.749-02	.205-02	.9619 04	.454 01	.3211 02
.46	.164 00	.259-01	.8124 04	.299 01	.4997 01	.34	.162-01	.538-02	.9161 04	.393 01	.1563 02	.26	.598-02	.152-02	.9722 04	.475 01	.3982 02
.44	.148 00	.214-01	.8254 04	.311 01	.5686 01	.32	.133-01	.418-02	.9272 04	.410 01	.1870 02	.24	.470-02	.110-02	.9823 04	.499 01	.5010 02
.42	.132 00	.176-01	.8381 04	.323 01	.6511 01	.30	.109-01	.320-02	.9381 04	.427 01	.2260 02	.22	.364-02	.781-03	.9921 04	.525 01	.6409 02
.40	.116 00	.143-01	.8505 04	.335 01	.7509 01	.28	.882-02	.242-02	.9488 04	.447 01	.2764 02	.20	.276-02	.539-03	.1002 05	.555 01	.8358 02
.38	.100 00	.115-01	.8627 04	.349 01	.8724 01	.26	.703-02	.179-02	.9592 04	.468 01	.3424 02	.18	.203-02	.360-03	.1011 05	.588 01	.1115 03
.36	.084 00	.921-02	.8747 04	.363 01	.1022 02	.24	.555-02	.130-02	.9694 04	.491 01	.4303 02	.16	.148-02	.231-03	.1020 05	.628 01	.1529 03
.34	.068 00	.727-02	.8865 04	.378 01	.1207 02	.22	.430-02	.925-03	.9793 04	.517 01	.5498 02	.14	.104-02	.141-03	.1029 05	.675 01	.2172 03
.32	.052 00	.566-02	.8979 04	.394 01	.1440 02	.20	.327-02	.640-03	.9890 04	.546 01	.7160 02	.12	.690-03	.807-04	.1038 05	.733 01	.3233 03
.30	.036 00	.435-02	.9092 04	.412 01	.1736 02	.18	.243-02	.428-03	.9985 04	.580 01	.9338 02	.10	.430-03	.420-04	.1046 05	.807 01	.5141 03
.28	.020 00	.329-02	.9202 04	.431 01	.2116 02	.16	.176-02	.275-03	.1008 05	.619 01	.1306 03	.09	.328-03	.288-04	.1050 05	.853 01	.6707 03
.26	.004 00	.244-02	.9310 04	.451 01	.2614 02	.14	.123-02	.168-03	.1017 05	.665 01	.1853 03	.08	.243-03	.190-04	.1054 05	.908 01	.8024 03
.24	.980-02	.178-02	.9415 04	.474 01	.3275 02	.12	.819-03	.962-04	.1025 05	.723 01	.2754 03	.07	.173-03	.118-04	.1058 05	.974 01	.1262 04
.22	.583-02	.127-02	.9518 04	.500 01	.4171 02	.10	.512-03	.501-04	.1034 05	.796 01	.4371 03	.06	.117-03	.685-05	.1063 05	.106 02	.1859 04
.20	.444-02	.878-03	.9618 04	.528 01	.5414 02	.09	.391-03	.345-04	.1038 05	.842 01	.5697 03	.05	.738-04	.360-05	.1067 05	.116 02	.2941 04
.18	.331-02	.589-03	.9715 04	.561 01	.7189 02	.08	.290-03	.227-04	.1047 05	.896 01	.7656 03	.04	.419-04	.163-05	.1071 05	.130 02	.5159 04
.16	.240-02	.380-03	.9810 04	.599 01	.9810 02	.07	.207-03	.142-04	.1047 05	.961 01	.1070 04	.03	.202-04	.592-06	.1075 05	.151 02	.1066 05
.14	.169-02	.233-03	.9903 04	.644 01	.1386 03	.06	.140-03	.822-05	.1051 05	.104 02	.1574 04	.02	.722-05	.139-06	.1079 05	.186 02	.2970 05
.12	.113-02	.134-03	.9993 04	.700 01	.2052 03	.05	.882-04	.432-05	.1055 05	.115 02	.2486 04	.01	.125-05	.113-07	.1083 05	.270 02	.1717 06
.10	.708-03	.700-04	.1008 05	.772 01	.3242 03	.04	.502-04	.197-05	.1059 05	.129 02	.4353 04						
.09	.542-03	.482-04	.1012 05	.816 01	.4216 03	.03	.243-04	.711-06	.1063 05	.149 02	.8974 04						
.08	.403-03	.318-04	.1017 05	.868 01	.5651 03	.02	.870-05	.171-06	.1067 05	.183 02	.2492 05						
.07	.288-03	.199-04	.1021 05	.932 01	.7872 03	.01	.151-05	.142-07	.1071 05	.263 02	.1432 06						
.06	.195-03	.116-04	.1025 05	.101 02	.1154 04												
.05	.124-03	.611-05	.1030 05	.111 02	.1816 04												
.04	.707-04	.279-05	.1034 05	.125 02	.3166 04												
.03	.343-04	.102-05	.1038 05	.145 02	.6489 04												
.02	.124-04	.102-06	.1042 05	.178 02	.1787 05												
.01	.218-05	.227-07	.1046 05	.249 02	.1013 06												

Table 17b. EXPANSION AFTER FREEZING - HYPERBOLIC NOZZLE

$T_0' = 4000^\circ K$, $P_0' = 100$ atm.

$l = 10$ $l = 1.0$ $l = 0.1$

T	P	P	u' (ft/sec.)	M	A	T	P	P	u' (ft/sec.)	M	A	T	P	P	u' (ft/sec.)	M	A	T	P	P	u' (ft/sec.)	M	A
.77	.282 00	.214 00	.5865 04	.170 01	.1396 01	.651	.137 00	.978-01	.7067 04	.223 01	.2384 01	.558	.740-01	.406-01	.7826 04	.265 01	.3084 01						
.76	.270 00	.202 00	.5958 04	.172 01	.1437 01	.64	.129 00	.915-01	.7151 04	.225 01	.2498 01	.54	.664-01	.353-01	.7047 04	.277 01	.4372 01						
.74	.246 00	.180 00	.6138 04	.179 01	.1529 01	.62	.116 00	.908-01	.7301 04	.223 01	.2726 01	.52	.585-01	.299-01	.6831 04	.282 01	.4884 01						
.72	.224 00	.159 00	.6313 04	.187 01	.1633 01	.60	.104 00	.913-01	.7448 04	.242 01	.2989 01	.50	.513-01	.252-01	.6612 04	.292 01	.5482 01						
.70	.203 00	.141 00	.6484 04	.195 01	.1751 01	.58	.924-01	.923-01	.7591 04	.251 01	.3290 01	.48	.447-01	.211-01	.6341 04	.302 01	.5884 01						
.68	.184 00	.124 00	.6649 04	.203 01	.1886 01	.56	.821-01	.923-01	.7732 04	.260 01	.3638 01	.46	.388-01	.176-01	.6167 04	.313 01	.6184 01						
.66	.166 00	.108 00	.6811 04	.211 01	.2039 01	.54	.725-01	.923-01	.7869 04	.269 01	.4042 01	.44	.335-01	.145-01	.5991 04	.325 01	.6424 01						
.64	.150 00	.946-01	.6968 04	.219 01	.2213 01	.52	.635-01	.923-01	.8004 04	.279 01	.4513 01	.42	.287-01	.119-01	.5813 04	.337 01	.6724 01						
.62	.134 00	.823-01	.7122 04	.227 01	.2413 01	.50	.569-01	.923-01	.8137 04	.289 01	.5064 01	.40	.245-01	.93-02	.5632 04	.350 01	.6968 01						
.60	.120 00	.712-01	.7272 04	.236 01	.2641 01	.48	.489-01	.923-01	.8267 04	.299 01	.5513 01	.38	.207-01	.714-02	.5390 04	.363 01	.7145 01						
.58	.107 00	.614-01	.7419 04	.244 01	.2904 01	.46	.424-01	.923-01	.8394 04	.310 01	.6483 01	.36	.174-01	.616-02	.5065 04	.378 01	.7484 01						
.56	.932-01	.526-01	.7563 04	.254 01	.3207 01	.44	.366-01	.923-01	.8519 04	.322 01	.7401 01	.34	.145-01	.485-02	.4918 04	.393 01	.7747 01						
.54	.842-01	.449-01	.7704 04	.263 01	.3559 01	.42	.314-01	.923-01	.8642 04	.334 01	.8505 01	.32	.119-01	.376-02	.4684 04	.410 01	.7980 01						
.52	.746-01	.381-01	.7842 04	.273 01	.3969 01	.40	.268-01	.923-01	.8763 04	.347 01	.9842 01	.30	.976-02	.288-02	.4397 04	.427 01	.8257 01						
.50	.651-01	.321-01	.7977 04	.283 01	.4449 01	.38	.227-01	.923-01	.8891 04	.360 01	.1147 02	.28	.788-02	.217-02	.4203 04	.447 01	.8582 01						
.48	.568-01	.269-01	.8109 04	.293 01	.5014 01	.36	.190-01	.923-01	.8997 04	.375 01	.1349 02	.26	.628-02	.161-02	.3907 04	.468 01	.8824 01						
.46	.483-01	.224-01	.8239 04	.304 01	.5684 01	.34	.158-01	.923-01	.9111 04	.390 01	.1599 02	.24	.494-02	.117-02	.3608 04	.491 01	.9012 01						
.44	.426-01	.185-01	.8367 04	.316 01	.6483 01	.32	.131-01	.923-01	.9222 04	.407 01	.1815 02	.22	.382-02	.821-03	.3208 04	.517 01	.9154 01						
.42	.365-01	.151-01	.8492 04	.328 01	.7442 01	.30	.107-01	.923-01	.9332 04	.424 01	.2317 02	.20	.290-02	.571-03	.2905 04	.546 01	.9303 01						
.40	.311-01	.123-01	.8615 04	.340 01	.8603 01	.28	.862-02	.923-01	.9439 04	.443 01	.2836 02	.18	.215-02	.381-03	.2699 04	.580 01	.9494 01						
.38	.264-01	.990-02	.8735 04	.354 01	.1002 02	.26	.688-02	.923-01	.9544 04	.465 01	.3316 02	.16	.156-02	.245-03	.2409 04	.618 01	.9671 01						
.36	.222-01	.788-02	.8852 04	.368 01	.1177 02	.24	.541-02	.923-01	.9666 04	.488 01	.4424 02	.14	.108-02	.149-03	.2118 04	.665 01	.9841 01						
.34	.185-01	.620-02	.8969 04	.383 01	.1394 02	.22	.419-02	.923-01	.9746 04	.513 01	.5660 02	.12	.722-03	.852-04	.1827 04	.722 01	.9994 01						
.32	.152-01	.482-02	.9082 04	.400 01	.1667 02	.20	.318-02	.923-01	.9843 04	.543 01	.7380 02	.10	.450-03	.443-04	.1335 04	.796 01	.9999 04						
.30	.125-01	.369-02	.9193 04	.417 01	.2016 02	.18	.236-02	.923-01	.9938 04	.576 01	.9846 02	.08	.343-03	.304-04	.1040 05	.841 01	.9999 04						
.28	.101-01	.278-02	.9302 04	.436 01	.2465 02	.16	.170-02	.923-01	.1003 05	.614 01	.1350 03	.06	.254-03	.200-04	.1044 05	.895 01	.9999 04						
.26	.84-02	.206-02	.9408 04	.457 01	.3054 02	.14	.119-02	.923-01	.1012 05	.661 01	.1919 03	.04	.181-03	.125-04	.1048 05	.960 01	.9999 04						
.24	.632-02	.150-02	.9512 04	.480 01	.3838 02	.12	.792-03	.923-01	.1021 05	.718 01	.2856 03	.02	.122-03	.722-05	.1042 05	.104 02	.9999 04						
.22	.490-02	.105-02	.9613 04	.506 01	.4905 02	.10	.494-03	.923-01	.1034 05	.791 01	.4342 03	.00	.769-04	.378-05	.1036 05	.114 02	.9999 04						
.20	.372-02	.735-03	.9712 04	.534 01	.6390 02	.08	.377-03	.923-01	.1034 05	.836 01	.5927 03	.00	.436-04	.172-05	.1060 05	.128 02	.9999 04						
.18	.276-02	.491-03	.9808 04	.567 01	.8515 02	.06	.279-03	.923-01	.1034 05	.890 01	.7976 03	.00	.210-04	.621-06	.1064 05	.149 02	.9999 04						
.16	.200-02	.316-03	.9902 04	.606 01	.1167 03	.04	.193-03	.923-01	.1034 05	.954 01	.1116 04	.00	.750-05	.147-06	.1068 05	.183 02	.9999 04						
.14	.143-02	.193-03	.9993 04	.651 01	.1656 03	.02	.134-03	.923-01	.1034 05	.1046 05	.1644 04	.00	.125-05	.115-07	.1072 05	.268 02	.9999 04						
.12	.930-03	.110-03	.1008 05	.708 01	.2462 03	.00	.843-04	.923-01	.1034 05	.1050 05	.2502 04	.00											
.10	.561-03	.573-04	.1017 05	.780 01	.3910 03	.04	.479-04	.923-01	.1034 05	.1054 05	.4566 04	.00											
.08	.443-03	.394-04	.1021 05	.824 01	.5098 03	.03	.231-04	.923-01	.1034 05	.1059 05	.9440 04	.00											
.06	.326-03	.259-04	.1026 05	.878 01	.6854 03	.02	.825-05	.923-01	.1034 05	.1063 05	.2632 05	.00											
.04	.234-03	.162-04	.1030 05	.941 01	.9580 03	.01	.142-05	.923-01	.1034 05	.1067 05	.1324 06	.00											
.02	.158-03	.938-05	.1034 05	.102 02	.1410 04	.00						.00											
.00	.998-04	.492-05	.1038 05	.112 02	.2229 04	.00						.00											
.00	.567-04	.224-05	.1042 05	.126 02	.3907 04	.00						.00											
.00	.274-04	.811-06	.1046 05	.146 02	.8064 04	.00						.00											
.00	.980-05	.193-06	.1051 05	.180 02	.2243 05	.00						.00											
.00	.169-05	.172-07	.1055 05	.253 02	.1294 06	.00						.00											

Table 18a. EXPANSION AFTER FREEZING - WEDGE NOZZLE

 $T_0' = 4000^\circ\text{K}$, $P_0' = 300 \text{ atm}$.

$\ell = 10$
 $\ell = 1.0$
 $\ell = 0.1$

T	ρ	P	u' (ft/sec.)	M	A	T	ρ	P	u' (ft/sec.)	M	A	T	ρ	P	u' (ft/sec.)	M	A	T	ρ	P	u' (ft/sec.)	M	A	T
.776	.293 00	.225 00	.5790 04	.167 01	.1361 01	.684	.167 0	.113 00	.6766 04	.208 01	.2037 01	.597	.960-01	.504-01	.7523 04	.247 01	.2197 01							
.76	.273 00	.205 00	.5939 04	.171 01	.1425 01	.68	.164 04	.110 00	.6798 04	.207 01	.2069 01	.58	.870-01	.497-01	.7643 04	.252 01	.3171 01							
.74	.249 00	.182 00	.6120 04	.179 01	.1516 01	.66	.143 00	.965-01	.6956 04	.215 01	.2239 01	.56	.773-01	.426-01	.7783 04	.261 01	.3810 01							
.72	.227 00	.161 00	.6296 04	.186 01	.1619 01	.64	.138 00	.841-01	.7110 04	.223 01	.2434 01	.54	.683-01	.363-01	.7920 04	.271 01	.4267 01							
.70	.206 00	.142 00	.6467 04	.194 01	.1736 01	.62	.120 00	.731-01	.7261 04	.232 01	.2656 01	.52	.602-01	.308-01	.8054 04	.281 01	.4764 01							
.68	.186 00	.125 00	.6633 04	.202 01	.1868 01	.60	.107 00	.633-01	.7408 04	.240 01	.2910 01	.50	.527-01	.259-01	.8185 04	.291 01	.4344 01							
.66	.168 00	.105 00	.6794 04	.210 01	.2020 01	.58	.954-01	.545-01	.7552 04	.249 01	.3203 01	.48	.460-01	.181-01	.8315 04	.301 01	.4034 01							
.64	.151 00	.938-01	.6952 04	.218 01	.2192 01	.56	.847-01	.468-01	.7694 04	.258 01	.3541 01	.46	.399-01	.131-01	.8441 04	.312 01	.4851 01							
.62	.136 00	.833-01	.7106 04	.227 01	.2389 01	.54	.749-01	.399-01	.7832 04	.268 01	.3933 01	.44	.344-01	.104-01	.8566 04	.324 01	.4783 01							
.60	.122 00	.721-01	.7257 04	.235 01	.2615 01	.52	.660-01	.338-01	.7968 04	.277 01	.4390 01	.42	.295-01	.91-02	.8668 04	.336 01	.4943 01							
.58	.108 00	.611-01	.7404 04	.244 01	.2875 01	.50	.579-01	.285-01	.8101 04	.287 01	.4925 01	.40	.252-01	.797-02	.8808 04	.349 01	.5141 01							
.56	.963-01	.533-01	.7548 04	.253 01	.3175 01	.48	.505-01	.230-01	.8231 04	.298 01	.5555 01	.38	.213-01	.797-02	.8926 04	.362 01	.5171 01							
.54	.852-01	.455-01	.7689 04	.262 01	.3522 01	.46	.438-01	.199-01	.8359 04	.309 01	.6302 01	.36	.179-01	.634-02	.9041 04	.377 01	.5171 01							
.52	.751-01	.366-01	.7828 04	.272 01	.3927 01	.44	.378-01	.164-01	.8485 04	.320 01	.7193 01	.34	.149-01	.499-02	.9154 04	.392 01	.5171 01							
.50	.659-01	.325-01	.7953 04	.282 01	.4402 01	.42	.325-01	.134-01	.8608 04	.333 01	.8264 01	.32	.123-01	.387-02	.9265 04	.404 01	.5171 01							
.48	.575-01	.273-01	.8096 04	.293 01	.4960 01	.40	.277-01	.109-01	.8729 04	.345 01	.9561 01	.30	.100-01	.286-02	.9374 04	.426 01	.5171 01							
.46	.495-01	.227-01	.8226 04	.304 01	.5622 01	.38	.234-01	.877-02	.8848 04	.359 01	.1114 02	.28	.811-02	.223-02	.9481 04	.446 01	.5171 01							
.44	.413-01	.187-01	.8354 04	.315 01	.6412 01	.36	.197-01	.697-02	.8964 04	.373 01	.1310 02	.26	.647-02	.165-02	.9585 04	.467 01	.5171 01							
.42	.330-01	.153-01	.8479 04	.327 01	.7360 01	.34	.164-01	.549-02	.9079 04	.389 01	.1553 02	.24	.508-02	.120-02	.9687 04	.490 01	.5171 01							
.40	.247-01	.125-01	.8602 04	.340 01	.8507 01	.32	.135-01	.428-02	.9191 04	.405 01	.1858 02	.22	.393-02	.852-03	.9786 04	.516 01	.5171 01							
.38	.267-01	.100-01	.8722 04	.353 01	.9907 01	.30	.110-01	.326-02	.9300 04	.423 01	.2248 02	.20	.299-02	.588-03	.9882 04	.545 01	.5171 01							
.36	.224-01	.798-02	.8841 04	.367 01	.1163 02	.28	.892-02	.246-02	.9408 04	.442 01	.2751 02	.18	.222-02	.393-03	.9978 04	.578 01	.5171 01							
.34	.187-01	.658-02	.8956 04	.383 01	.1378 02	.26	.711-02	.183-02	.9513 04	.463 01	.3411 02	.16	.160-02	.282-03	.1007 05	.617 01	.5171 01							
.32	.154-01	.488-02	.9070 04	.399 01	.1648 02	.24	.560-02	.132-02	.9616 04	.486 01	.4290 02	.14	.112-02	.184-03	.1016 05	.663 01	.5171 01							
.30	.126-01	.374-02	.9181 04	.417 01	.1992 02	.22	.429-02	.933-03	.9716 04	.512 01	.5487 02	.12	.743-03	.878-04	.1025 05	.721 01	.5171 01							
.28	.102-01	.282-02	.9290 04	.436 01	.2435 02	.20	.329-02	.648-03	.9814 04	.541 01	.7154 02	.10	.463-03	.456-04	.1033 05	.794 01	.5171 01							
.26	.814-02	.209-02	.9397 04	.456 01	.3016 02	.18	.244-02	.433-03	.9909 04	.574 01	.9541 02	.09	.354-03	.313-04	.1037 05	.839 01	.5171 01							
.24	.641-02	.152-02	.9501 04	.479 01	.3791 02	.16	.176-02	.278-03	.1000 05	.612 01	.1308 03	.08	.262-03	.206-04	.1042 05	.893 01	.5171 01							
.22	.496-02	.108-02	.9602 04	.505 01	.4844 02	.14	.123-02	.170-03	.1009 05	.659 01	.1858 03	.07	.186-03	.128-04	.1046 05	.958 01	.5171 01							
.20	.377-02	.745-03	.9701 04	.534 01	.6309 02	.12	.820-03	.969-04	.1018 05	.715 01	.2766 03	.06	.126-03	.743-05	.1050 05	.104 02	.5171 01							
.18	.280-02	.498-03	.9797 04	.566 01	.8406 02	.10	.511-03	.504-04	.1027 05	.788 01	.4397 03	.05	.792-04	.390-05	.1054 05	.114 02	.5171 01							
.16	.203-02	.320-03	.9891 04	.605 01	.1152 03	.09	.390-03	.346-04	.1031 05	.833 01	.5736 03	.04	.450-04	.177-05	.1058 05	.128 02	.5171 01							
.14	.142-02	.196-03	.9983 04	.650 01	.1634 03	.08	.289-03	.228-04	.1035 05	.887 01	.7717 03	.03	.217-04	.638-06	.1062 05	.149 02	.5171 01							
.12	.944-03	.112-03	.1007 05	.707 01	.2429 03	.07	.206-03	.142-04	.1039 05	.931 01	.1079 04	.02	.773-05	.152-06	.1066 05	.183 02	.5171 01							
.10	.586-03	.582-04	.1016 05	.779 01	.3856 03	.06	.139-03	.823-05	.1043 05	.103 02	.1590 04	.01	.133-05	.143-07	.1070 05	.253 02	.5171 01							
.09	.450-03	.400-04	.1020 05	.823 01	.5027 03	.05	.876-04	.432-05	.1048 05	.113 02	.2515 04													
.08	.333-03	.264-04	.1024 05	.876 01	.6758 03	.04	.497-04	.196-05	.1052 05	.127 02	.4413 04													
.07	.238-03	.164-04	.1029 05	.940 01	.9444 03	.03	.240-04	.708-06	.1056 05	.168 02	.9120 04													
.06	.161-03	.993-05	.1033 05	.102 02	.1390 04	.02	.857-05	.170-06	.1060 05	.181 02	.2541 05													
.05	.101-03	.500-05	.1037 05	.112 02	.2197 04	.01	.148-05	.143-07	.1064 05	.259 02	.1470 06													
.04	.576-04	.228-05	.1041 05	.126 02	.3849 04																			
.03	.278-04	.825-06	.1045 05	.146 02	.7941 04																			
.02	.996-05	.196-06	.1050 05	.180 02	.2208 05																			
.01	.172-05	.172-07	.1054 05	.254 02	.1272 06																			

Table 18b. EXPANSION AFTER FREEZING - HYPERBOLIC NOZZLE
 $T_0' = 4000^\circ\text{K}$, $p_0' = 300 \text{ atm}$.

$l = 0.1$															$l = 1.0$															$l = 10$														
T	P	P	u'	M	A	T	P	P	u'	M	A	T	P	P	u'	M	A	T	P	P	u'	M	A	T	P	P	u'	M	A															
			(ft./sec.)						(ft./sec.)						(ft./sec.)						(ft./sec.)																							
.668	.168 00	.112 00	.6737 04	.209 01	.2032 01	.57	.895-01	.506-01	.7563 04	.253 01	.3407 01	.495	.533-01	.261-01	.8115 04	.290 01	.3407 01	.495	.533-01	.261-01	.8115 04	.290 01	.3407 01	.495	.533-01	.261-01	.8115 04	.290 01	.3407 01															
.66	.162 00	.106 00	.6801 04	.210 01	.2097 01	.56	.843-01	.468-01	.7633 04	.256 01	.3585 01	.48	.481-01	.220-01	.8315 04	.307 01	.3585 01	.48	.481-01	.220-01	.8315 04	.307 01	.3585 01	.48	.481-01	.220-01	.8315 04	.307 01	.3585 01															
.64	.146 00	.923-01	.6959 04	.218 01	.2277 01	.54	.743-01	.399-01	.7773 04	.265 01	.3891 01	.46	.417-01	.160-01	.8335 04	.319 01	.3891 01	.46	.417-01	.160-01	.8335 04	.319 01	.3891 01	.46	.417-01	.160-01	.8335 04	.319 01	.3891 01															
.62	.131 00	.803-01	.7113 04	.226 01	.2483 01	.52	.650-01	.338-01	.7909 04	.275 01	.4444 01	.34	.360-01	.157-01	.8445 04	.337 01	.4444 01	.34	.360-01	.157-01	.8445 04	.337 01	.4444 01	.34	.360-01	.157-01	.8445 04	.337 01	.4444 01															
.60	.117 00	.595-01	.7263 04	.235 01	.2719 01	.50	.575-01	.285-01	.8143 04	.285 01	.4584 01	.32	.309-01	.125-01	.8538 04	.344 01	.4584 01	.32	.309-01	.125-01	.8538 04	.344 01	.4584 01	.32	.309-01	.125-01	.8538 04	.344 01	.4584 01															
.58	.104 00	.598-01	.7410 04	.244 01	.2991 01	.48	.502-01	.239-01	.8304 04	.295 01	.4622 01	.30	.263-01	.103-01	.8709 04	.344 01	.4622 01	.30	.263-01	.103-01	.8709 04	.344 01	.4622 01	.30	.263-01	.103-01	.8709 04	.344 01	.4622 01															
.56	.924-01	.513-01	.7554 04	.253 01	.3305 01	.46	.435-01	.199-01	.8430 04	.306 01	.4678 01	.28	.187-01	.067-02	.8928 04	.371 01	.4678 01	.28	.187-01	.067-02	.8928 04	.371 01	.4678 01	.28	.187-01	.067-02	.8928 04	.371 01	.4678 01															
.54	.817-01	.437-01	.7695 04	.262 01	.3669 01	.44	.376-01	.164-01	.8530 04	.318 01	.4780 01	.26	.156-01	.055-02	.9050 04	.387 01	.4780 01	.26	.156-01	.055-02	.9050 04	.387 01	.4780 01	.26	.156-01	.055-02	.9050 04	.387 01	.4780 01															
.52	.719-01	.371-01	.7833 04	.272 01	.4093 01	.42	.322-01	.134-01	.8634 04	.330 01	.4865 01	.24	.128-01	.037-02	.9171 04	.404 01	.4865 01	.24	.128-01	.037-02	.9171 04	.404 01	.4865 01	.24	.128-01	.037-02	.9171 04	.404 01	.4865 01															
.50	.631-01	.313-01	.7968 04	.282 01	.4590 01	.40	.275-01	.109-01	.8736 04	.342 01	.4975 01	.22	.105-01	.025-02	.9281 04	.421 01	.4975 01	.22	.105-01	.025-02	.9281 04	.421 01	.4975 01	.22	.105-01	.025-02	.9281 04	.421 01	.4975 01															
.48	.550-01	.262-01	.8101 04	.292 01	.5175 01	.38	.233-01	.075-02	.8795 04	.356 01	.5128 02	.20	.084-02	.013-03	.9386 04	.440 01	.5128 02	.20	.084-02	.013-03	.9386 04	.440 01	.5128 02	.20	.084-02	.013-03	.9386 04	.440 01	.5128 02															
.46	.477-01	.218-01	.8231 04	.303 01	.5869 01	.36	.195-01	.066-02	.8913 04	.370 01	.5326 02	.18	.067-02	.009-04	.9494 04	.461 01	.5326 02	.18	.067-02	.009-04	.9494 04	.461 01	.5326 02	.18	.067-02	.009-04	.9494 04	.461 01	.5326 02															
.44	.412-01	.189-01	.8358 04	.315 01	.6697 01	.34	.163-01	.048-02	.9027 04	.385 01	.5722 02	.16	.051-02	.006-06	.9681 04	.482 01	.5722 02	.16	.051-02	.006-06	.9681 04	.482 01	.5722 02	.16	.051-02	.006-06	.9681 04	.482 01	.5722 02															
.42	.353-01	.147-01	.8483 04	.327 01	.7692 01	.32	.134-01	.025-02	.9140 04	.402 01	.5882 02	.14	.031-02	.002-08	.9870 04	.508 01	.5882 02	.14	.031-02	.002-08	.9870 04	.508 01	.5882 02	.14	.031-02	.002-08	.9870 04	.508 01	.5882 02															
.40	.301-01	.119-01	.8606 04	.339 01	.8697 01	.30	.110-01	.026-02	.9230 04	.419 01	.6277 02	.12	.023-02	.001-10	.9994 04	.537 01	.6277 02	.12	.023-02	.001-10	.9994 04	.537 01	.6277 02	.12	.023-02	.001-10	.9994 04	.537 01	.6277 02															
.38	.255-01	.090-02	.8727 04	.353 01	.1037 02	.28	.088-02	.013-03	.9358 04	.438 01	.6566 02	.10	.017-02	.000-12	.9994 04	.570 01	.6566 02	.10	.017-02	.000-12	.9994 04	.570 01	.6566 02	.10	.017-02	.000-12	.9994 04	.570 01	.6566 02															
.36	.214-01	.044-02	.8845 04	.367 01	.1218 02	.26	.070-02	.012-02	.9464 04	.459 01	.6856 02	.08	.016-02	.000-14	.9994 04	.608 01	.6856 02	.08	.016-02	.000-14	.9994 04	.608 01	.6856 02	.08	.016-02	.000-14	.9994 04	.608 01	.6856 02															
.34	.178-01	.031-02	.8960 04	.382 01	.1444 02	.24	.054-02	.012-02	.9568 04	.482 01	.7152 02	.06	.012-02	.000-16	.9994 04	.654 01	.7152 02	.06	.012-02	.000-16	.9994 04	.654 01	.7152 02	.06	.012-02	.000-16	.9994 04	.654 01	.7152 02															
.32	.147-01	.027-02	.9074 04	.399 01	.1728 02	.22	.042-02	.013-03	.9668 04	.508 01	.7453 02	.04	.010-03	.000-18	.9994 04	.710 01	.7453 02	.04	.010-03	.000-18	.9994 04	.710 01	.7453 02	.04	.010-03	.000-18	.9994 04	.710 01	.7453 02															
.30	.120-01	.017-02	.9185 04	.416 01	.2090 02	.20	.032-02	.013-03	.9766 04	.537 01	.7753 02	.02	.009-03	.000-20	.9994 04	.750 01	.7753 02	.02	.009-03	.000-20	.9994 04	.750 01	.7753 02	.02	.009-03	.000-20	.9994 04	.750 01	.7753 02															
.28	.973-02	.269-02	.9294 04	.435 01	.2558 02	.18	.024-02	.013-03	.9862 04	.570 01	.8047 02	.00	.008-03	.000-22	.9994 04	.808 01	.8047 02	.00	.008-03	.000-22	.9994 04	.808 01	.8047 02	.00	.008-03	.000-22	.9994 04	.808 01	.8047 02															
.26	.775-02	.193-02	.9400 04	.456 01	.3171 02	.16	.017-02	.013-03	.9958 04	.608 01	.8349 02	.00	.007-03	.000-24	.9994 04	.854 01	.8349 02	.00	.007-03	.000-24	.9994 04	.854 01	.8349 02	.00	.007-03	.000-24	.9994 04	.854 01	.8349 02															
.24	.653-02	.145-02	.9504 04	.479 01	.3989 02	.14	.012-02	.013-03	.9958 04	.654 01	.8640 02	.00	.006-03	.000-26	.9994 04	.908 01	.8640 02	.00	.006-03	.000-26	.9994 04	.908 01	.8640 02	.00	.006-03	.000-26	.9994 04	.908 01	.8640 02															
.22	.471-02	.103-02	.9604 04	.504 01	.5102 02	.12	.010-03	.013-03	.9958 04	.710 01	.8937 02	.00	.005-03	.000-28	.9994 04	.954 01	.8937 02	.00	.005-03	.000-28	.9994 04	.954 01	.8937 02	.00	.005-03	.000-28	.9994 04	.954 01	.8937 02															
.20	.357-02	.078-03	.9704 04	.533 01	.5652 02	.10	.008-03	.013-03	.9958 04	.750 01	.9237 02	.00	.004-03	.000-30	.9994 04	.1005 05	.9237 02	.00	.004-03	.000-30	.9994 04	.1005 05	.9237 02	.00	.004-03	.000-30	.9994 04	.1005 05	.9237 02															
.18	.265-02	.043-03	.9800 04	.566 01	.6872 02	.08	.005-03	.013-03	.9958 04	.808 01	.9529 03	.00	.003-03	.000-32	.9994 04	.1026 05	.9529 03	.00	.003-03	.000-32	.9994 04	.1026 05	.9529 03	.00	.003-03	.000-32	.9994 04	.1026 05	.9529 03															
.16	.192-02	.034-03	.9894 04	.604 01	.8217 03	.06	.003-03	.013-03	.9958 04	.854 01	.9844 03	.00	.002-03	.000-34	.9994 04	.1031 05	.9844 03	.00	.002-03	.000-34	.9994 04	.1031 05	.9844 03	.00	.002-03	.000-34	.9994 04	.1031 05	.9844 03															
.14	.134-02	.018-03	.9985 04	.650 01	.9278 03	.04	.002-03	.013-03	.9958 04	.908 01	.1005 05	.00	.001-03	.000-36	.9994 04	.1035 05	.1005 05	.00	.001-03	.000-36	.9994 04	.1035 05	.1005 05	.00	.001-03	.000-36	.9994 04	.1035 05	.1005 05															
.12	.890-03	.105-03	.1007 05	.706 01	.2573 03	.00	.000-03	.013-03	.9958 04	.1005 05	.1005 05	.00	.000-03	.000-38	.9994 04	.1035 05	.1005 05	.00	.000-03	.000-38	.9994 04	.1035 05	.1005 05	.00	.000-03	.000-38	.9994 04	.1035 05	.1005 05															
.10	.555-03	.050-04	.1016 05	.778 01	.4091 03	.00	.000-03	.013-03	.9958 04	.1005 05	.1005 05	.00	.000-03	.000-40	.9994 04	.1035 05	.1005 05	.00	.000-03	.000-40	.9994 04	.1035 05	.1005 05	.00	.000-03	.000-40	.9994 04	.1035 05	.1005 05															
.08	.423-03	.037-04	.1020 05	.822 01	.5338 03	.00	.000-03	.013-03	.9958 04	.1005 05	.1005 05	.00	.000-03	.000-42	.9994 04	.1035 05	.1005 05	.00	.000-03	.000-42	.9994 04	.1035 05	.1005 05	.00	.000-03	.000-42	.9994 04	.1035 05	.1005 05															
.06	.313-03	.024-04	.1025 05	.875 01	.6183 03	.00	.000-03	.013-03	.9958 04	.1005 05	.1005 05	.00	.000-03	.000-44	.9994 04	.1035 05	.1005 05	.00	.000-03	.000-44	.9994 04	.1035 05	.1005 05	.00	.000-03	.000-44	.9994 04	.1035 05	.1005 05															
.04	.223-03	.015-04	.1029 05	.939 01	.7005 04	.00	.000-03	.013-03	.9958 04	.1005 05	.1005 05	.00	.000-03	.000-46	.9994 04	.1035 05	.1005 05	.00	.000-03	.000-46	.9994 04	.1035 05	.1005 05	.00	.000-03	.																		

$l = 0.1$ $l = 1.0$ $l = 10$

T	P	P	u' (ft/sec.)	M	A	T	P	P	u' (ft/sec.)	M	A	T	P	P	u' (ft/sec.)	M	A
.697	.200 00	.138 00	.6463 04	.196 01	.1789 01	.609	.116 00	.700-01	.7251 04	.235 01	.2742 01	.536	.714-01	.379-01	.7819 04	.269 01	.4131 01
.68	.183 00	.124 00	.6604 04	.201 01	.1905 01	.60	.110 00	.656-01	.7318 04	.237 01	.2858 01	.52	.645-01	.339-01	.7828 04	.275 01	.4513 01
.66	.165 00	.108 00	.6767 04	.209 01	.2060 01	.58	.983-01	.565-01	.7464 04	.246 01	.3145 01	.50	.565-01	.280-01	.8062 04	.285 01	.5063 01
.64	.149 00	.945-01	.6925 04	.217 01	.2236 01	.56	.872-01	.484-01	.7606 04	.255 01	.3476 01	.48	.693-01	.234-01	.8193 04	.296 01	.5711 01
.62	.134 00	.822-01	.7080 04	.225 01	.2438 01	.54	.771-01	.413-01	.7746 04	.264 01	.3860 01	.46	.429-01	.193-01	.8321 04	.307 01	.6480 01
.60	.120 00	.711-01	.7231 04	.234 01	.2669 01	.52	.679-01	.350-01	.7884 04	.274 01	.4307 01	.44	.369-01	.161-01	.8448 04	.318 01	.7307 01
.58	.107 00	.613-01	.7378 04	.243 01	.2935 01	.50	.585-01	.295-01	.8018 04	.284 01	.4831 01	.42	.317-01	.131-01	.8571 04	.330 01	.8500 01
.56	.946-01	.525-01	.7523 04	.252 01	.3242 01	.48	.519-01	.247-01	.8150 04	.294 01	.5449 01	.40	.270-01	.107-01	.8693 04	.343 01	.8835 01
.54	.836-01	.448-01	.7664 04	.261 01	.3598 01	.46	.451-01	.205-01	.8279 04	.305 01	.6181 01	.38	.228-01	.855-02	.8812 04	.356 01	.1147 02
.52	.738-01	.380-01	.7803 04	.271 01	.4013 01	.44	.389-01	.170-01	.8406 04	.317 01	.7054 01	.36	.192-01	.538-02	.8929 04	.371 01	.1348 02
.50	.646-01	.320-01	.7939 04	.281 01	.4500 01	.42	.334-01	.139-01	.8530 04	.329 01	.8104 01	.34	.160-01	.417-02	.9044 04	.386 01	.1598 02
.48	.563-01	.268-01	.8072 04	.291 01	.5073 01	.40	.284-01	.113-01	.8652 04	.341 01	.9375 01	.32	.132-01	.321-02	.9156 04	.407 01	.1913 02
.46	.489-01	.223-01	.8202 04	.302 01	.5752 01	.38	.241-01	.906-02	.8772 04	.355 01	.1093 02	.30	.108-01	.241-02	.9266 04	.420 01	.2315 02
.44	.422-01	.184-01	.8330 04	.314 01	.6562 01	.36	.202-01	.721-02	.8890 04	.369 01	.1284 02	.28	.888-02	.241-02	.9374 04	.439 01	.2834 02
.42	.363-01	.151-01	.8456 04	.326 01	.7536 01	.34	.168-01	.567-02	.9005 04	.384 01	.1523 02	.26	.662-02	.178-02	.9479 04	.460 01	.3514 02
.40	.308-01	.122-01	.8579 04	.338 01	.8715 01	.32	.139-01	.440-02	.9118 04	.401 01	.1623 02	.24	.544-02	.129-02	.9582 04	.483 01	.4422 02
.38	.261-01	.984-02	.8700 04	.352 01	.1016 02	.30	.113-01	.337-02	.9228 04	.418 01	.2205 02	.22	.421-02	.918-03	.9683 04	.509 01	.5657 02
.36	.219-01	.783-02	.8818 04	.366 01	.1193 02	.28	.915-02	.254-02	.9336 04	.437 01	.2698 02	.20	.370-02	.631-03	.9781 04	.538 01	.7378 02
.34	.183-01	.616-02	.8934 04	.381 01	.1414 02	.26	.730-02	.188-02	.9442 04	.458 01	.3346 02	.18	.237-02	.423-03	.9876 04	.570 01	.9843 02
.32	.151-01	.478-02	.9048 04	.397 01	.1692 02	.24	.574-02	.137-02	.9546 04	.481 01	.4209 02	.16	.171-02	.272-03	.9969 04	.609 01	.1350 03
.30	.123-01	.366-02	.9160 04	.415 01	.2046 02	.22	.444-02	.968-03	.9647 04	.507 01	.5384 02	.14	.119-02	.166-03	.1006 05	.655 01	.1918 03
.28	.994-02	.276-02	.9269 04	.434 01	.2504 02	.20	.337-02	.668-03	.9745 04	.535 01	.7021 02	.12	.795-03	.945-04	.1015 05	.711 01	.2858 03
.26	.793-02	.204-02	.9375 04	.455 01	.3103 02	.18	.250-02	.446-03	.9841 04	.568 01	.9366 02	.10	.466-03	.491-04	.1023 05	.784 01	.4546 03
.24	.623-02	.148-02	.9479 04	.478 01	.3903 02	.16	.181-02	.287-03	.9934 04	.607 01	.1285 03	.08	.378-03	.331-04	.1028 05	.839 01	.5933 03
.22	.482-02	.105-02	.9581 04	.503 01	.6506 02	.14	.126-02	.175-03	.1003 05	.652 01	.1825 03	.06	.280-03	.222-04	.1032 05	.882 01	.7985 03
.20	.366-02	.726-03	.9680 04	.532 01	.6506 02	.12	.839-03	.998-04	.1011 05	.709 01	.2718 03	.04	.199-03	.138-04	.1036 05	.946 01	.1117 04
.18	.272-02	.485-03	.9777 04	.564 01	.8677 02	.10	.523-03	.518-04	.1024 05	.781 01	.4322 03	.02	.135-03	.800-05	.1040 05	.103 02	.1647 04
.16	.196-02	.312-03	.9871 04	.602 01	.1190 03	.08	.399-03	.356-04	.1029 05	.826 01	.5640 03	.00	.847-04	.420-05	.1044 05	.113 02	.2607 04
.14	.137-02	.190-03	.9962 04	.648 01	.1690 03	.06	.285-03	.234-04	.1033 05	.879 01	.7590 03	.04	.481-04	.190-05	.1049 05	.127 02	.4577 04
.12	.912-03	.109-03	.1005 05	.704 01	.2515 03	.04	.210-03	.146-04	.1037 05	.943 01	.1062 04	.02	.231-04	.686-06	.1053 05	.147 02	.9468 04
.10	.565-03	.564-04	.1014 05	.776 01	.3998 03	.02	.142-03	.845-05	.1041 05	.102 02	.1565 04	.00	.828-05	.163-06	.1057 05	.180 02	.2647 05
.08	.434-03	.388-04	.1018 05	.820 01	.5216 03	.00	.894-04	.443-05	.1045 05	.112 02	.2477 04	.01	.142-05	.144-07	.1061 05	.234 02	.1531 06
.06	.321-03	.255-04	.1022 05	.873 01	.7018 03	.04	.508-04	.201-05	.1045 05	.126 02	.4347 04						
.04	.229-03	.159-04	.1027 05	.937 01	.9816 03	.02	.244-04	.727-06	.1049 05	.146 02	.8990 04						
.02	.155-03	.521-05	.1031 05	.102 02	.1446 04	.00	.873-05	.174-06	.1054 05	.180 02	.2508 05						
.00	.974-04	.483-05	.1035 05	.112 02	.2288 04	.01	.150-05	.144-07	.1058 05	.236 02	.1453 06						
.00	.553-04	.219-05	.1039 05	.125 02	.4015 04												
.00	.266-04	.794-06	.1043 05	.145 02	.8299 04												
.00	.952-05	.188-06	.1048 05	.179 02	.2314 05												
.01	.164-05	.173-07	.1052 05	.250 02	.1339 06												

Table 19b. EXPANSION AFTER FREEZING - HYPERBOLIC NOZZLE

 $T_0' = 4000^\circ\text{K}$, $P_0' = 1000$ atm.

$\ell = 0.1$ $\ell = 1.0$ $\ell = 10$

T	P	P	u'	M	A	T	ρ	P	u'	M	A	T	ρ	P	u'	M	A
.766	.215 00	.158 00	.7360 04	.190 01	.1690 01	.648	.766-01	.460-01	.9041 04	.259 01	.3893 01	.553	.310-01	.158-01	.1008 05	.313 01	.8559 01
.76	.209 00	.152 00	.7416 04	.193 01	.1722 01	.64	.731-01	.437-01	.9098 04	.262 01	.4024 01	.54	.287-01	.143-01	.1017 05	.322 01	.8176 01
.74	.192 00	.136 00	.7600 04	.201 01	.1835 01	.62	.656-01	.380-01	.9247 04	.271 01	.4407 01	.52	.252-01	.121-01	.1040 05	.332 01	.7830 02
.72	.175 00	.121 00	.7779 04	.208 01	.1962 01	.60	.588-01	.330-01	.9393 04	.280 01	.4845 01	.50	.227-01	.102-01	.1043 05	.343 01	.7461 02
.70	.160 00	.107 00	.7954 04	.216 01	.2106 01	.58	.524-01	.284-01	.9537 04	.289 01	.5348 01	.48	.192-01	.091-01	.1056 05	.354 01	.7117 02
.68	.145 00	.945-01	.8125 04	.224 01	.2268 01	.56	.466-01	.244-01	.9679 04	.298 01	.5929 01	.46	.167-01	.076-02	.1068 05	.366 01	.6772 02
.66	.132 00	.831-01	.8293 04	.232 01	.2452 01	.54	.413-01	.208-01	.9819 04	.308 01	.6602 01	.44	.144-01	.052-02	.1081 05	.379 01	.6424 02
.64	.119 00	.729-01	.8457 04	.240 01	.2660 01	.52	.364-01	.177-01	.9936 04	.318 01	.7387 01	.42	.123-01	.035-02	.1093 05	.392 01	.6091 02
.62	.107 00	.638-01	.8617 04	.249 01	.2897 01	.50	.319-01	.149-01	.1009 05	.329 01	.8307 01	.40	.104-01	.021-02	.1105 05	.406 01	.5767 02
.60	.962-01	.553-01	.8775 04	.257 01	.3168 01	.48	.279-01	.125-01	.1022 05	.340 01	.9392 01	.38	.081-02	.008-02	.1117 05	.420 01	.5444 02
.58	.861-01	.478-01	.8930 04	.266 01	.3478 01	.46	.242-01	.104-01	.1036 05	.352 01	.1068 02	.36	.073-02	.005-02	.1129 05	.436 01	.5124 02
.56	.748-01	.412-01	.9081 04	.276 01	.3835 01	.44	.209-01	.089-02	.1049 05	.364 01	.1222 02	.34	.061-02	.001-02	.1141 05	.453 01	.4806 02
.54	.682-01	.353-01	.9231 04	.285 01	.4249 01	.42	.179-01	.073-02	.1074 05	.377 01	.1407 02	.32	.050-02	.000-02	.1152 05	.471 01	.4482 02
.52	.603-01	.300-01	.9377 04	.295 01	.4729 01	.40	.153-01	.057-02	.1074 05	.391 01	.1632 02	.30	.040-02	.000-02	.1164 05	.491 01	.4164 02
.50	.531-01	.254-01	.9521 04	.306 01	.5289 01	.38	.129-01	.038-02	.1086 05	.405 01	.1908 02	.28	.037-02	.000-02	.1175 05	.513 01	.3867 02
.48	.465-01	.214-01	.9663 04	.316 01	.5948 01	.36	.108-01	.024-02	.1098 05	.421 01	.2248 02	.26	.029-02	.000-02	.1185 05	.536 01	.3573 02
.46	.406-01	.179-01	.9802 04	.328 01	.6727 01	.34	.090-02	.016-02	.1110 05	.437 01	.2675 02	.24	.020-02	.000-02	.1196 05	.562 01	.3280 03
.44	.352-01	.148-01	.9939 04	.340 01	.7654 01	.32	.074-02	.012-02	.1122 05	.455 01	.3214 02	.22	.016-02	.000-02	.1217 05	.591 01	.2986 03
.42	.303-01	.122-01	.1007 05	.352 01	.8767 01	.30	.064-02	.009-02	.1134 05	.474 01	.3905 02	.20	.011-02	.000-02	.1236 05	.624 01	.2698 03
.40	.259-01	.093-02	.1021 05	.366 01	.1011 02	.28	.046-02	.007-02	.1145 05	.495 01	.4803 02	.18	.008-03	.000-02	.1256 05	.661 01	.2424 03
.38	.220-01	.801-02	.1034 05	.380 01	.1175 02	.26	.030-02	.006-02	.1156 05	.518 01	.5990 02	.16	.003-03	.000-02	.1276 05	.705 01	.2155 03
.36	.186-01	.640-02	.1046 05	.395 01	.1377 02	.24	.030-02	.006-02	.1167 05	.544 01	.7586 02	.14	.003-03	.000-02	.1296 05	.756 01	.1888 03
.34	.155-01	.505-02	.1059 05	.411 01	.1628 02	.22	.032-02	.007-03	.1178 05	.572 01	.9778 02	.12	.003-03	.000-02	.1316 05	.801 01	.1624 03
.32	.128-01	.384-02	.1071 05	.428 01	.1944 02	.20	.015-02	.007-03	.1188 05	.604 01	.1286 03	.10	.003-03	.000-02	.1336 05	.850 01	.1371 04
.30	.105-01	.302-02	.1084 05	.446 01	.2347 02	.18	.019-02	.007-03	.1199 05	.640 01	.1732 03	.09	.003-03	.000-02	.1356 05	.901 01	.1121 04
.28	.082-02	.228-02	.1096 05	.467 01	.2867 02	.16	.023-03	.008-03	.1209 05	.683 01	.2398 03	.08	.003-03	.000-02	.1376 05	.950 01	.0878 04
.26	.060-02	.169-02	.1107 05	.489 01	.3551 02	.14	.038-03	.008-03	.1218 05	.733 01	.3440 03	.07	.003-03	.000-02	.1396 05	.101 02	.0624 04
.24	.036-02	.123-02	.1119 05	.513 01	.4464 02	.12	.042-03	.008-03	.1228 05	.795 01	.5161 03	.06	.003-03	.000-02	.1416 05	.117 02	.0378 04
.22	.015-02	.083-03	.1130 05	.540 01	.5711 02	.10	.062-03	.008-03	.1237 05	.872 01	.8243 03	.05	.003-03	.000-02	.1436 05	.129 02	.0124 04
.20	.015-02	.063-03	.1141 05	.571 01	.7451 02	.09	.062-03	.008-03	.1241 05	.922 01	.1076 04	.04	.003-03	.000-02	.1456 05	.144 02	.0071 05
.18	.033-02	.402-03	.1152 05	.606 01	.9948 02	.08	.063-03	.008-03	.1250 05	.980 01	.1445 04	.03	.003-03	.000-02	.1476 05	.167 02	.0024 05
.16	.169-02	.259-03	.1162 05	.646 01	.1366 03	.07	.106-03	.009-03	.1260 05	.105 02	.2016 04	.02	.003-03	.000-02	.1496 05	.206 02	.0006 05
.14	.118-02	.158-03	.1172 05	.695 01	.1940 03	.06	.0721-04	.009-03	.1265 05	.114 02	.2957 04	.01	.003-03	.000-02	.1516 05	.279 02	.0006 06
.12	.085-03	.092-04	.1182 05	.754 01	.2881 03	.05	.0451-04	.009-03	.1269 05	.125 02	.4647 04						
.10	.0493-03	.0472-04	.1192 05	.829 01	.4550 03	.04	.0262-04	.009-03	.1269 05	.140 02	.8086 04						
.09	.039-03	.0326-04	.1196 05	.875 01	.5903 03	.03	.0262-04	.009-03	.1269 05	.163 02	.1652 05						
.08	.033-03	.0217-04	.1201 05	.931 01	.7879 03	.02	.0464-05	.009-03	.1269 05	.200 02	.4532 05						
.07	.0303-03	.0136-04	.1206 05	.998 01	.1091 04	.01	.0822-06	.009-03	.1269 05	.281 02	.2550 06						
.06	.0193-03	.0060-05	.1210 05	.108 02	.1588 04												
.05	.0891-04	.0426-05	.1215 05	.119 02	.2472 04												
.04	.0316-04	.0197-05	.1219 05	.133 02	.4254 04												
.03	.0355-04	.0132-06	.1224 05	.155 02	.8571 04												
.02	.0455-05	.0182-06	.1228 05	.190 02	.2304 05												
.01	.1173-05	.160-07	.1233 05	.272 02	.1253 06												

Table 20a. EXPANSION AFTER FREEZING - WEDGE NOZZLE

 $T_0' = 5000^\circ\text{K}$, $P_0' = 100$ atm.

$l = 10$
$$l = 1.0$$
$$\ell = 0.1$$

Table 20b. EXPANSION AFTER FREEZING - HYPERBOLIC NOZZLE

$$T_0' = 5000^\circ\text{K}, p_0' = 100 \text{ atm.}$$

T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A
		(ft./sec.)					(ft./sec.)					(ft./sec.)		
255 00	193 00	7003 04	178 01	1498 01	692	115 00	8458 04	233 01	2750 01	608	520-01	293-01	282 01	507 01
247 00	185 00	7077 04	182 01	1530 01	68	108 00	8555 04	238 01	2883 01	60	499-01	271-01	280 01	508 01
227 00	166 00	7289 04	189 01	1622 01	66	981-01	8714 04	246 01	3128 01	58	444-01	235-01	305 01	600 01
228 00	148 00	7457 04	197 01	1726 01	64	885-01	8870 04	254 01	3406 01	56	395-01	203-01	305 01	683 01
190 00	132 00	7640 04	204 01	1833 01	62	796-01	9023 04	263 01	3723 01	54	349-01	173-01	315 01	777 01
173 00	117 00	7818 04	212 01	1975 01	60	714-01	9173 04	272 01	4085 01	52	307-01	148-01	325 01	882 01
158 00	103 00	7932 04	220 01	2125 01	58	638-01	9321 04	281 01	4501 01	50	269-01	125-01	336 01	1008 01
135 00	90 00	8162 04	228 01	2284 01	56	567-01	9466 04	290 01	4980 01	46	235-01	105-01	347 01	1147 01
126 00	80 00	8339 04	236 01	2466 01	54	503-01	9609 04	300 01	5355 01	44	204-01	865-02	359 01	1240 01
116 00	69 00	8492 04	244 01	2704 01	52	444-01	9749 04	310 01	5818 01	42	176-01	711-02	371 01	1431 01
105 00	64 00	8652 04	253 01	2954 01	50	390-01	9888 04	320 01	6398 01	40	128-01	589-02	388 01	1591 01
931-01	523-01	8809 04	262 01	3240 01	48	341-01	1002 05	331 01	6929 01	38	108-01	475-02	413 01	1746 01
836-01	451-01	8963 04	271 01	3568 01	46	286-01	1016 05	343 01	7484 01	36	906-02	381-02	428 01	1943 01
743-01	386-01	9114 04	281 01	3948 01	44	256-01	1029 05	355 01	8014 02	34	752-02	303-02	445 01	2148 02
658-01	329-01	9263 04	291 01	4390 01	42	220-01	1042 05	368 01	8616 02	32	619-02	243-02	463 01	2380 02
580-01	279-01	9409 04	301 01	4905 01	40	188-01	1055 05	381 01	9218 02	30	503-02	184-02	483 01	2632 02
508-01	235-01	9552 04	312 01	5510 01	38	159-01	1068 05	396 01	9742 02	28	404-02	135-02	504 01	2902 02
438-01	195-01	9693 04	323 01	6225 01	36	134-01	1080 05	411 01	1031 02	26	320-02	773-03	527 01	3133 02
383-01	163-01	9831 04	335 01	7075 01	34	111-01	1092 05	427 01	1094 02	24	250-02	551-03	553 01	3404 02
332-01	134-01	9988 04	347 01	8094 01	32	920-02	1104 05	445 01	1263 02	22	192-02	392-03	581 01	3710 03
284-01	109-01	1010 05	361 01	9324 01	30	751-02	1117 05	466 01	1393 02	20	144-02	268-03	614 01	4042 03
232-01	883-02	1023 05	375 01	1082 02	28	606-02	1130 05	485 01	1574 02	18	106-02	177-03	643 01	4383 03
224-01	708-02	1036 05	390 01	1267 02	26	482-02	1143 05	507 01	1811 02	16	757-03	112-03	693 01	4811 03
170-01	558-02	1049 05	403 01	1496 02	24	378-02	1159 05	533 01	2031 02	14	522-03	678-04	744 01	5158 03
141-01	435-02	1062 05	425 01	1784 02	22	291-02	1161 05	560 01	2291 02	12	345-03	383-04	807 01	5625 03
130 00	334-02	1074 05	441 01	2151 02	20	220-02	1171 05	592 01	2537 02	10	213-03	198-04	887 01	6003 04
95-01	253-02	1086 05	461 01	2635 02	18	163-02	1182 05	628 01	2833 03	8	163-03	136-04	936 01	6314 04
751-02	188-02	1098 05	483 01	3246 02	16	117-02	1192 05	669 01	3122 03	6	120-03	893-05	995 01	6766 04
592-02	137-02	1109 05	507 01	4075 02	14	810-03	1202 05	719 01	3477 03	4	858-04	551-05	107 02	7247 04
459-02	971-03	1121 05	534 01	5205 02	12	538-03	1211 05	780 01	4106 03	2	582-04	324-05	115 02	7628 04
349-02	671-03	1132 05	564 01	6780 02	10	336-03	1216 04	857 01	4530 03	06	368-04	171-05	127 02	7816 04
259-02	449-03	1143 05	599 01	8038 02	09	257-03	1218 05	905 01	4850 03	05	310-04	780-06	142 02	7983 04
187-02	288-03	1153 05	639 01	9298 01	08	191-03	1230 05	962 01	5139 04	04	212-04	283-06	165 02	8049 04
131-02	176-03	1163 05	687 01	11756 03	07	137-03	1239 05	1032 02	5385 04	03	102-04	1280 05	202 02	8265 05
87-02	101-03	1173 05	746 01	12603 03	06	932-04	1247 05	1122 05	5689 04	02	646-06	530-08	297 02	8214 06
51-03	531-04	1183 05	820 01	14101 03	05	593-04	1253 05	1232 05	6030 04	01				
42-03	367-04	1198 05	866 01	15313 04	04	341-04	1267 05	1338 02	6289 04					
37-03	244-04	1212 05	921 01	17033 03	03	167-04	1278 05	160 02	6479 05					
228-03	154-04	1227 05	988 01	19795 03	02	612-05	1286 05	196 02	6795 05					
100-03	923-05	1240 05	107 02	21423 04	01	110-05	1295 05	277 02	7193 06					
582-04	224-05	1251 05	132 02	23797 02										
289-04	834-06	1255 05	153 02	2628 04										
15-05	206-06	1260 05	188 02	2042 05										
15-05	187-07	1275 05	269 02	1103 06										

$\ell = 0.1$ $\ell = 1.0$ $\ell = 10$

T	ρ	P	u' (ft./sec.)	M	A	T	ρ	P	u' (ft./sec.)	M	A	T	ρ	P	u' (ft./sec.)	M	A
.714	.150 00	.102 00	.7846 04	.216 01	.2223 01	.60	.576-01	.325-01	.9301 04	.277 01	.4946 01	.513	.270-01	.129-01	.1011 05	.324 01	.2704 01
.73	.143 00	.938-01	.8666 04	.220 01	.2341 01	.58	.513-01	.280-01	.9446 04	.286 01	.5467 01	.50	.248-01	.116-01	.1020 05	.333 01	.1047 02
.68	.127 00	.826-01	.8235 04	.228 01	.2529 01	.56	.455-01	.240-01	.9589 04	.295 01	.5069 01	.48	.216-01	.099-02	.1033 05	.344 01	.1168 03
.66	.115 00	.725-01	.8399 04	.236 01	.2743 01	.54	.402-01	.204-01	.9730 04	.305 01	.4678 01	.46	.187-01	.080-02	.1046 05	.350 01	.1355 04
.64	.104 00	.634-01	.8561 04	.244 01	.2985 01	.52	.354-01	.174-01	.9868 04	.315 01	.4584 01	.44	.161-01	.062-02	.1058 05	.358 01	.1555 04
.62	.932-01	.552-01	.8719 04	.253 01	.3261 01	.50	.310-01	.148-01	.1000 05	.325 01	.4543 01	.42	.138-01	.040-02	.1071 05	.363 01	.1768 02
.60	.835-01	.478-01	.8875 04	.261 01	.3577 01	.48	.270-01	.122-01	.1014 05	.337 01	.4675 01	.40	.117-01	.037-02	.1083 05	.365 01	.2034 03
.58	.743-01	.412-01	.9027 04	.270 01	.3941 01	.46	.234-01	.101-01	.1027 05	.348 01	.4802 01	.38	.095-02	.030-02	.1096 05	.410 01	.2486 02
.56	.662-01	.354-01	.9177 04	.280 01	.4360 01	.44	.202-01	.0835-02	.1040 05	.360 01	.4932 02	.36	.0823-02	.027-02	.1108 05	.425 01	.2906 02
.54	.585-01	.302-01	.9324 04	.289 01	.4846 01	.42	.173-01	.0682-02	.1053 05	.373 01	.5158 02	.34	.0682-02	.0217-02	.1120 05	.442 01	.3472 02
.52	.517-01	.257-01	.9469 04	.299 01	.5412 01	.40	.147-01	.055-02	.1066 05	.387 01	.5694 02	.32	.0559-02	.0167-02	.1131 05	.460 01	.4186 02
.50	.454-01	.217-01	.9611 04	.310 01	.6075 01	.38	.124-01	.043-02	.1078 05	.401 01	.6395 02	.30	.0453-02	.0127-02	.1143 05	.479 01	.5114 02
.48	.396-01	.182-01	.9751 04	.321 01	.6857 01	.36	.104-01	.031-02	.1090 05	.417 01	.7345 02	.28	.0363-02	.0091-03	.1154 05	.509 01	.6421 02
.46	.344-01	.151-01	.9888 04	.332 01	.7784 01	.34	.0860-02	.0273-02	.1102 05	.433 01	.8296 02	.26	.0287-02	.0068-03	.1165 05	.523 01	.7923 02
.44	.297-01	.125-01	.1002 05	.344 01	.8891 01	.32	.0706-02	.0213-02	.1114 05	.451 01	.9369 02	.24	.0223-02	.0051-03	.1176 05	.549 01	.1004 03
.42	.255-01	.102-01	.1016 05	.357 01	.1022 02	.30	.0573-02	.0162-02	.1126 05	.470 01	.1040 02	.22	.0171-02	.0031-03	.1186 05	.577 01	.1207 03
.40	.217-01	.081-02	.1029 05	.370 01	.1184 02	.28	.0460-02	.0121-02	.1137 05	.491 01	.0563 02	.20	.0138-02	.0023-03	.1197 05	.610 01	.1730 04
.38	.184-01	.0668-02	.1042 05	.384 01	.1382 02	.26	.0364-02	.0091-03	.1148 05	.514 01	.0333 02	.18	.0066-03	.0016-04	.1207 05	.646 01	.2344 04
.36	.154-01	.0531-02	.1054 05	.399 01	.1627 02	.24	.0284-02	.0064-03	.1159 05	.539 01	.0141 03	.16	.0046 02	.0009-04	.1217 05	.689 01	.3268 04
.34	.129-01	.0417-02	.1067 05	.415 01	.1932 02	.22	.0218-02	.0030-03	.1170 05	.567 01	.0041 03	.14	.0030-03	.0003-04	.1226 05	.740 01	.4771 04
.32	.106-01	.0324-02	.1079 05	.433 01	.2319 02	.20	.0163-02	.0020-03	.1181 05	.599 01	.001374 03	.12	.0030-03	.0003-04	.1236 05	.802 01	.7140 04
.30	.083-02	.0247-02	.1091 05	.451 01	.2814 02	.18	.0120-02	.0013-03	.1191 05	.635 01	.001857 03	.10	.00185-03	.0001-04	.1245 05	.881 01	.1150 04
.28	.055-02	.0186-02	.1103 05	.472 01	.3456 02	.16	.0084-03	.0009-03	.1201 05	.677 01	.002583 03	.09	.00141-03	.0001-04	.1253 05	.950 01	.1508 04
.26	.032-02	.0137-02	.1115 05	.494 01	.4305 02	.14	.0058-03	.0007-03	.1211 05	.727 01	.003721 03	.08	.00104-03	.0001-04	.1263 05	.989 01	.2086 04
.24	.014-01	.0099-03	.1126 05	.519 01	.5445 02	.12	.00387-03	.0004-04	.1220 05	.788 01	.005610 03	.07	.000738-04	.0001-04	.1272 05	.106 02	.2845 04
.22	.002-01	.0068-03	.1137 05	.546 01	.7009 02	.10	.00239-03	.0002-04	.1229 05	.867 01	.00907 03	.06	.000498-04	.0001-04	.1282 05	.115 02	.4714 04
.20	.001-02	.0049-03	.1148 05	.577 01	.9206 02	.09	.00182-03	.0001-04	.1234 05	.915 01	.01179 04	.05	.000314-04	.0001-04	.1292 05	.126 02	.6671 04
.18	.001-02	.00317-03	.1158 05	.612 01	.1238 03	.08	.00135-03	.0001-04	.1238 05	.972 01	.01589 04	.04	.000178-04	.0001-04	.1302 05	.141 02	.1171 05
.16	.001-02	.00202-03	.1169 05	.653 01	.1712 03	.07	.000959-04	.0001-05	.1243 05	.104 02	.02224 04	.03	.0000859-05	.0001-04	.1312 05	.164 02	.2414 05
.14	.001-03	.00122-03	.1179 05	.701 01	.2453 03	.06	.000649-04	.0001-05	.1247 05	.113 02	.03274 04	.02	.0000307-05	.0001-04	.1323 05	.202 02	.4738 05
.12	.001-03	.00095-04	.1189 05	.761 01	.3675 03	.05	.000410-04	.0001-06	.1251 05	.124 02	.05169 04	.01	.0000530-06	.0001-04	.1333 05	.276 02	.8497 05
.10	.001-03	.00060-04	.1198 05	.837 01	.5860 03	.04	.000233-04	.0001-06	.1256 05	.139 02	.09043 04						
.08	.001-03	.00048-04	.1203 05	.884 01	.7642 03	.03	.000113-04	.0001-06	.1260 05	.161 02	.1861 05						
.06	.001-03	.00036-04	.1207 05	.939 01	.1026 04	.02	.0000407-05	.0001-06	.1264 05	.186 02	.5155 05						
.04	.001-03	.00024-04	.1213 05	.101 02	.1430 04	.01	.0000708-06	.0001-06	.1269 05	.301 02	.2950 06						
.02	.001-03	.00012-04	.1216 05	.109 02	.2095 04												
.00	.001-03	.00006-04	.1221 05	.120 02	.3289 04												
.00	.001-04	.00003-05	.1225 05	.135 02	.5715 04												
.00	.001-04	.00002-06	.1230 05	.156 02	.1166 05												
.00	.001-05	.00001-07	.1234 05	.192 02	.3192 05												
.00	.001-05	.00000-08	.1239 05	.277 02	.1790 06												

Table 21a. EXPANSION AFTER FREEZING - WEDGE NOZZLE

 $T_0' = 5000^\circ\text{K}$, $p_0' = 200$ atm.

$l = 0.1$ $l = 1.0$ $l = 10$

T	ρ	P	u' (ft/sec.)	M	A	T	ϵ	P	u' (ft/sec.)	M	A	T	ρ	P	u' (ft/sec.)	M	A
.746	.194 00	.140 00	.7495 04	.198 01	.1804 01	.652	.905	.559-01	.8723 04	.248 01	.3356 01	.57	.447-01	.236-01	.9001 04	.201 01	.624-01
.74	.191 00	.136 00	.7550 04	.200 01	.1840 01	.64	.850-01	.515-01	.8817 04	.253 01	.3537 01	.56	.410-01	.219-01	.9071 04	.208 01	.638-01
.72	.174 00	.120 00	.7730 04	.207 01	.1970 01	.62	.763-01	.448-01	.8971 04	.261 01	.3872 01	.57	.395-01	.195-01	.9111 04	.209 01	.784-01
.70	.158 00	.106 00	.7906 04	.215 01	.2118 01	.60	.682-01	.389-01	.9122 04	.270 01	.4256 01	.58	.374-01	.184-01	.9149 04	.219 01	.824-01
.68	.144 00	.937-01	.8077 04	.223 01	.2285 01	.58	.608-01	.334-01	.9270 04	.279 01	.4698 01	.59	.354-01	.173-01	.9198 04	.229 01	.877-01
.66	.130 00	.823-01	.8245 04	.231 01	.2474 01	.56	.540-01	.287-01	.9416 04	.288 01	.5208 01	.60	.334-01	.162-01	.9257 04	.239 01	.937-01
.64	.117 00	.720-01	.8410 04	.239 01	.2689 01	.54	.478-01	.240-01	.9559 04	.298 01	.5799 01	.61	.314-01	.151-01	.9316 04	.249 01	.1000 02
.62	.105 00	.627-01	.8571 04	.248 01	.2934 01	.52	.421-01	.207-01	.9700 04	.308 01	.6493 01	.62	.294-01	.140-01	.9375 04	.259 01	.1500 02
.60	.944-01	.544-01	.8729 04	.256 01	.3214 01	.50	.369-01	.175-01	.9839 04	.319 01	.7299 01	.63	.274-01	.129-01	.9434 04	.269 01	.1848 02
.58	.830-01	.470-01	.8885 04	.265 01	.3536 01	.48	.322-01	.146-01	.9975 04	.330 01	.8254 01	.64	.254-01	.118-01	.9493 04	.279 01	.2167 02
.56	.713-01	.404-01	.9037 04	.274 01	.3907 01	.46	.279-01	.122-01	.1011 05	.341 01	.9389 01	.65	.234-01	.107-01	.9552 04	.289 01	.2503 02
.54	.595-01	.345-01	.9186 04	.284 01	.4337 01	.44	.241-01	.100-01	.1024 05	.353 01	.1075 02	.66	.214-01	.096-01	.9611 04	.299 01	.2847 02
.52	.587-01	.293-01	.9333 04	.294 01	.4837 01	.42	.206-01	.821-02	.1037 05	.366 01	.1238 02	.67	.194-01	.085-01	.9670 04	.309 01	.3206 02
.50	.515-01	.247-01	.9478 04	.304 01	.5423 01	.40	.176-01	.665-02	.1050 05	.380 01	.1437 02	.68	.174-01	.074-01	.9729 04	.319 01	.3606 02
.48	.451-01	.208-01	.9619 04	.315 01	.6114 01	.38	.148-01	.524-02	.1063 05	.394 01	.1681 02	.69	.154-01	.063-01	.9788 04	.329 01	.3968 02
.46	.392-01	.173-01	.9759 04	.327 01	.6932 01	.36	.124-01	.428-02	.1075 05	.409 01	.1982 02	.70	.134-01	.052-01	.9847 04	.339 01	.4347 02
.44	.339-01	.143-01	.9896 04	.339 01	.7908 01	.34	.103-01	.332-02	.1087 05	.426 01	.2360 02	.71	.114-01	.041-01	.9906 04	.349 01	.4747 02
.42	.291-01	.117-01	.1003 05	.351 01	.9081 01	.32	.849-02	.257-02	.1099 05	.443 01	.2838 02	.72	.094-01	.030-01	.9965 04	.359 01	.5147 02
.40	.248-01	.953-02	.1016 05	.364 01	.1050 02	.30	.691-02	.196-02	.1111 05	.462 01	.3452 02	.73	.074-01	.019-01	.1000 05	.369 01	.5547 02
.38	.210-01	.767-02	.1029 05	.378 01	.1224 02	.28	.555-02	.147-02	.1123 05	.483 01	.4230 02	.74	.054-01	.008-01	.1059 05	.379 01	.5947 02
.36	.177-01	.611-02	.1042 05	.393 01	.1439 02	.26	.440-02	.108-02	.1134 05	.505 01	.5306 02	.75	.034-01	.007-01	.1118 05	.389 01	.6347 02
.34	.147-01	.485-02	.1055 05	.409 01	.1707 02	.24	.344-02	.781-03	.1145 05	.530 01	.6729 02	.76	.014-01	.006-01	.1177 05	.399 01	.6747 02
.32	.121-01	.373-02	.1067 05	.426 01	.2046 02	.22	.264-02	.550-03	.1156 05	.558 01	.8684 02	.77	.004-01	.005-01	.1236 05	.409 01	.7147 02
.30	.990-02	.285-02	.1079 05	.445 01	.2478 02	.20	.199-02	.376-03	.1167 05	.589 01	.1144 03	.78	.004-01	.004-01	.1295 05	.419 01	.7547 02
.28	.799-02	.215-02	.1091 05	.465 01	.3040 02	.18	.146-02	.249-03	.1177 05	.625 01	.1543 03	.79	.004-01	.003-01	.1354 05	.429 01	.7947 02
.26	.635-02	.159-02	.1103 05	.487 01	.3780 02	.16	.104-02	.158-03	.1187 05	.667 01	.2140 03	.80	.004-01	.002-01	.1413 05	.439 01	.8347 02
.24	.498-02	.115-02	.1115 05	.512 01	.4773 02	.14	.720-03	.950-04	.1197 05	.716 01	.3076 03	.81	.004-01	.001-01	.1472 05	.449 01	.8747 02
.22	.384-02	.810-03	.1126 05	.539 01	.6133 02	.12	.475-03	.540-04	.1207 05	.777 01	.4624 03	.82	.004-01	.001-01	.1531 05	.459 01	.9147 02
.20	.290-02	.557-03	.1137 05	.569 01	.8041 02	.10	.294-03	.270-04	.1216 05	.854 01	.7402 03	.83	.004-01	.001-01	.1590 05	.469 01	.9547 02
.18	.214-02	.370-03	.1148 05	.604 01	.1079 03	.09	.224-03	.181-04	.1221 05	.901 01	.9672 03	.84	.004-01	.001-01	.1649 05	.479 01	.9947 02
.16	.154-02	.236-03	.1158 05	.645 01	.1490 03	.08	.166-03	.126-04	.1225 05	.958 01	.1301 04	.85	.004-01	.001-01	.1708 05	.489 01	.1034 03
.14	.107-02	.143-03	.1168 05	.693 01	.2129 03	.07	.119-03	.786-05	.1230 05	.103 02	.1817 04	.86	.004-01	.001-01	.1767 05	.499 01	.1094 03
.12	.797-03	.814-04	.1178 05	.752 01	.3182 03	.06	.804-04	.457-05	.1234 05	.111 02	.2670 04	.87	.004-01	.001-01	.1826 05	.509 01	.1154 03
.10	.445-03	.423-04	.1188 05	.827 01	.5061 03	.05	.509-04	.241-05	.1238 05	.122 02	.4205 04	.88	.004-01	.001-01	.1885 05	.519 01	.1214 03
.08	.337-03	.291-04	.1192 05	.893 01	.6591 03	.04	.291-04	.110-05	.1243 05	.137 02	.7333 04	.89	.004-01	.001-01	.1946 05	.529 01	.1274 03
.06	.231-03	.192-04	.1197 05	.928 01	.8834 03	.03	.141-04	.401-06	.1247 05	.159 02	.1503 05	.90	.004-01	.001-01	.2007 05	.539 01	.1334 03
.04	.175-03	.121-04	.1202 05	.995 01	.1229 04	.02	.211-05	.980-07	.1252 05	.195 02	.4141 05	.91	.004-01	.001-01	.2068 05	.549 01	.1394 03
.02	.122-03	.764-05	.1206 05	.108 02	.1797 04	.01	.899-06	.810-08	.1256 05	.281 02	.2347 06	.92	.004-01	.001-01	.2129 05	.559 01	.1454 03
.00	.777-04	.378-05	.1211 05	.118 02	.4881 04												
.04	.447-04	.172-05	.1215 05	.133 02	.4881 04												
.08	.219-04	.629-06	.1220 05	.154 02	.9926 04												
.12	.801-05	.153-06	.1224 05	.190 02	.2704 05												
.16	.143-05	.136-07	.1229 05	.270 02	.1504 06												

Table 21b. EXPANSION AFTER FREEZING - HYPERBOLIC NOZZLE

 $T_o' = 5000^\circ K, p_o' = 200 \text{ atm.}$

$\ell = 10$ $\ell = 1.0$ $\ell = 0.1$

T	ρ	P	u' (ft./sec.)	M	A	T	ρ	P	u' (ft./sec.)	M	A	T	ρ	P	u' (ft./sec.)	M	A
.679	.121-00	.785-01	.8255-04	.230-01	.2639-01	.575	.510-01	.278-01	.9380-04	.284-01	.5511-01	.49	.255-01	.118-01	.1011-05	.330-01	.1023-02
.66	.110-00	.693-01	.8411-04	.236-01	.2853-01	.56	.466-01	.247-01	.9488-04	.291-01	.5964-01	.48	.238-01	.109-01	.1018-05	.338-01	.1090-02
.64	.989-01	.605-01	.8572-04	.245-01	.3109-01	.54	.411-01	.210-01	.9630-04	.301-01	.6653-01	.46	.206-01	.992-02	.1031-05	.349-01	.1243-02
.62	.888-01	.526-01	.8730-04	.253-01	.3401-01	.52	.362-01	.178-01	.9769-04	.311-01	.7458-01	.44	.177-01	.734-02	.1044-05	.362-01	.1427-02
.60	.794-01	.456-01	.8885-04	.262-01	.3735-01	.50	.317-01	.150-01	.9907-04	.322-01	.8404-01	.42	.151-01	.599-02	.1057-05	.375-01	.1649-02
.58	.708-01	.393-01	.9037-04	.271-01	.4119-01	.48	.276-01	.125-01	.1004-05	.333-01	.9321-01	.40	.126-01	.484-02	.1069-05	.388-01	.1920-02
.56	.629-01	.337-01	.9187-04	.280-01	.4563-01	.46	.239-01	.104-01	.1018-05	.344-01	.1085-02	.38	.108-01	.388-02	.1081-05	.403-01	.2252-02
.54	.556-01	.287-01	.9334-04	.290-01	.5078-01	.44	.206-01	.856-02	.1031-05	.356-01	.1244-02	.36	.904-02	.307-02	.1094-05	.418-01	.2666-02
.52	.490-01	.243-01	.9478-04	.300-01	.5679-01	.42	.176-01	.699-02	.1044-05	.369-01	.1437-02	.34	.749-02	.240-02	.1106-05	.435-01	.3184-02
.50	.429-01	.205-01	.9620-04	.310-01	.6384-01	.40	.149-01	.566-02	.1056-05	.383-01	.1671-02	.32	.614-02	.185-02	.1117-05	.452-01	.3843-01
.48	.374-01	.172-01	.9759-04	.321-01	.7215-01	.38	.126-01	.453-02	.1069-05	.397-01	.1958-02	.30	.498-02	.141-02	.1129-05	.471-01	.4041-02
.46	.325-01	.143-01	.9897-04	.333-01	.8203-01	.36	.105-01	.359-02	.1081-05	.412-01	.2315-02	.28	.398-02	.105-02	.1143-05	.492-01	.5748-02
.44	.280-01	.118-01	.1003-05	.345-01	.9385-01	.34	.873-02	.281-02	.1093-05	.429-01	.2763-02	.26	.315-02	.772-03	.1152-05	.515-01	.7268-02
.42	.240-01	.964-02	.1016-05	.357-01	.1081-02	.32	.716-02	.217-02	.1105-05	.446-01	.3331-02	.24	.245-02	.554-03	.1163-05	.541-01	.8255-02
.40	.204-01	.781-02	.1030-05	.371-01	.1254-02	.30	.581-02	.165-02	.1117-05	.465-01	.4061-02	.22	.187-02	.368-03	.1173-05	.569-01	.1200-03
.38	.173-01	.637-02	.1042-05	.385-01	.1456-02	.28	.466-02	.123-02	.1128-05	.486-01	.5014-02	.20	.140-02	.265-03	.1184-05	.600-01	.1588-03
.36	.145-01	.497-02	.1055-05	.400-01	.1729-02	.26	.368-02	.907-03	.1140-05	.509-01	.6277-02	.18	.103-02	.174-03	.1194-05	.637-01	.2152-03
.34	.120-01	.390-02	.1067-05	.416-01	.2057-02	.24	.287-02	.682-03	.1151-05	.534-01	.7084-02	.16	.730-03	.110-03	.1204-05	.679-01	.3301-03
.32	.987-02	.302-02	.1080-05	.433-01	.2474-02	.22	.220-02	.437-03	.1162-05	.562-01	.8034-02	.14	.501-03	.661-04	.1214-05	.729-01	.4337-03
.30	.803-02	.230-02	.1092-05	.452-01	.3008-02	.20	.165-02	.312-03	.1172-05	.593-01	.1366-03	.12	.329-03	.372-04	.1223-05	.790-01	.6561-03
.28	.645-02	.173-02	.1104-05	.472-01	.3703-02	.18	.121-02	.205-03	.1183-05	.629-01	.1849-03	.10	.202-03	.191-04	.1232-05	.869-01	.1057-04
.26	.512-02	.137-02	.1115-05	.495-01	.4622-02	.16	.858-03	.130-03	.1193-05	.671-01	.2575-03	.09	.154-03	.130-04	.1237-05	.917-01	.1387-04
.24	.399-02	.916-03	.1126-05	.519-01	.5860-02	.14	.590-03	.782-04	.1202-05	.720-01	.3715-03	.08	.118-03	.856-05	.1241-05	.974-01	.1873-04
.22	.305-02	.645-03	.1138-05	.547-01	.7562-02	.12	.388-03	.440-04	.1212-05	.781-01	.5010-03	.07	.806-04	.532-05	.1245-05	.104-02	.2627-04
.20	.230-02	.441-03	.1148-05	.577-01	.9960-02	.10	.239-03	.226-04	.1221-05	.859-01	.9023-03	.06	.544-04	.368-05	.1250-05	.113-02	.3878-04
.18	.169-02	.291-03	.1159-05	.613-01	.1343-03	.09	.182-03	.155-04	.1226-05	.906-01	.1182-04	.05	.342-04	.161-05	.1254-05	.124-02	.6142-04
.16	.121-02	.185-03	.1169-05	.653-01	.1864-03	.08	.134-03	.102-04	.1230-05	.964-01	.1595-04	.04	.194-04	.732-06	.1258-05	.139-02	.1078-05
.14	.835-03	.112-03	.1179-05	.702-01	.2678-03	.07	.955-04	.633-05	.1235-05	.103-02	.2335-04	.03	.936-05	.264-06	.1263-05	.162-02	.2230-05
.12	.551-03	.632-04	.1189-05	.762-01	.4027-03	.06	.646-04	.367-05	.1239-05	.112-02	.3295-04	.02	.335-05	.633-07	.1267-05	.199-02	.6219-05
.10	.341-03	.326-04	.1198-05	.838-01	.6446-03	.05	.407-04	.193-05	.1243-05	.123-02	.5210-04	.01	.576-06	.550-08	.1271-05	.280-02	.3602-06
.08	.260-03	.224-04	.1203-05	.884-01	.8424-03	.04	.231-04	.875-06	.1248-05	.138-02	.9132-04						
.06	.193-03	.147-04	.1208-05	.940-01	.1133-04	.03	.112-04	.317-06	.1252-05	.160-02	.1884-05						
.04	.931-04	.534-05	.1217-05	.109-02	.2327-04	.02	.401-05	.770-07	.1256-05	.196-02	.5237-05						
.02	.589-04	.282-05	.1221-05	.120-02	.3665-04	.01	.694-06	.550-08	.1261-05	.294-02	.3014-06						
.00	.163-04	.468-06	.1230-05	.135-02	.6394-04												
	.591-05	.113-06	.1234-05	.192-02	.3615-05												
	.104-05	.110-07	.1239-05	.264-02	.2053-06												

Table 22a. EXPANSION AFTER FREEZING - WEDGE NOZZLE

 $T_o' = 5000^\circ K, p_o' = 300 \text{ atm.}$

$l = 0.1$

$l = 1.0$

$l = 10$

T	P	P	u' (ft/sec.)	M	A	T	P	P	u' (ft/sec.)	M	A	T	P	P	u' (ft/sec.)	M	A
718	145 00	114 00	7765 04	210 01	2058 01	625	775-01	460-01	8872 04	258 01	3834 01	546	402-01	207-01	9646 04	300 01	6799 01
715	152 00	122 00	7918 04	216 01	2194 01	62	754-01	444-01	8911 04	259 01	3924 01	54	387-01	197-01	9688 04	303 01	7030 01
68	138 00	896-01	8089 04	223 01	2370 01	60	674-01	384-01	9063 04	268 01	4316 01	52	340-01	167-01	9827 04	313 01	7884 01
55	124 00	788-01	8257 04	231 01	2569 01	58	600-01	331-01	9212 04	277 01	4767 01	50	298-01	141-01	9963 04	324 01	8484 01
54	112 00	689-01	8421 04	240 01	2796 01	56	533-01	284-01	9358 04	286 01	5288 01	48	259-01	118-01	1010 05	335 01	9007 02
62	101 00	599-01	8582 04	248 01	3054 01	54	471-01	242-01	9502 04	296 01	5893 01	46	224-01	975-02	1023 05	346 01	9642 02
60	900-01	519-01	8740 04	257 01	3350 01	52	414-01	205-01	9644 04	306 01	6599 01	44	193-01	650-02	1036 05	359 01	1036 05
55	802-01	448-01	8895 04	266 01	3680 01	50	363-01	175-01	9783 04	317 01	7429 01	42	165-01	530-02	1049 05	371 01	1152 02
55	712-01	384-01	9047 04	275 01	4083 01	48	316-01	145-01	9921 04	328 01	8408 01	40	140-01	424-02	1062 05	385 01	1277 02
57	612-01	348-01	9196 04	285 01	4539 01	46	274-01	125-01	1006 05	339 01	9573 01	38	118-01	424-02	1074 05	399 01	1477 02
57	517-01	278-01	9342 04	295 01	5070 01	44	226-01	987-02	1019 05	351 01	1097 02	36	988-02	336-02	1086 05	415 01	1771 02
52	468-01	236-01	9486 04	305 01	5692 01	42	202-01	806-02	1032 05	364 01	1469 02	34	818-02	203-02	1110 05	440 01	2077 02
50	423-01	196-01	9628 04	316 01	6426 01	40	172-01	653-02	1045 05	377 01	1721 02	32	671-02	154-02	1122 05	468 01	2456 02
49	370-01	163-01	9767 04	327 01	7298 01	38	145-01	523-02	1057 05	392 01	2032 02	30	545-02	114-02	1133 05	489 01	2930 02
47	319-01	135-01	9904 04	339 01	8339 01	36	121-01	415-02	1070 05	407 01	2322 02	28	436-02	847-03	1145 05	512 01	3677 02
47	274-01	110-01	1004 05	352 01	9592 01	34	101-01	325-02	1082 05	423 01	2422 02	26	345-02	609-03	1156 05	537 01	4494 02
40	233-01	896-02	1017 05	365 01	1111 02	32	826-02	251-02	1094 05	441 01	2797 02	24	289-02	427-03	1166 05	565 01	5494 03
36	197-01	720-02	1030 05	379 01	1298 02	30	671-02	191-02	1106 05	460 01	3552 02	22	245-02	291-03	1177 05	596 01	6854 03
34	165-01	572-02	1043 05	394 01	1517 02	28	538-02	143-02	1118 05	480 01	4381 02	20	174-02	192-03	1187 05	632 01	8494 03
32	137-01	449-02	1056 05	410 01	1817 02	26	426-02	105-02	1129 05	503 01	5478 02	18	113-02	121-03	1197 05	674 01	1065 04
30	113-01	348-02	1068 05	427 01	2182 02	24	332-02	758-03	1140 05	528 01	6958 02	16	801-03	121-03	1207 05	724 01	1265 04
28	921-02	266-02	1080 05	446 01	2649 02	22	255-02	532-03	1151 05	555 01	8997 02	14	531-03	121-03	1216 05	785 01	1654 04
26	841-02	199-02	1092 05	466 01	3256 02	20	191-02	363-03	1162 05	586 01	1187 03	12	361-03	140-04	1226 05	863 01	2065 04
24	740-02	166-02	1104 05	488 01	4059 02	18	140-02	240-03	1172 05	622 01	1605 03	10	233-03	144-04	1230 05	911 01	2744 03
22	640-02	137-02	1115 05	512 01	5138 02	16	999-03	152-03	1183 05	663 01	2231 03	08	169-03	144-04	1235 05	968 01	3967 03
20	544-02	106-02	1126 05	540 01	6021 02	14	688-03	915-04	1192 05	713 01	3214 03	06	125-03	121-03	1239 05	104 02	5946 03
18	460-02	87-02	1137 05	570 01	7075 02	12	453-03	516-04	1202 05	773 01	4845 03	04	888-04	121-03	1243 05	112 02	8654 03
16	365-02	612-03	1148 05	605 01	8172 03	10	280-03	266-04	1211 05	850 01	7778 03	02	600-04	121-03	1248 05	123 02	1265 04
14	265-02	439-03	1158 05	645 01	9423 03	08	213-03	182-04	1216 05	897 01	1018 04	00	378-04	121-03	1252 05	139 02	1704 04
12	143-02	216-03	1169 05	694 01	1123 03	06	137-03	120-04	1220 05	954 01	1372 04	00	215-04	121-03	1256 05	160 02	2026 05
10	945-03	130-03	1178 05	752 01	13493 03	04	112-03	746-05	1225 05	102 02	1920 04	00	104-04	121-03	1261 05	198 02	2564 05
8	641-03	739-04	1188 05	828 01	1578 03	02	79-04	432-05	1234 05	111 02	2827 04	00	371-05	121-03	1265 05	290 02	3257 06
6	398-03	392-04	1193 05	874 01	1728 03	00	479-04	227-05	1234 05	122 02	4463 04	00	640-06	121-03	1265 05	390 02	4257 06
4	225-03	173-04	1197 05	929 01	1982 03	00	273-04	104-05	1234 05	137 02	7808 04	00	160 02	121-03	1265 05	490 02	5257 06
2	161-03	138-04	1202 05	996 01	2365 04	00	145-05	908-07	1247 05	158 02	1607 05	00	160 02	121-03	1265 05	590 02	6257 06
0	109-03	639-05	1206 05	108 02	2802 04	00	827-06	820-08	1251 05	273 02	2548 06	00	160 02	121-03	1265 05	690 02	7257 06
0	692-04	332-05	1211 05	119 02	3148 04	00	132-04	377-06	1251 05	273 02	2548 06	00	160 02	121-03	1265 05	790 02	8257 06
0	396-04	132-05	1215 05	133 02	3479 04	00	132-04	377-06	1251 05	273 02	2548 06	00	160 02	121-03	1265 05	890 02	9257 06
0	193-04	566-06	1220 05	154 02	3976 05	00	132-04	377-06	1251 05	273 02	2548 06	00	160 02	121-03	1265 05	990 02	10257 06
0	700-05	135-06	1224 05	190 02	4307 05	00	132-04	377-06	1251 05	273 02	2548 06	00	160 02	121-03	1265 05	1090 02	11257 06
0	124-05	110-07	1229 05	276 02	4733 06	00	132-04	377-06	1251 05	273 02	2548 06	00	160 02	121-03	1265 05	1190 02	12257 06

Table 22b. EXPANSION AFTER FREEZING - HYPERBOLIC NOZZLE

$T_0' = 5000^\circ K$, $p_0' = 300$ atm.

$l = 0.1$																$l = 1.0$																$l = 10$																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																														
T	P	P	u'	M	A	T	P	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'	M	A	T	P	u'

Table 23a. EXPANSION AFTER FREEZING - WEDGE NOZZLE

 $T_0' = 5000^\circ K$, $p_0' = 1000$ atm.

$l = 10$

$l = 1.0$

$l = 0.1$

T	P	P	u'	M	A	T	P	P	u'	M	A	T	P	P	u'	M	A
(ft/sec.)																	
638	107.00	665-01	8335.04	339.01	2913.01	552	563-01	300-01	9164.04	282.01	5059.01	481	318-01	148-01	9750.04	314.01	8-17.01
62	974-01	586-01	8481.04	345.01	3159.01	54	522-01	272-01	9252.04	286.01	5402.01	48	316-01	146-01	9757.04	320.01	8-71.01
60	870-01	506-01	8640.04	353.01	3470.01	52	459-01	230-01	9397.04	296.01	6052.01	46	273-01	121-01	9894.04	332.01	9-06-01
58	775-01	436-01	8796.04	362.01	3828.01	50	401-01	194-01	9540.04	307.01	6815.01	44	235-01	99-03	1003.05	344.01	11-07.01
56	687-01	373-01	8950.04	371.01	4243.01	48	349-01	162-01	9681.04	317.01	7718.01	42	201-01	81-02	1216.05	356.01	12-70.01
54	607-01	318-01	9100.04	381.01	4725.01	46	302-01	134-01	9819.04	329.01	8792.01	40	170-01	65-02	1429.05	368.01	14-87.01
52	534-01	265-01	9248.04	391.01	5288.01	44	260-01	111-01	9955.04	341.01	1008.02	38	144-01	52-02	1620.05	384.01	17-44.01
50	467-01	224-01	9393.04	401.01	5948.01	42	222-01	92-02	1039.05	353.01	1163.02	36	120-01	41-02	1855.05	404.01	20-52.01
48	407-01	189-01	9536.04	412.01	6729.01	40	189-01	79-02	1072.05	367.01	1353.02	34	99-02	32-02	2067.05	424.01	24-62.01
46	352-01	157-01	9676.04	423.01	7658.01	38	159-01	58-02	1095.05	381.01	1586.02	32	81-02	25-02	2279.05	444.01	29-70.01
44	303-01	129-01	9814.04	435.01	8771.01	36	133-01	46-02	1128.05	396.01	1874.02	30	66-02	19-02	2491.05	464.01	36-24.01
42	259-01	106-01	9950.04	448.01	1011.02	34	110-01	36-02	1160.05	412.01	2237.02	28	52-02	14-02	2711.05	484.01	44-77.01
40	220-01	85-02	1008.05	461.01	1175.02	32	90-02	27-02	1192.05	429.01	2697.02	26	47-02	10-02	2937.05	504.01	54-10.01
38	186-01	68-02	1021.05	475.01	1376.02	30	73-02	21-02	1224.05	447.01	3289.02	24	35-02	7-02	3167.05	524.01	64-74.01
36	155-01	54-02	1034.05	490.01	1624.02	28	58-02	16-02	1256.05	468.01	4062.02	22	28-02	5-02	3407.05	544.01	76-40.01
34	129-01	42-02	1047.05	506.01	1937.02	26	46-02	11-02	1288.05	490.01	5088.02	20	22-02	4-02	3647.05	564.01	90-26.01
32	106-01	32-02	1059.05	523.01	2332.02	24	36-02	8-03	1319.05	514.01	6473.02	18	18-02	3-03	3887.05	584.01	105-13.01
30	85-02	24-02	1072.05	541.01	2841.02	22	27-02	5-03	1350.05	541.01	8386.02	16	15-02	2-03	4127.05	604.01	121-43.01
28	67-02	18-02	1084.05	561.01	3505.02	20	20-02	3-03	1381.05	572.01	1109.03	14	12-02	1-03	4367.05	624.01	138-43.01
26	54-02	13-02	1095.05	583.01	4385.02	18	15-02	2-03	1412.05	607.01	1502.03	12	10-02	0-04	4607.05	644.01	156-43.01
24	43-02	9-03	1107.05	607.01	5573.02	16	10-02	1-03	1443.05	647.01	2094.03	10	8-02	0-04	4847.05	664.01	175-43.01
22	34-02	6-03	1118.05	634.01	7211.02	14	7-03	0-04	1474.05	696.01	3023.03	8	6-02	0-04	5087.05	684.01	194-43.01
20	24-02	4-03	1129.05	665.01	9523.02	12	4-03	0-04	1505.05	755.01	4571.03	6	5-02	0-04	5327.05	704.01	213-43.01
18	17-02	3-03	1140.05	699.01	1288.03	10	2-03	0-04	1536.05	830.01	7361.03	4	4-02	0-04	5567.05	724.01	232-43.01
16	12-02	2-03	1150.05	739.01	1793.03	8	1-03	0-04	1567.05	931.01	1033.04	2	3-02	0-04	5807.05	744.01	251-43.01
14	7-02	1-03	1160.05	787.01	2585.03	6	0-03	0-04	1598.05	1082.01	1828.04	0	2-02	0-04	6047.05	764.01	270-43.01
12	2-02	0-04	1170.05	845.01	3902.03	4	0-04	0-04	1629.05	1233.01	2999.04	0	1-02	0-04	6287.05	784.01	289-43.01
10	0-02	0-04	1180.05	920.01	6274.03	2	0-04	0-04	1660.05	1384.01	4799.04	0	0-02	0-04	6527.05	804.01	308-43.01
8	0-03	0-04	1185.05	1066.01	8218.03	0	0-04	0-04	1691.05	1535.01	7499.04	0	0-02	0-04	6767.05	824.01	327-43.01
6	0-03	0-04	1194.05	1200.01	1108.04	0	0-04	0-04	1722.05	1686.01	1033.05	0	0-02	0-04	6997.05	844.01	346-43.01
4	0-03	0-04	1199.05	1387.01	1553.04	0	0-04	0-04	1753.05	1837.01	1550.05	0	0-02	0-04	7227.05	864.01	365-43.01
2	0-03	0-04	1204.05	1619.01	2289.04	0	0-04	0-04	1784.05	2088.01	2301.06	0	0-02	0-04	7457.05	884.01	384-43.01
0	0-03	0-04	1209.05	1882.01	3619.04	0	0-04	0-04	1815.05	2339.01	3051.06	0	0-02	0-04	7677.05	904.01	403-43.01
	0-03	0-04	1214.05	2199.01	5342.04	0	0-04	0-04	1846.05	2589.01	3801.06	0	0-02	0-04	7897.05	924.01	422-43.01
	0-03	0-04	1219.05	2599.01	8342.04	0	0-04	0-04	1877.05	2839.01	4551.06	0	0-02	0-04	8117.05	944.01	441-43.01
	0-03	0-04	1224.05	3099.01	12042.04	0	0-04	0-04	1908.05	3089.01	5291.06	0	0-02	0-04	8337.05	964.01	460-43.01
	0-03	0-04	1229.05	3699.01	17042.04	0	0-04	0-04	1939.05	3339.01	6031.06	0	0-02	0-04	8557.05	984.01	479-43.01
	0-03	0-04	1234.05	4399.01	23042.04	0	0-04	0-04	1970.05	3589.01	6771.06	0	0-02	0-04	8777.05	1004.01	498-43.01
	0-03	0-04	1239.05	5199.01	30042.04	0	0-04	0-04	2001.05	3839.01	7511.06	0	0-02	0-04	8997.05	1024.01	517-43.01
	0-03	0-04	1244.05	6099.01	38042.04	0	0-04	0-04	2032.05	4089.01	8251.06	0	0-02	0-04	9217.05	1044.01	536-43.01
	0-03	0-04	1249.05	7099.01	47042.04	0	0-04	0-04	2063.05	4339.01	8991.06	0	0-02	0-04	9437.05	1064.01	555-43.01
	0-03	0-04	1254.05	8299.01	57042.04	0	0-04	0-04	2094.05	4589.01	9731.06	0	0-02	0-04	9657.05	1084.01	574-43.01
	0-03	0-04	1259.05	9599.01	68042.04	0	0-04	0-04	2125.05	4839.01	10471.06	0	0-02	0-04	9877.05	1104.01	593-43.01
	0-03	0-04	1264.05	1099.01	80042.04	0	0-04	0-04	2156.05	5089.01	11211.06	0	0-02	0-04	10097.05	1124.01	612-43.01
	0-03	0-04	1269.05	1219.01	93042.04	0	0-04	0-04	2187.05	5339.01	11951.06	0	0-02	0-04	10317.05	1144.01	631-43.01
	0-03	0-04	1274.05	1359.01	107042.04	0	0-04	0-04	2218.05	5589.01	12691.06	0	0-02	0-04	10537.05	1164.01	650-43.01
	0-03	0-04	1279.05	1519.01	125042.04	0	0-04	0-04	2249.05	5839.01	13471.06	0	0-02	0-04	10757.05	1184.01	669-43.01
	0-03	0-04	1284.05	1699.01	145042.04	0	0-04	0-04	2280.05	6089.01	14251.06	0	0-02	0-04	10977.05	1204.01	688-43.01
	0-03	0-04	1289.05	1899.01	167042.04	0	0-04	0-04	2311.05	6339.01	15031.06	0	0-02	0-04	11197.05	1224.01	707-43.01
	0-03	0-04	1294.05	2119.01	191042.04	0	0-04	0-04	2342.05	6589.01	15811.06	0	0-02	0-04	11417.05	1244.01	726-43.01
	0-03	0-04	1299.05	2359.01	217042.04	0	0-04	0-04	2373.05	6839.01	16591.06	0	0-02	0-04	11637.05	1264.01	745-43.01
	0-03	0-04	1304.05	2619.01	245042.04	0	0-04	0-04	2404.05	7089.01	17371.06	0	0-02	0-04	11857.05	1284.01	764-43.01
	0-03	0-04	1309.05	2899.01	275042.04	0	0-04	0-04	2435.05	7339.01	18151.06	0	0-02	0-04	12077.05	1304.01	783-43.01
	0-03	0-04	1314.05	3199.01	307042.04	0	0-04	0-04	2466.05	7589.01	18931.06	0	0-02	0-04	12297.05	1324.01	802-43.01
	0-03	0-04	1319.05	3519.01	341042.04	0	0-04	0-04	2497.05	7839.01	19711.06	0	0-02	0-04	12517.05	1344.01	821-43.01
	0-03	0-04	1324.05	3859.01	377042.04	0	0-04	0-04	2528.05	8089.01	20491.06	0	0-02	0-04	12737.05	1364.01	840-43.01
	0-03	0-04	1329.05	4219.01	415042.04	0	0-04	0-04	2559.05	8339.01	21271.06	0	0-02	0-04	12957.05	1384.01	859-43.01
	0-03	0-04	1334.05	4599.01	455042.04	0	0-04	0-04	2590.05	8589.01	22051.06	0	0-02	0-04	13177.05	1404.01	878-43.01
	0-03	0-04	1339.05	5019.01	497042.04	0	0-04	0-04	2621.05	8839.01	22831.06	0	0-02	0-04	13397.05	1424.01	897-43.01
	0-03	0-04	1344.05	5459.01	541042.04	0	0-04	0-04	2652.05	9089.01	23611.06	0	0-02	0-04	13617.05	1444.01	916-43.01
	0-03	0-04	1349.05	5919.01	587042.04	0	0-04	0-04	2683.05	9339.01	24391.06	0	0-02	0-04	13837.05	1464.01	935-43.01
	0-03	0-04	1354.05	6399.01	635042.04	0	0-04	0-04	2714.05	9589.01	25171.06	0	0-02	0-04	14057.05	1484.01	954-43.01
	0-03	0-04	1359.05	6899.01	685042.04	0	0-04	0-04	2745.05	9839.01	25951.06	0	0-02	0-04	14277.05	1504.01	973-43.01
	0-03	0-04	1364.05	7419.01	737042.04	0	0-04	0-04									

$l = 0.1$ $l = 1.0$ $l = 10$

T	P	P	T	A	M	u' (ft/sec.)	P	T	P	M	A	T	P	u' (ft/sec.)	M	A
49.5	1.18 00	368-01	9085 04	219 01	2265 01	575	415-01	217-01	1593 05	300 01	6691 01	495	147-01	1703 05	502 01	1710 01
50	1.18 00	384-01	9222 04	232 01	2390 01	56	381-01	194-01	1104 05	314 01	7220 01	48	137-01	1715 05	479 01	1687 01
51	1.18 00	781-01	9451 04	241 01	2576 01	54	358-01	166-01	1119 05	324 01	8022 01	46	115-01	1727 05	451 01	1651 01
52	1.14 00	687-01	9576 04	249 01	2785 01	52	339-01	141-01	1134 05	334 01	8956 01	44	99-02	1740 05	424 01	1616 01
53	1.10 00	503-01	9749 04	257 01	3023 01	50	283-01	120-01	1148 05	345 01	1005 02	42	85-02	1753 05	418 01	1582 01
54	99-01	526-01	9918 04	266 01	3293 01	48	230-01	101-01	1163 05	357 01	1153 02	40	72-02	1765 05	412 01	1547 01
55	83-01	457-01	1009 05	275 01	3602 01	46	201-01	84-02	1177 05	369 01	1285 02	38	51-02	1778 05	406 01	1512 01
56	73-01	396-01	1025 05	284 01	3957 01	44	174-01	69-02	1191 05	381 01	1456 02	36	44-02	1791 05	400 01	1477 01
57	64-01	341-01	1041 05	294 01	4365 01	42	150-01	57-02	1204 05	395 01	1684 02	34	34-02	1804 05	394 01	1442 01
58	56-01	291-01	1057 05	304 01	4837 01	40	128-01	465-02	1218 05	409 01	1948 02	32	28-02	1816 05	388 01	1407 01
59	50-01	245-01	1072 05	315 01	5387 01	38	109-01	375-02	1231 05	424 01	2271 02	30	22-02	1828 05	382 01	1372 01
60	45-01	213-01	1087 05	326 01	6030 01	36	914-02	296-02	1244 05	440 01	2670 02	28	17-02	1840 05	376 01	1337 01
61	40-01	176-01	1103 05	337 01	6787 01	34	762-02	236-02	1257 05	457 01	3168 02	26	13-02	1851 05	370 01	1302 01
62	35-01	147-01	1117 05	349 01	7686 01	32	659-02	183-02	1270 05	475 01	3798 02	24	10-02	1864 05	364 01	1267 01
63	30-01	121-01	1132 05	362 01	8759 01	30	514-02	140-02	1282 05	496 01	4606 02	22	7-02	1876 05	358 01	1232 01
64	25-01	994-02	1146 05	376 01	1005 02	28	415-02	106-02	1295 05	517 01	5655 02	20	5-02	1888 05	352 01	1197 01
65	20-01	807-02	1161 05	390 01	1162 02	26	330-02	780-03	1307 05	541 01	7043 02	18	4-02	1899 05	346 01	1162 01
66	15-01	648-02	1175 05	405 01	1355 02	24	258-02	564-03	1319 05	568 01	8912 02	16	3-02	1910 05	340 01	1127 01
67	10-01	514-02	1188 05	422 01	1593 02	22	199-02	397-03	1331 05	598 01	1149 03	14	2-02	1921 05	334 01	1092 01
68	5-01	403-02	1202 05	439 01	1893 02	20	150-02	272-03	1343 05	631 01	1512 03	12	1-02	1932 05	328 01	1057 01
69	0-01	311-02	1215 05	459 01	2273 02	18	110-02	180-03	1355 05	669 01	2038 03	10	0-02	1943 05	322 01	1022 01
70	54-02	236-02	1228 05	479 01	2764 02	16	786-03	114-03	1364 05	713 01	2830 03	08	84-04	1954 05	316 01	987 01
71	48-02	176-02	1241 05	502 01	3406 02	14	582-03	690-04	1375 05	766 01	4070 03	06	64-04	1965 05	310 01	952 01
72	42-02	129-02	1254 05	527 01	4263 02	12	388-03	590-04	1385 05	820 01	5123 03	04	50-04	1976 05	304 01	917 01
73	36-02	915-03	1266 05	555 01	5431 02	10	222-03	502-04	1395 05	912 01	6784 03	02	38-04	1987 05	298 01	882 01
74	30-02	635-03	1279 05	587 01	7061 02	08	170-03	439-04	1405 05	962 01	9275 03	00	27-04	1998 05	292 01	847 01
75	24-02	425-03	1290 05	623 01	9399 02	06	126-03	320-05	1405 05	102 02	1708 04	04	11-05	2009 05	286 01	812 01
76	18-02	275-03	1302 05	665 01	1287 03	04	908-04	578-05	1410 05	109 02	2371 04	02	0-05	2020 05	280 01	777 01
77	12-02	157-03	1313 05	714 01	1824 03	02	622-04	339-05	1415 05	118 02	3451 04	00	19-05	2031 05	274 01	742 01
78	6-02	962-04	1324 05	775 01	2700 03	00	388-04	181-05	1419 05	130 02	5372 04	00	14-05	2042 05	268 01	707 01
79	0-02	606-04	1335 05	852 01	4239 03	04	338-04	939-06	1424 05	146 02	9221 04	00	11-05	2053 05	262 01	672 01
80	44-03	351-04	1340 05	899 01	5473 03	02	231-04	313-06	1429 05	169 02	1852 05	00	8-05	2064 05	256 01	637 01
81	38-03	235-04	1345 05	955 01	7259 03	00	428-05	781-07	1433 05	207 02	4953 05	00	5-05	2075 05	250 01	602 01
82	32-03	125-04	1350 05	102 02	9966 03	00	790-06	750-08	1438 05	291 02	2672 06	00	349-06	2086 05	244 01	567 01
83	26-03	0-04	1355 05	111 02	1434 04	00						00				
84	20-03	685-05	1360 05	122 02	2201 04	00						00				
85	14-03	475-05	1365 05	137 02	3718 04	00						00				
86	8-03	266-05	1370 05	158 02	7315 04	00						00				
87	2-03	657-05	1375 05	195 02	1901 05	00						00				
88	0-03	219-05	1380 05	281 02	9753 05	00						00				

Table 24a. EXPANSION AFTER FREEZING - WEDGE NOZZLE
T₀' = 6000°K, P₀' = 100 atm.

$l = 0.1$
 $l = 1.0$
 $l = 10$

T	P	P	u' (ft/sec.)	M	A	T	P	P	u' (ft/sec.)	M	A	T	P	P	u' (ft/sec.)	M	A
.73	.198 00	.138 00	.8502 04	.198 01	.1803 01	.63	.780 00	.455-01	.1010 05	.261 01	.3851 01	.55	.300-01	.149-01	.1130 05	.319 01	.8957 01
.72	.190 00	.130 00	.8500 04	.209 01	.1861 01	.62	.741-01	.425-01	.1019 05	.272 01	.4021 01	.54	.282-01	.137-01	.1137 05	.331 01	.8455 01
.70	.174 00	.116 00	.8792 04	.217 01	.1987 01	.60	.687-01	.371-01	.1035 05	.281 01	.4398 01	.52	.249-01	.117-01	.1151 05	.341 01	.1087 02
.68	.159 00	.103 00	.8980 04	.225 01	.2129 01	.58	.598-01	.321-01	.1051 05	.290 01	.4828 01	.50	.219-01	.097-01	.1165 05	.352 01	.1188 02
.66	.145 00	.0910-01	.9164 04	.233 01	.2290 01	.56	.535-01	.277-01	.1066 05	.300 01	.5322 01	.48	.192-01	.089-01	.1179 05	.364 01	.1382 02
.64	.131 00	.802-01	.9344 04	.241 01	.2471 01	.54	.476-01	.238-01	.1082 05	.309 01	.5893 01	.46	.167-01	.081-01	.1193 05	.376 01	.1525 02
.62	.119 00	.704-01	.9521 04	.250 01	.2676 01	.52	.422-01	.203-01	.1097 05	.320 01	.6555 01	.44	.144-01	.072-01	.1207 05	.389 01	.1783 02
.60	.108 00	.615-01	.9695 04	.258 01	.2910 01	.50	.373-01	.173-01	.1112 05	.330 01	.7328 01	.42	.124-01	.060-01	.1220 05	.402 01	.2036 02
.58	.969-01	.535-01	.9865 04	.267 01	.3177 01	.48	.327-01	.145-01	.1127 05	.342 01	.8234 01	.40	.106-01	.051-01	.1234 05	.416 01	.2375 02
.56	.876-01	.464-01	.1003 05	.277 01	.3482 01	.46	.286-01	.122-01	.1141 05	.353 01	.9305 01	.38	.0896-02	.043-02	.1247 05	.432 01	.2716 02
.54	.776-01	.400-01	.1020 05	.286 01	.3834 01	.44	.248-01	.101-01	.1156 05	.366 01	.1058 02	.36	.0753-02	.034-02	.1260 05	.448 01	.3230 02
.52	.691-01	.343-01	.1036 05	.296 01	.4240 01	.42	.214-01	.083-02	.1170 05	.379 01	.1210 02	.34	.0627-02	.027-02	.1272 05	.465 01	.3806 02
.50	.612-01	.292-01	.1052 05	.307 01	.4713 01	.40	.184-01	.0681-02	.1184 05	.393 01	.1395 02	.32	.0517-02	.021-02	.1285 05	.484 01	.4472 02
.48	.540-01	.247-01	.1067 05	.318 01	.5265 01	.38	.157-01	.0551-02	.1197 05	.407 01	.1619 02	.30	.0421-02	.014-02	.1297 05	.504 01	.5557 02
.46	.474-01	.208-01	.1083 05	.329 01	.5915 01	.36	.132-01	.0441-02	.1211 05	.423 01	.1896 02	.28	.0339-02	.011-02	.1310 05	.526 01	.6840 02
.44	.414-01	.173-01	.1098 05	.341 01	.6684 01	.34	.111-01	.0349-02	.1224 05	.440 01	.2240 02	.26	.0269-02	.009-02	.1322 05	.551 01	.8260 02
.42	.359-01	.144-01	.1113 05	.354 01	.7602 01	.32	.0918-02	.0272-02	.1237 05	.458 01	.2673 02	.24	.0210-02	.004-02	.1333 05	.578 01	.1084 03
.40	.309-01	.118-01	.1127 05	.367 01	.8705 01	.30	.0753-02	.0209-02	.1250 05	.477 01	.3226 02	.22	.0161-02	.003-02	.1345 05	.608 01	.1400 03
.38	.265-01	.099-02	.1142 05	.381 01	.1004 02	.28	.0610-02	.0159-02	.1263 05	.499 01	.3942 02	.20	.0121-02	.002-02	.1356 05	.641 01	.1869 03
.36	.225-01	.0771-02	.1156 05	.397 01	.1168 02	.26	.0487-02	.0117-02	.1275 05	.522 01	.4884 02	.18	.0087-03	.001-03	.1368 05	.680 01	.2501 03
.34	.189-01	.0613-02	.1170 05	.413 01	.1371 02	.24	.0383-02	.0085-03	.1288 05	.548 01	.6148 02	.16	.00632-03	.000-03	.1378 05	.725 01	.3484 04
.32	.158-01	.0481-02	.1184 05	.430 01	.1624 02	.22	.0296-02	.0060-03	.1300 05	.577 01	.7877 02	.14	.00434-03	.000-03	.1389 05	.778 01	.5030 04
.30	.130-01	.0372-02	.1197 05	.449 01	.1946 02	.20	.0225-02	.0041-03	.1312 05	.609 01	.1030 03	.12	.00286-03	.000-03	.1399 05	.844 01	.7507 03
.28	.106-01	.0284-02	.1211 05	.470 01	.2360 02	.18	.0166-02	.00277-03	.1323 05	.647 01	.1381 03	.10	.00177-03	.000-03	.1409 05	.926 01	.1219 04
.26	.085-02	.0212-02	.1224 05	.492 01	.2900 02	.16	.0119-02	.00177-03	.1334 05	.690 01	.1904 03	.09	.00135-03	.000-03	.1414 05	.977 01	.1593 04
.24	.068-02	.0155-02	.1237 05	.517 01	.3620 02	.14	.00830-03	.00108-03	.1345 05	.741 01	.2718 03	.08	.00100-03	.000-03	.1419 05	.104 02	.2139 04
.22	.044-02	.0111-02	.1249 05	.544 01	.4598 02	.12	.00552-03	.000613-04	.1356 05	.803 01	.4056 03	.07	.000716-04	.000-04	.1424 05	.111 02	.2978 04
.20	.024-02	.00770-03	.1262 05	.576 01	.5958 02	.10	.00346-03	.000320-04	.1366 05	.883 01	.6426 03	.06	.000489-04	.000-04	.1428 05	.120 02	.4284 04
.18	.0301-02	.00517-03	.1274 05	.611 01	.7904 02	.09	.00266-03	.000221-04	.1371 05	.931 01	.8336 03	.05	.000312-04	.000-04	.1433 05	.132 02	.6796 04
.16	.019-02	.00334-03	.1286 05	.652 01	.1078 03	.08	.00198-03	.000147-04	.1376 05	.989 01	.1111 04	.04	.000180-04	.000-04	.1437 05	.148 02	.1171 05
.14	.0134-02	.00205-03	.1297 05	.701 01	.1522 03	.07	.00143-03	.0000925-05	.1381 05	.106 02	.1534 04	.03	.0000891-05	.000-05	.1442 05	.171 02	.2363 05
.12	.008-02	.00118-03	.1308 05	.761 01	.2244 03	.06	.00086-04	.0000548-05	.1386 05	.115 02	.2220 04	.02	.0000330-05	.000-05	.1447 05	.210 02	.6368 05
.10	.00650-03	.000625-04	.1319 05	.837 01	.3508 03	.05	.00036-04	.0000294-05	.1391 05	.126 02	.3431 04	.01	.0000601-06	.000-06	.1451 05	.305 02	.3481 06
.08	.00382-03	.000291-04	.1330 05	.938 01	.4518 03	.04	.000187-04	.0000139-06	.1396 05	.141 02	.5843 04						
.07	.00278-03	.000185-04	.1345 05	.101 02	.8183 03	.03	.0000705-05	.00000131-06	.1401 05	.164 02	.1161 05						
.06	.00193-03	.000110-04	.1362 05	.109 02	.1173 04	.02	.0000134-05	.000000126-07	.1410 05	.284 02	.1610 06						
.05	.00126-03	.0000599-05	.1345 05	.120 02	.1795 04												
.04	.000745-04	.0000108-05	.1350 05	.134 02	.3020 04												
.03	.000379-04	.00000108-05	.1355 05	.156 02	.5908 04												
.02	.000146-04	.000000280-06	.1360 05	.191 02	.1524 05												
.01	.0000288-05	.000000277-07	.1365 05	.271 02	.7719 05												

Table 24b. EXPANSION AFTER FREEZING - HYPERBOLIC NOZZLE
 $T_0' = 6000^\circ K$, $p_0' = 100$ atm.

$l = 0.1$ $l = 1.0$ $l = 10$

T	P	P	u'	M	A	T	P	P	u'	M	A	T	P	P	u'	M	A
69	550-01	367-01	1020 05	273 01	4501 01	514	230-01	107-01	1141 05	336 01	1138 02	44	990-02	390-02	1217 05	386 01	2478 02
59	519-01	345-01	1027 05	280 01	4694 01	50	209-01	950-02	1151 05	348 01	1238 02	42	845-02	318-02	1230 05	408 01	2872 02
58	553-01	297-01	1043 05	289 01	5176 01	48	182-01	784-02	1155 05	360 01	1404 02	40	716-02	257-02	1243 05	423 01	3353 02
56	592-01	251-01	1059 05	299 01	5729 01	46	158-01	659-02	1179 05	372 01	1602 02	38	602-02	205-02	1256 05	438 01	3940 02
54	536-01	218-01	1075 05	309 01	6373 01	44	136-01	543-02	1183 05	384 01	1839 02	36	502-02	162-02	1269 05	454 01	4685 02
52	585-01	180-01	1090 05	319 01	7117 01	42	116-01	443-02	1206 05	398 01	2126 02	34	413-02	126-02	1281 05	472 01	5617 02
50	538-01	151-01	1105 05	330 01	7992 01	40	989-02	359-02	1220 05	412 01	2475 02	32	339-02	972-03	1294 05	491 01	6808 02
48	596-01	123-01	1120 05	341 01	9022 01	38	833-02	287-02	1233 05	427 01	2906 02	30	274-02	736-03	1306 05	511 01	8240 02
46	527-01	110-01	1134 05	352 01	1025 02	36	696-02	227-02	1246 05	443 01	3441 02	28	218-02	546-03	1318 05	534 01	1038 03
44	522-01	908-02	1148 05	365 01	1171 02	34	577-02	178-02	1259 05	460 01	4114 02	26	172-02	400-03	1330 05	558 01	1300 03
42	591-01	742-02	1163 05	378 01	1347 02	32	472-02	137-02	1271 05	479 01	4971 02	24	153-02	285-03	1342 05	586 01	1678 03
40	563-01	603-02	1177 05	392 01	1560 02	30	383-02	104-02	1284 05	499 01	6077 02	22	101-02	198-03	1353 05	616 01	2194 03
38	534-01	485-02	1190 05	406 01	1822 02	28	306-02	777-03	1296 05	521 01	7528 02	20	746-03	134-03	1364 05	650 01	2832 03
36	516-01	385-02	1204 05	422 01	2147 02	26	241-02	568-03	1308 05	545 01	9463 02	18	540-03	872-04	1375 05	689 01	3614 03
34	561-02	303-02	1217 05	439 01	2555 02	24	187-02	407-03	1320 05	572 01	1209 03	16	379-03	544-04	1386 05	735 01	5677 03
32	591-02	233-02	1230 05	457 01	3088 02	22	142-02	284-03	1332 05	602 01	1575 03	14	257-03	322-04	1396 05	788 01	8321 03
30	544-02	179-02	1243 05	477 01	3729 02	20	106-02	192-03	1343 05	635 01	2096 03	12	166-03	179-04	1406 05	854 01	1278 04
28	518-02	134-02	1256 05	498 01	4591 02	18	770-03	126-03	1354 05	674 01	2862 03	10	101-03	905-05	1416 05	938 01	2588 04
26	510-02	989-03	1268 05	521 01	5724 02	16	543-03	788-04	1365 05	718 01	4025 03	09	763-04	615-05	1426 05	105 02	3738 04
24	520-02	713-03	1281 05	547 01	7279 02	14	370-03	469-04	1376 05	771 01	5870 03	08	560-04	402-05	1436 05	112 02	5261 04
22	545-02	550-03	1293 05	576 01	9413 02	12	240-03	261-04	1386 05	836 01	8956 03	07	397-04	249-05	1430 05	122 02	7780 04
20	518-02	341-03	1305 05	609 01	1243 03	10	147-03	133-04	1396 05	918 01	1457 04	06	267-04	144-05	1439 05	133 02	1233 05
18	535-02	225-03	1316 05	646 01	1684 03	09	111-03	908-05	1401 05	968 01	1916 04	05	168-04	753-06	1439 05	150 02	2167 05
16	556-03	145-03	1327 05	689 01	2348 03	08	819-04	594-05	1406 05	103 02	2592 04	04	936-05	344-06	1449 05	173 02	4461 05
14	557-03	653-04	1338 05	740 01	3393 03	07	582-04	370-05	1410 05	110 02	3635 04	03	462-05	125-06	1453 05	214 02	1240 06
12	521-03	489-04	1349 05	802 01	5132 03	06	394-04	214-05	1415 05	119 02	5354 04	02	166-05	286-07	1453 05	303 02	7140 06
10	566-03	247-04	1359 05	881 01	8249 03	05	249-04	113-05	1420 05	131 02	8445 04	01	287-06	260-08	1457 05		
09	503-03	165-04	1364 05	929 01	1079 04	04	142-04	516-06	1424 05	147 02	1473 05						
08	515-03	115-04	1369 05	987 01	1450 04	03	692-05	188-06	1429 05	170 02	3020 05						
07	508-03	669-05	1374 05	106 02	2020 04	02	251-05	442-07	1434 05	210 02	8314 05						
06	533-04	408-05	1379 05	114 02	2953 04	01	441-06	520-08	1438 05	277 02	4710 06						
05	567-04	217-05	1384 05	126 02	4618 04												
04	570-04	100-05	1389 05	141 02	7974 04												
03	533-04	379-06	1393 05	163 02	1613 05												
02	590-05	912-07	1398 05	200 02	4362 05												
01	588-06	780-08	1403 05	289 02	2396 06												

Table 25a. EXPANSION AFTER FREEZING - WEDGE NOZZLE

 $T_0' = 6000^\circ\text{K}$, $p_0' = 300$ atm.

$l = 10$

$l = 1.0$

$l = 0.1$

T	P	P	u'	M	A	T	P	P	u'	M	A	T	P	P	u'	M	A
			(ft/sec.)						(ft/sec.)						(ft/sec.)		
.663	.110 00	.684-01	.9442 04	.239 01	.2875 01	.575	.450-01	.238-01	.1066 05	.295 01	.6221 01	.502	.200-01	.908-02	.1154 05	.343 01	.1293 02
.66	.109 00	.673-01	.9462 04	.244 01	.2904 01	.56	.412-01	.212-01	.1078 05	.306 01	.6724 01	.50	.197-01	.893-02	.1156 05	.350 01	.1303 02
.64	.082-01	.593-01	.9635 04	.252 01	.3155 01	.54	.365-01	.181-01	.1093 05	.316 01	.7490 01	.48	.172-01	.746-02	.1170 05	.362 01	.1485 02
.62	.085-01	.515-01	.9806 04	.261 01	.3440 01	.54	.365-01	.181-01	.1093 05	.316 01	.7490 01	.46	.149-01	.619-02	.1184 05	.374 01	.1695 02
.63	.0795-01	.448-01	.9973 04	.270 01	.3767 01	.52	.321-01	.134-01	.1108 05	.326 01	.8382 01	.44	.128-01	.510-02	.1197 05	.386 01	.1847 02
.58	.0711-01	.387-01	.1014 05	.279 01	.4141 01	.50	.282-01	.130-01	.1123 05	.337 01	.9429 01	.42	.109-01	.416-02	.1211 05	.400 01	.2052 02
.56	.0634-01	.333-01	.1030 05	.288 01	.4571 01	.48	.246-01	.109-01	.1137 05	.348 01	.1066 02	.40	.092-02	.336-02	.1224 05	.414 01	.2263 02
.54	.0563-01	.283-01	.1046 05	.298 01	.5070 01	.46	.214-01	.093-02	.1152 05	.360 01	.1213 02	.38	.0785-02	.269-02	.1236 05	.429 01	.2491 02
.52	.0498-01	.243-01	.1062 05	.308 01	.5649 01	.44	.184-01	.074-02	.1166 05	.372 01	.1389 02	.36	.064-02	.213-02	.1250 05	.445 01	.2740 02
.50	.0438-01	.203-01	.1077 05	.319 01	.6327 01	.42	.158-01	.061-02	.1180 05	.386 01	.1621 02	.34	.051-02	.166-02	.1263 05	.463 01	.3007 02
.48	.0384-01	.173-01	.1092 05	.330 01	.7124 01	.40	.135-01	.0495-02	.1193 05	.400 01	.1859 02	.32	.043-02	.128-02	.1276 05	.481 01	.3280 02
.46	.034-01	.144-01	.1107 05	.341 01	.8068 01	.38	.114-01	.0397-02	.1207 05	.414 01	.2175 02	.30	.039-02	.0973-03	.1288 05	.502 01	.3570 02
.44	.029-01	.120-01	.1122 05	.354 01	.9193 01	.36	.093-02	.0315-02	.1220 05	.430 01	.2567 02	.28	.0287-02	.0728-03	.1301 05	.524 01	.3864 02
.42	.0249-01	.098-02	.1136 05	.366 01	.1035 02	.34	.0791-02	.0247-02	.1233 05	.447 01	.3060 02	.26	.026-02	.0531-03	.1313 05	.548 01	.4167 02
.40	.0213-01	.080-02	.1151 05	.380 01	.1218 02	.32	.0650-02	.0191-02	.1246 05	.466 01	.3585 02	.24	.0175-02	.0380-03	.1324 05	.575 01	.4524 02
.38	.0181-01	.0645-02	.1165 05	.394 01	.1419 02	.30	.0528-02	.0146-02	.1259 05	.485 01	.4089 02	.22	.0133-02	.0263-03	.1336 05	.605 01	.4907 03
.36	.0152-01	.0514-02	.1178 05	.410 01	.1666 02	.28	.0424-02	.0109-02	.1272 05	.507 01	.4541 02	.20	.0090-03	.0179-03	.1347 05	.638 01	.5297 03
.34	.0127-01	.0403-02	.1192 05	.426 01	.1975 02	.26	.0355-02	.0081-03	.1284 05	.531 01	.5038 02	.18	.0071-03	.0117-03	.1359 05	.677 01	.5703 03
.32	.0105-01	.0314-02	.1205 05	.444 01	.2366 02	.24	.0281-02	.0057-03	.1308 05	.557 01	.5832 02	.16	.00507-03	.00733-04	.1369 05	.722 01	.6163 03
.30	.00855-02	.0241-02	.1219 05	.464 01	.2866 02	.22	.0199-02	.0033-03	.1320 05	.586 01	.6318 02	.14	.00344-03	.00436-04	.1380 05	.775 01	.6783 04
.28	.0069-02	.0181-02	.1232 05	.484 01	.3517 02	.20	.0109-02	.0023-03	.1331 05	.619 01	.6961 03	.12	.00236-03	.00243-04	.1390 05	.849 01	.7507 04
.26	.0048-02	.0138-02	.1245 05	.507 01	.4376 02	.18	.00527-03	.0016-04	.1342 05	.700 01	.7885 03	.10	.00133-03	.001841-05	.1405 05	.912 01	.8357 04
.24	.0029-02	.0067-03	.1257 05	.533 01	.5534 02	.16	.00345-03	.0008-04	.1353 05	.752 01	.8585 03	.08	.000760-04	.000550-05	.1410 05	.103 02	.9784 04
.22	.0020-02	.00467-03	.1269 05	.561 01	.7127 02	.14	.00212-03	.00057-04	.1363 05	.815 01	.9354 03	.07	.000540-04	.000342-05	.1414 05	.111 02	.9908 04
.20	.00146-02	.00318-03	.1281 05	.593 01	.9374 02	.12	.00161-03	.00037-04	.1374 05	.896 01	.1026 04	.06	.000365-04	.000198-05	.1419 05	.120 02	.9971 04
.18	.00083-02	.00209-03	.1293 05	.629 01	.1264 03	.10	.00119-03	.000287-05	.1383 05	.100 02	.1812 04	.05	.000230-04	.000104-05	.1424 05	.131 02	.9971 04
.16	.00033-02	.00130-03	.1305 05	.671 01	.1754 03	.08	.000850-04	.000158-05	.1388 05	.107 02	.2531 04	.04	.000132-04	.0000677-06	.1428 05	.147 02	.9985 05
.14	.00003-03	.00083-04	.1316 05	.721 01	.2521 03	.06	.000577-04	.00009318-05	.1393 05	.116 02	.3711 04	.03	.0000639-05	.0000175-06	.1433 05	.170 02	.9985 05
.12	.00000-03	.000593-04	.1327 05	.783 01	.3792 03	.04	.000211-04	.0000774-06	.1398 05	.128 02	.5825 04	.02	.0000221-05	.00000415-07	.1437 05	.209 02	.9985 05
.10	.00000-03	.00034-04	.1337 05	.860 01	.6057 03	.03	.000103-04	.0000284-06	.1402 05	.143 02	.1010 05	.01	.00000405-06	.00000260-08	.1442 05	.330 02	.9985 05
.09	.00000-03	.000282-04	.1342 05	.907 01	.7895 03	.02	.0000378-05	.00000703-07	.1412 05	.166 02	.2054 05						
.07	.00000-03	.000238-04	.1342 05	.907 01	.7895 03	.01	.00000677-06	.00000520-08	.1416 05	.304 02	.3114 06						
.06	.00000-03	.000198-05	.1347 05	.964 01	.1057 04												
.05	.00000-03	.000157-04	.1352 05	.103 02	.1468 04												
.04	.00000-03	.000103-03	.1357 05	.112 02	.2137 04												
.03	.00000-03	.0000580-05	.1362 05	.123 02	.3326 04												
.02	.00000-03	.0000309-05	.1367 05	.138 02	.5711 04												
.01	.00000-03	.0000144-05	.1372 05	.160 02	.1147 05												
.00	.00000-03	.00000534-06	.1377 05	.196 02	.3071 05												
.00	.00000-03	.00000133-06	.1382 05	.274 02	.1658 06												

Table 25b. EXPANSION AFTER FREEZING - HYPERBOLIC NOZZLE

$T_0' = 6000^\circ\text{K}$, $P_0' = 300$ atm.

$l = 0.1$ $l = 1.0$ $l = 10$

T	P	P	u'	M	A	T	P	P	u'	M	A	T	P	P	u'	M	A	T	P	P	u'	M	A
.523	.375-01	.184-01	.1073	.329	.7388	.1451	.165-01	.690-02	.1152	.359	.1542	.394	.920-02	.335-02	.1190	.347	.2658	.394	.920-02	.335-02	.1190	.347	.2658
.52	.345-01	.185-01	.1082	.317	.7066	.144	.152-01	.618-02	.1160	.370	.1667	.38	.814-02	.286-02	.1198	.415	.2681	.38	.814-02	.286-02	.1198	.415	.2681
.51	.291-01	.139-01	.1037	.328	.8983	.142	.129-01	.504-02	.1174	.383	.1930	.36	.678-02	.226-02	.1222	.431	.271	.36	.678-02	.226-02	.1222	.431	.271
.48	.250-01	.116-01	.1117	.339	.1019	.40	.110-01	.406-02	.1188	.397	.2252	.34	.559-02	.176-02	.1235	.448	.274	.34	.559-02	.176-02	.1235	.448	.274
.46	.224-01	.961-02	.1127	.350	.1162	.38	.921-02	.325-02	.1201	.412	.2649	.32	.456-02	.135-02	.1248	.466	.274	.32	.456-02	.135-02	.1248	.466	.274
.44	.193-01	.791-02	.1141	.363	.1334	.36	.768-02	.250-02	.1214	.428	.3144	.30	.366-02	.107-02	.1260	.486	.274	.30	.366-02	.107-02	.1260	.486	.274
.42	.164-01	.645-02	.1155	.376	.1543	.32	.634-02	.200-02	.1228	.445	.3769	.28	.293-02	.759-03	.1273	.508	.274	.28	.293-02	.759-03	.1273	.508	.274
.40	.145-01	.521-02	.1169	.390	.1797	.32	.518-02	.166-02	.1241	.463	.4566	.26	.230-02	.553-03	.1285	.532	.274	.26	.230-02	.553-03	.1285	.532	.274
.38	.127-01	.416-02	.1183	.404	.2110	.28	.433-02	.133-02	.1253	.483	.5599	.22	.178-02	.394-03	.1297	.558	.274	.22	.178-02	.394-03	.1297	.558	.274
.36	.980-02	.329-02	.1197	.420	.2501	.28	.361-02	.106-02	.1266	.504	.6957	.22	.135-02	.274-03	.1309	.587	.274	.22	.135-02	.274-03	.1309	.587	.274
.34	.810-02	.257-02	.1210	.437	.2992	.26	.261-02	.849-03	.1278	.524	.8774	.20	.937-03	.184-03	.1321	.620	.274	.20	.937-03	.184-03	.1321	.620	.274
.32	.662-02	.198-02	.1223	.455	.3618	.24	.202-02	.640-03	.1290	.554	.1125	.18	.720-03	.120-03	.1343	.658	.274	.18	.720-03	.120-03	.1343	.658	.274
.30	.536-02	.150-02	.1236	.474	.4428	.22	.153-02	.512-03	.1302	.563	.1471	.16	.505-03	.747-04	.1354	.701	.274	.16	.505-03	.747-04	.1354	.701	.274
.28	.428-02	.112-02	.1249	.496	.5491	.20	.113-02	.437-03	.1314	.581	.2656	.12	.341-03	.442-04	.1374	.816	.274	.12	.341-03	.442-04	.1374	.816	.274
.26	.334-02	.816-03	.1261	.519	.6911	.18	.820-03	.337-03	.1325	.593	.3811	.10	.230-03	.244-04	.1384	.845	.274	.10	.230-03	.244-04	.1384	.845	.274
.24	.260-02	.531-03	.1274	.545	.8845	.16	.576-03	.280-04	.1336	.597	.5568	.08	.161-03	.123-04	.1393	.884	.274	.08	.161-03	.123-04	.1393	.884	.274
.22	.195-02	.406-03	.1286	.573	.1154	.14	.389-03	.200-04	.1347	.611	.8585	.06	.101-03	.546-05	.1384	.108	.274	.06	.101-03	.546-05	.1384	.108	.274
.20	.137-02	.274-03	.1298	.606	.1538	.12	.252-03	.142-04	.1358	.628	.1404	.04	.07	.338-05	.1393	.128	.274	.04	.07	.338-05	.1393	.128	.274
.18	.105-02	.179-03	.1309	.643	.2104	.10	.153-03	.961-05	.1373	.639	.1853	.02	.351-04	.195-05	.1398	.143	.274	.02	.351-04	.195-05	.1398	.143	.274
.16	.749-03	.112-03	.1320	.686	.2965	.08	.115-03	.627-05	.1378	.658	.2516	.06	.221-04	.102-05	.1403	.166	.274	.06	.221-04	.102-05	.1403	.166	.274
.14	.503-03	.664-04	.1331	.735	.4336	.06	.846-04	.389-05	.1382	.672	.3543	.04	.125-04	.462-06	.1403	.203	.274	.04	.125-04	.462-06	.1403	.203	.274
.12	.329-03	.368-04	.1342	.798	.6640	.04	.599-04	.224-05	.1387	.688	.5243	.02	.063-05	.167-06	.1403	.250	.274	.02	.063-05	.167-06	.1403	.250	.274
.10	.200-03	.187-04	.1352	.877	.1082	.02	.403-04	.116-05	.1392	.705	.8313	.00	.215-05	.405-07	.1412	.310	.274	.00	.215-05	.405-07	.1412	.310	.274
.08	.152-03	.127-04	.1367	.925	.1426	.00	.253-04	.537-06	.1397	.722	.1459	.00	.370-06	.270-08	.1417	.360	.274	.00	.370-06	.270-08	.1417	.360	.274
.06	.111-03	.932-05	.1382	.983	.1932	.04	.144-04	.592-06	.1401	.745	.3015	.00						.00					
.04	.750-04	.513-05	.1397	.105	.2715	.04	.094-05	.192-06	.1406	.765	.8997	.00						.00					
.02	.336-04	.299-05	.1372	.114	.4008	.02	.246-05	.459-07	.1406	.783	.8997	.00						.00					
.00	.191-04	.713-06	.1381	.125	.6339	.00	.429-06	.270-08	.1410	.802	.4847	.00						.00					
.00	.927-05	.259-06	.1386	.163	.2283	.00												.00					
.00	.623-05	.621-07	.1391	.200	.5324	.00												.00					
.00	.581-05	.540-08	.1395	.284	.3616	.00												.00					

Table 26a. EXPANSION AFTER FREEZING - WEDGE NOZZLE

 $T_o' = 6000^\circ K, P_o' = 1000 \text{ atm.}$

$l = 0.1$ $l = 1.0$ $l = 10$

T	P	P	u'	M	A	T	P	P	u'	M	A	T	P	P	u'	M	A
(ft/sec.)																	
.591	.643-01	.357-01	.1006 05	.274 01	.532 01	.513	.302-01	.144-01	.1094 05	.320 01	.8874 01	.441	.149-01	.609-02	.1161 05	.366 01	.1695 02
.58	.603-01	.329-01	.1015 05	.280 01	.4789 01	.50	.278-01	.129-01	.1103 05	.339 01	.9576 01	.44	.148-01	.603-02	.1162 05	.371 01	.1707 02
.56	.523-01	.282-01	.1031 05	.289 01	.5312 01	.48	.241-01	.108-01	.1118 05	.351 01	.1086 02	.42	.126-01	.591-02	.1165 05	.384 01	.1977 02
.54	.475-01	.240-01	.1047 05	.299 01	.5920 01	.46	.209-01	.089-02	.1133 05	.353 01	.1240 02	.40	.107-01	.596-02	.1186 05	.398 01	.2307 02
.52	.416-01	.204-01	.1062 05	.309 01	.6629 01	.44	.179-01	.073-02	.1147 05	.365 01	.1424 02	.38	.098-02	.516-02	.1203 05	.413 01	.2715 02
.51	.365-01	.171-01	.1078 05	.320 01	.7463 01	.42	.153-01	.059-02	.1161 05	.378 01	.1648 02	.36	.0748-02	.449-02	.1216 05	.429 01	.3223 02
.48	.318-01	.143-01	.1093 05	.331 01	.8448 01	.40	.130-01	.048-02	.1175 05	.392 01	.1920 02	.34	.0617-02	.394-02	.1233 05	.446 01	.3864 02
.46	.275-01	.119-01	.1108 05	.343 01	.9621 01	.38	.109-01	.0387-02	.1188 05	.407 01	.2256 02	.32	.0504-02	.369-02	.1242 05	.464 01	.3482 02
.44	.237-01	.091-02	.1122 05	.355 01	.1103 02	.36	.0912-02	.0308-02	.1202 05	.422 01	.2675 02	.30	.0407-02	.341-02	.1255 05	.484 01	.3147 02
.42	.203-01	.061-02	.1137 05	.368 01	.1273 02	.34	.073-02	.0230-02	.1215 05	.439 01	.3202 02	.28	.0324-02	.313-03	.1267 05	.505 01	.2832 02
.41	.175-01	.046-02	.1151 05	.382 01	.1480 02	.32	.0616-02	.0184-02	.1228 05	.457 01	.3875 02	.26	.0254-02	.281-03	.1280 05	.524 01	.2502 02
.38	.143-01	.0319-02	.1165 05	.396 01	.1734 02	.30	.0498-02	.0139-02	.1241 05	.477 01	.4745 02	.24	.0197-02	.243-03	.1292 05	.555 01	.2155 03
.36	.121-01	.0211-02	.1179 05	.412 01	.2051 02	.28	.0397-02	.0104-02	.1254 05	.498 01	.5888 02	.22	.0149-02	.203-03	.1304 05	.584 01	.1510 03
.34	.100-01	.021-02	.1192 05	.428 01	.2449 02	.26	.0312-02	.0075-03	.1266 05	.522 01	.7415 02	.20	.0110-02	.205-03	.1316 05	.617 01	.2016 03
.32	.0823-02	.0248-02	.1206 05	.446 01	.2955 02	.24	.0241-02	.00540-03	.1279 05	.547 01	.9497 02	.18	.00798-03	.133-03	.1327 05	.654 01	.2769 03
.31	.0673-02	.0166-02	.1219 05	.465 01	.3608 02	.22	.0183-02	.00375-03	.1291 05	.576 01	.1240 03	.16	.00560-03	.130-04	.1338 05	.698 01	.3215 03
.28	.0533-02	.0140-02	.1232 05	.486 01	.4465 02	.20	.0136-02	.00253-03	.1302 05	.609 01	.1654 03	.14	.00379-03	.127-04	.1343 05	.743 01	.3742 03
.26	.0423-02	.0103-02	.1245 05	.509 01	.5606 02	.18	.0086-03	.00165-03	.1314 05	.646 01	.2264 03	.12	.00244-03	.137-04	.1352 05	.812 01	.4425 03
.24	.0326-02	.00736-03	.1257 05	.535 01	.7157 02	.16	.00693-03	.00103-03	.1325 05	.689 01	.3194 03	.10	.00148-03	.137-04	.1369 05	.892 01	.5444 04
.22	.0245-02	.00513-03	.1269 05	.563 01	.9311 02	.14	.00469-03	.00061-04	.1336 05	.740 01	.4674 03	.09	.00112-03	.137-04	.1374 05	.941 01	.6965 04
.20	.0185-02	.00348-03	.1281 05	.595 01	.1238 03	.12	.00304-03	.00034-04	.1347 05	.802 01	.7165 03	.08	.000821-04	.137-04	.1384 05	.999 01	.8646 04
.18	.0134-02	.00227-03	.1293 05	.631 01	.1688 03	.10	.00185-03	.00017-04	.1357 05	.882 01	.1169 04	.07	.000581-04	.137-04	.1388 05	.107 02	.1067 04
.16	.00848-03	.00143-03	.1304 05	.674 01	.2372 03	.09	.00116-03	.00007-05	.1367 05	.939 01	.1541 04	.06	.000391-04	.137-04	.1393 05	.116 02	.1397 04
.14	.00643-03	.00093-04	.1315 05	.724 01	.3457 03	.08	.000727-04	.0000474-05	.1372 05	.987 01	.2089 04	.05	.000246-04	.137-04	.1398 05	.127 02	.1860 04
.12	.00419-03	.000473-04	.1326 05	.785 01	.5275 03	.07	.000491-04	.0000274-05	.1376 05	.106 02	.2938 04	.04	.000140-04	.137-04	.1398 05	.143 02	.2503 05
.11	.00236-03	.000241-04	.1337 05	.863 01	.8563 03	.06	.000309-04	.0000144-05	.1381 05	.114 02	.4341 04	.03	.0000673-05	.137-04	.1402 05	.165 02	.3106 05
.09	.00164-03	.000164-04	.1342 05	.910 01	.1126 04	.05	.000176-04	.00000654-06	.1386 05	.126 02	.6870 04	.02	.0000241-05	.137-04	.1407 05	.204 02	.4661 05
.08	.00113-03	.000106-04	.1347 05	.967 01	.1522 04	.04	.0000890-05	.00000238-06	.1391 05	.141 02	.1203 05	.01	.0000415-06	.137-04	.1412 05	.324 02	.5908 06
.07	.0007102-03	.0000669-05	.1352 05	.104 02	.2134 04	.03	.0000306-05	.000000567-07	.1395 05	.163 02	.2480 05						
.06	.0004688-04	.0000205-05	.1356 05	.112 02	.3141 04	.02	.00002306-05	.0000000540-08	.1400 05	.201 02	.6879 05						
.05	.0002683-04	.00001205-05	.1361 05	.123 02	.4452 04	.01	.00001530-06			.277 02	.3953 06						
.04	.0001640-04	.00000635-06	.1366 05	.138 02	.8634 04												
.03	.0001104-04	.00000340-06	.1371 05	.160 02	.1769 05												
.02	.0000438-05	.000000810-07	.1376 05	.197 02	.4867 05												
.01	.00000771-06	.0000000810-08	.1380 05	.271 02	.2756 06												

Table 26b. EXPANSION AFTER FREEZING - HYPERBOLIC NOZZLE

 $T_0' = 6000^\circ K, p_0' = 1000 \text{ atm.}$

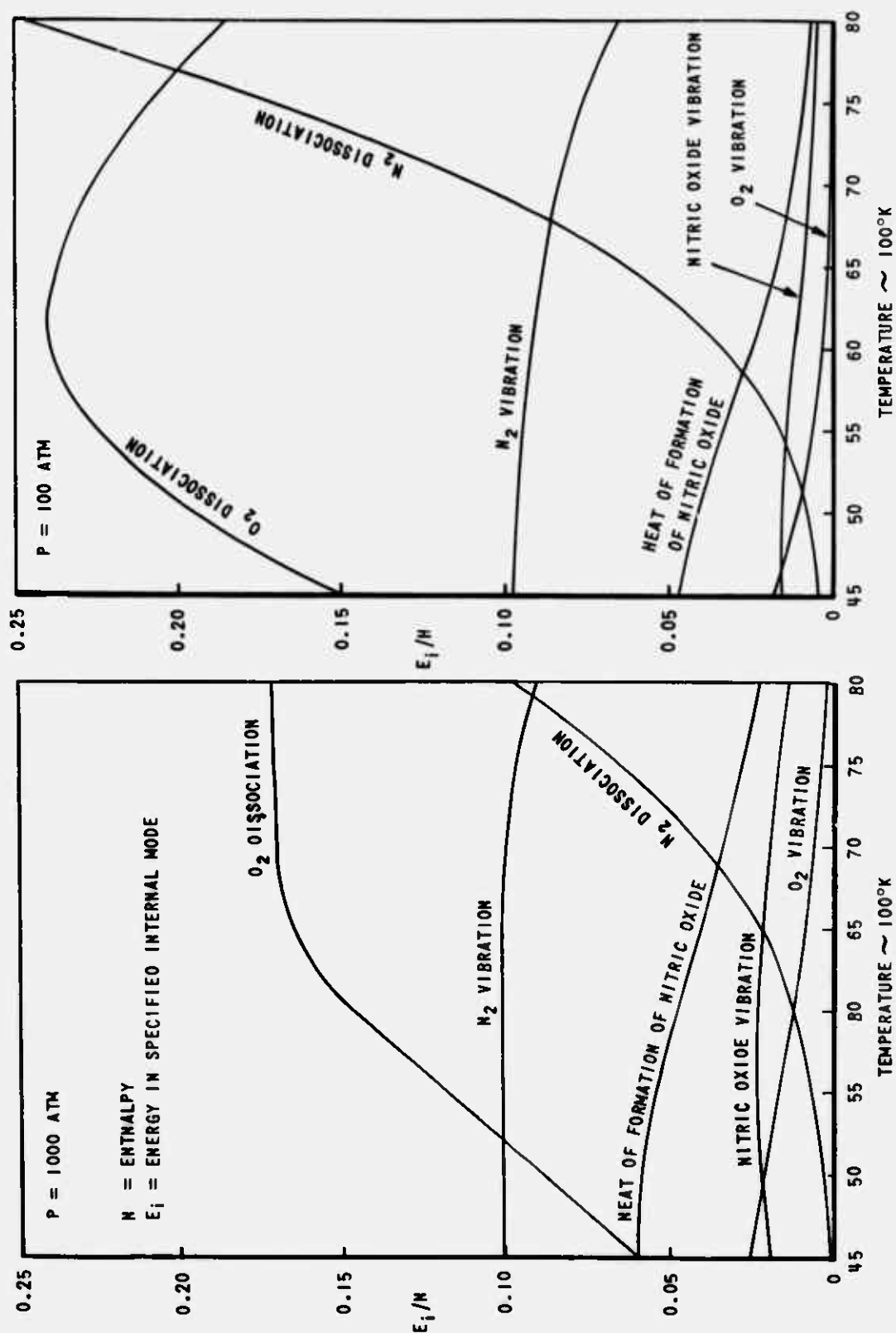
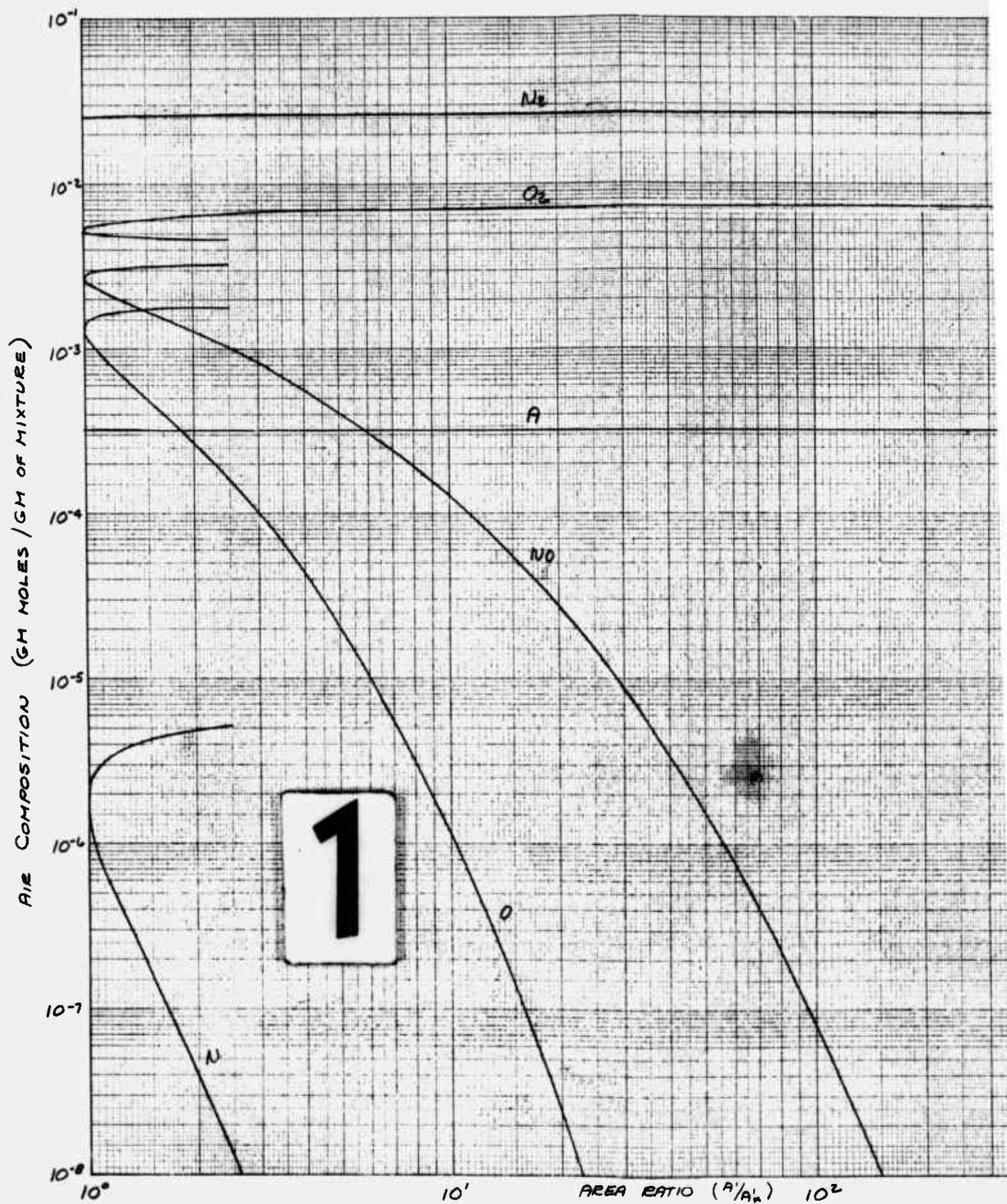


Figure 1 FRACTION OF THE ENTHALPY OF EQUILIBRIUM AIR WHICH IS IN THE SPECIFIED ENERGY MODE
(REFERENCE: CAL REPORT NO. BE-1007-A-3)



AIR COMPOSITION VS AREA RATIO IN AN ISENTROPIC

FIGURE NO. 2
 $T_0 = 4000^\circ K$
 $P_0 = 100 \text{ ATM}$

10

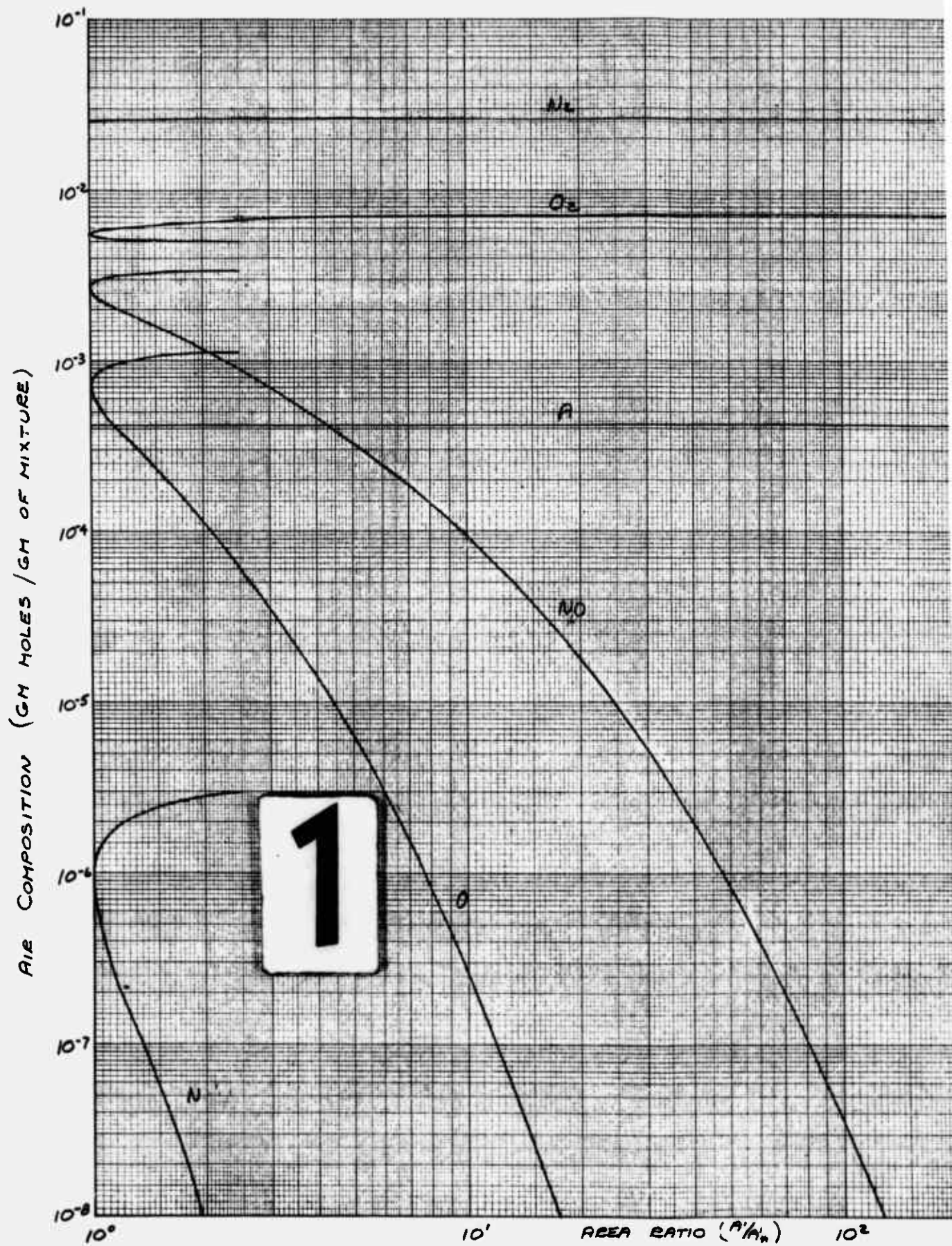
2

AREA RATIO (A'/A_*) 10^2

10^3

10^4

VS AREA RATIO IN AN ISENTROPIC EQUILIBRIUM EXPANSION



AIR COMPOSITION VS AREA RATIO IN AN

FIGURE NO. 3
 $T_0 = 4000^\circ\text{K}$
 $P_0 = 300\text{ ATM}$

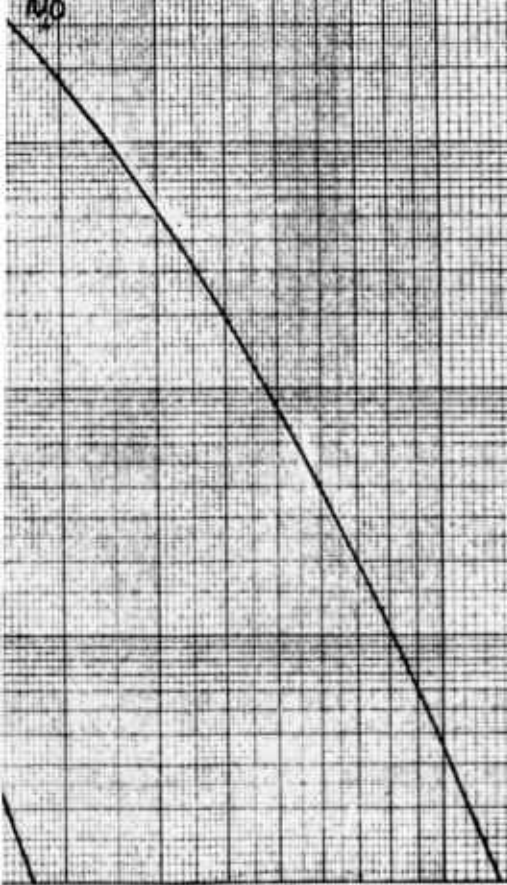
2

U_c

O_2

A

NO

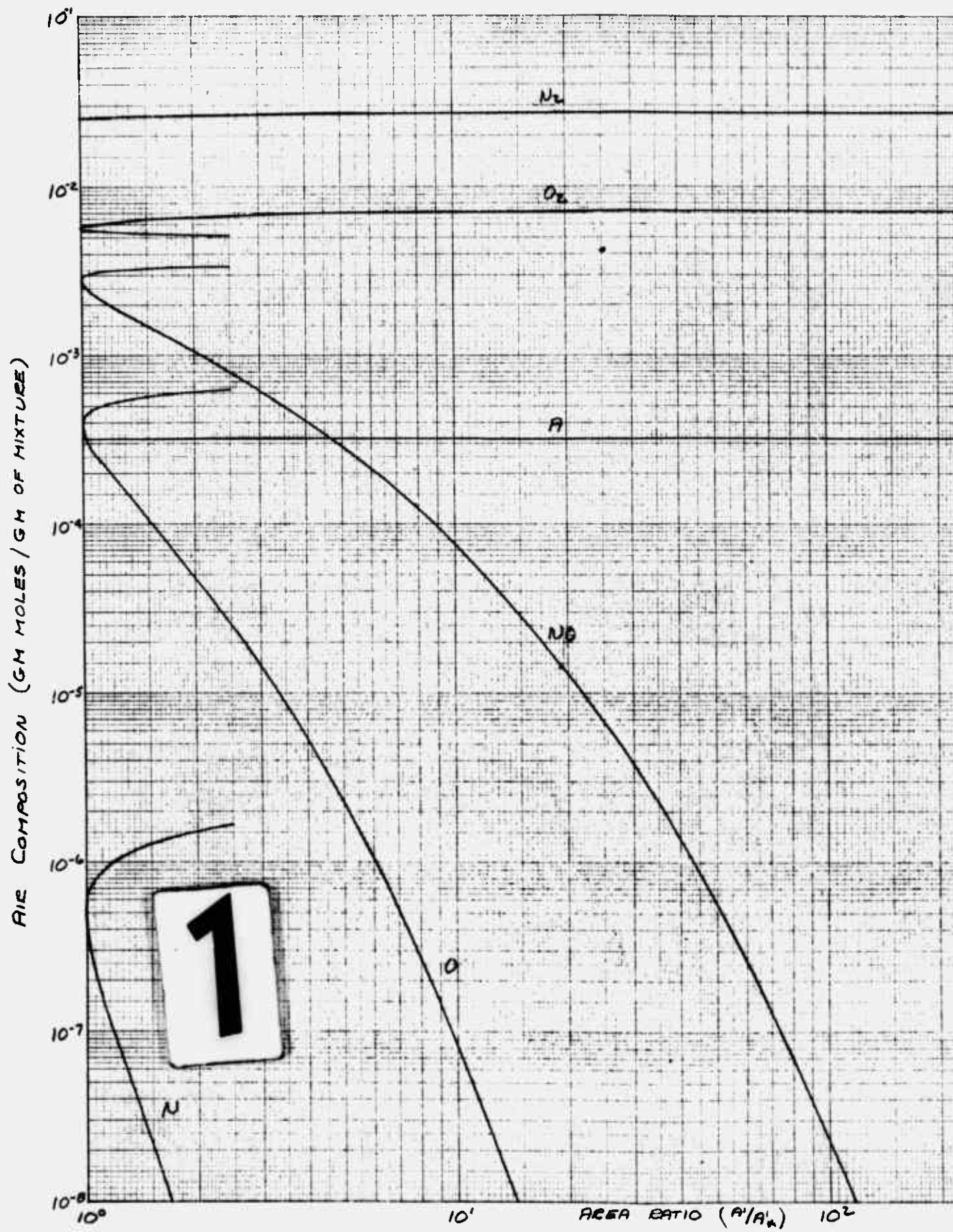


AREA RATIO (A/A_0) 10^2

10^3

10^4

ION VS AREA RATIO IN AN ISENTROPIC EQUILIBRIUM EXPANSION

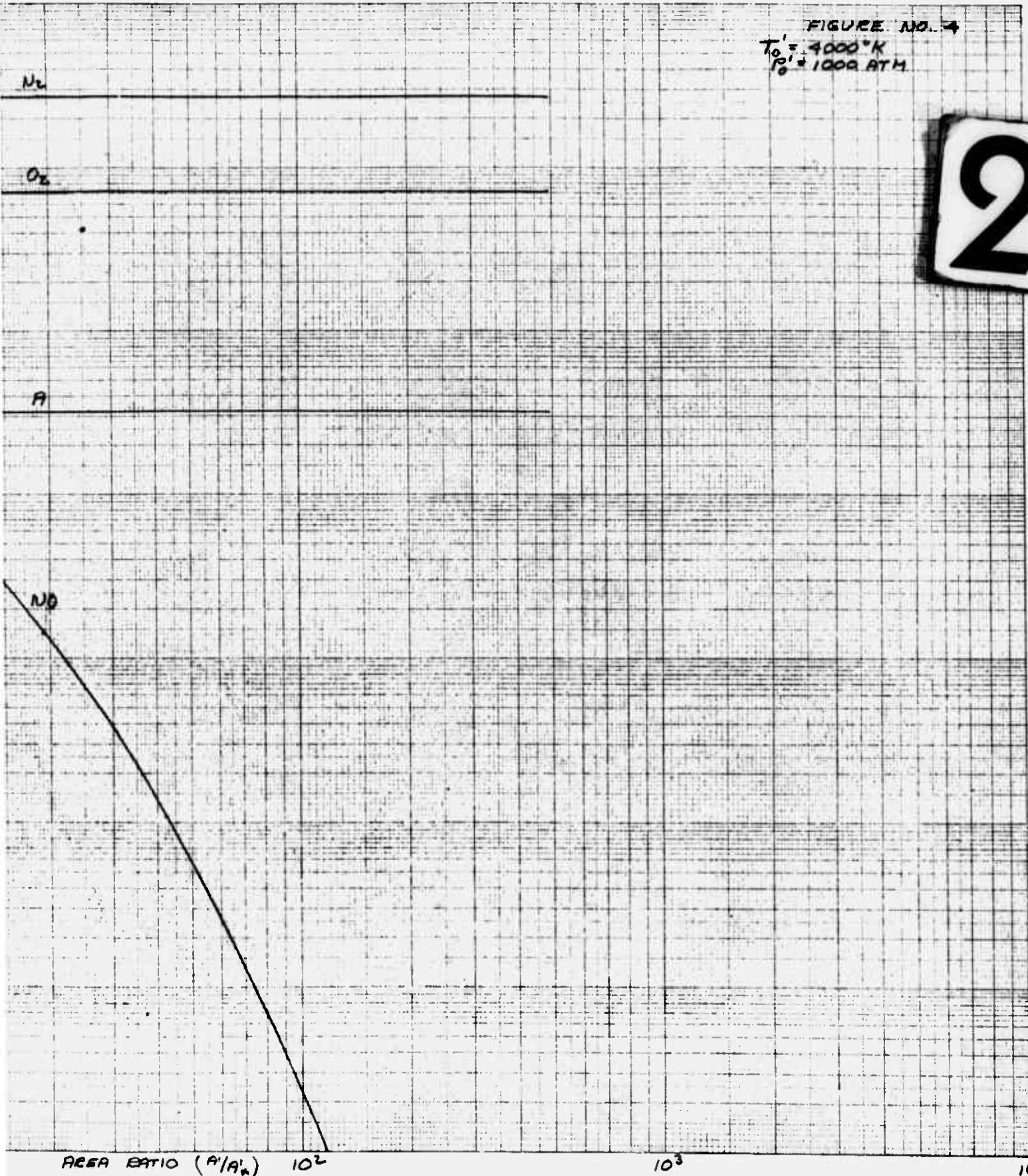


AIR COMPOSITION VS AREA RATIO IN AN J.S

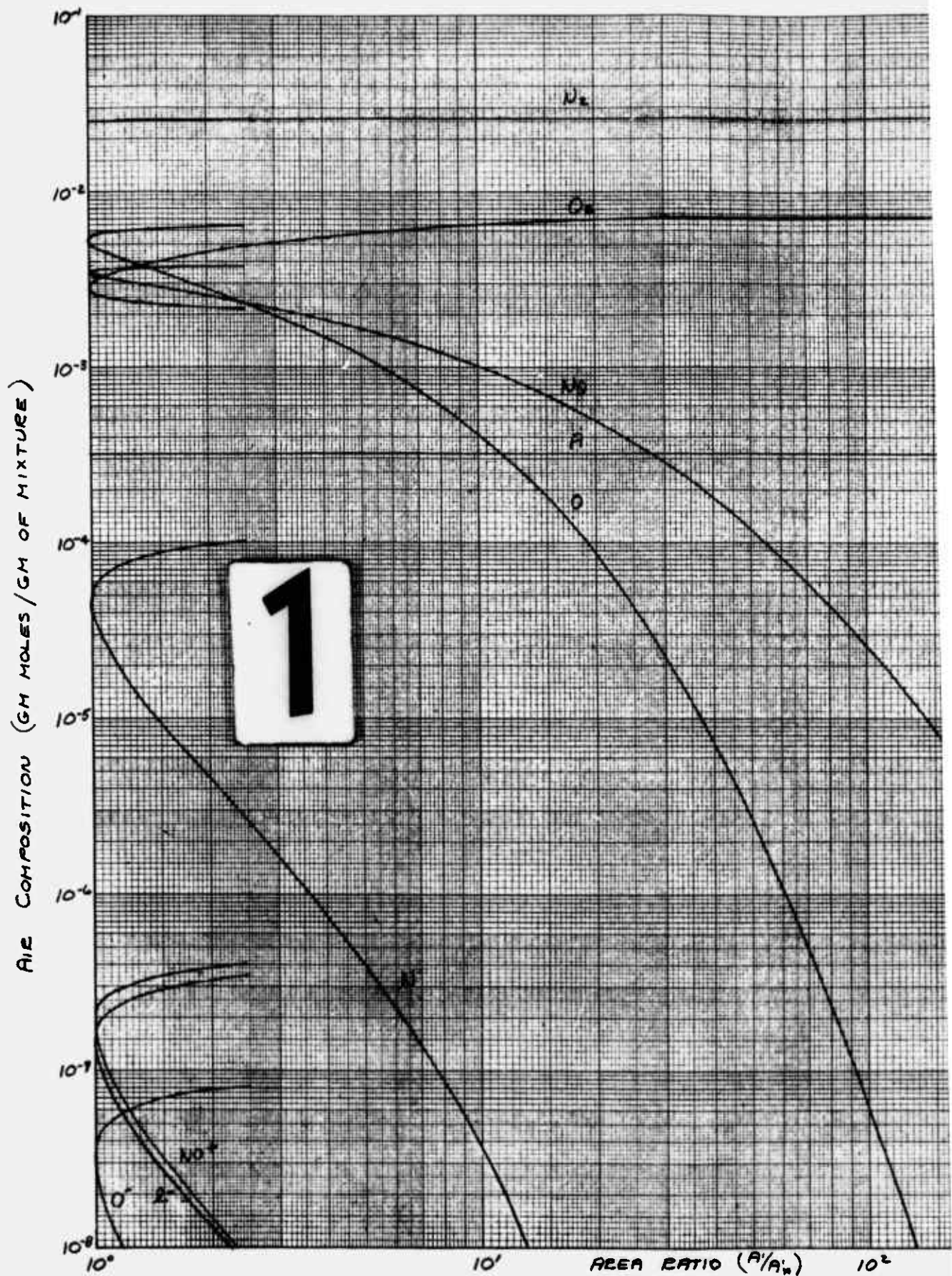
FIGURE NO. 4

$T_0 = 4000^\circ K$
 $P_0 = 1000 \text{ ATM}$

2

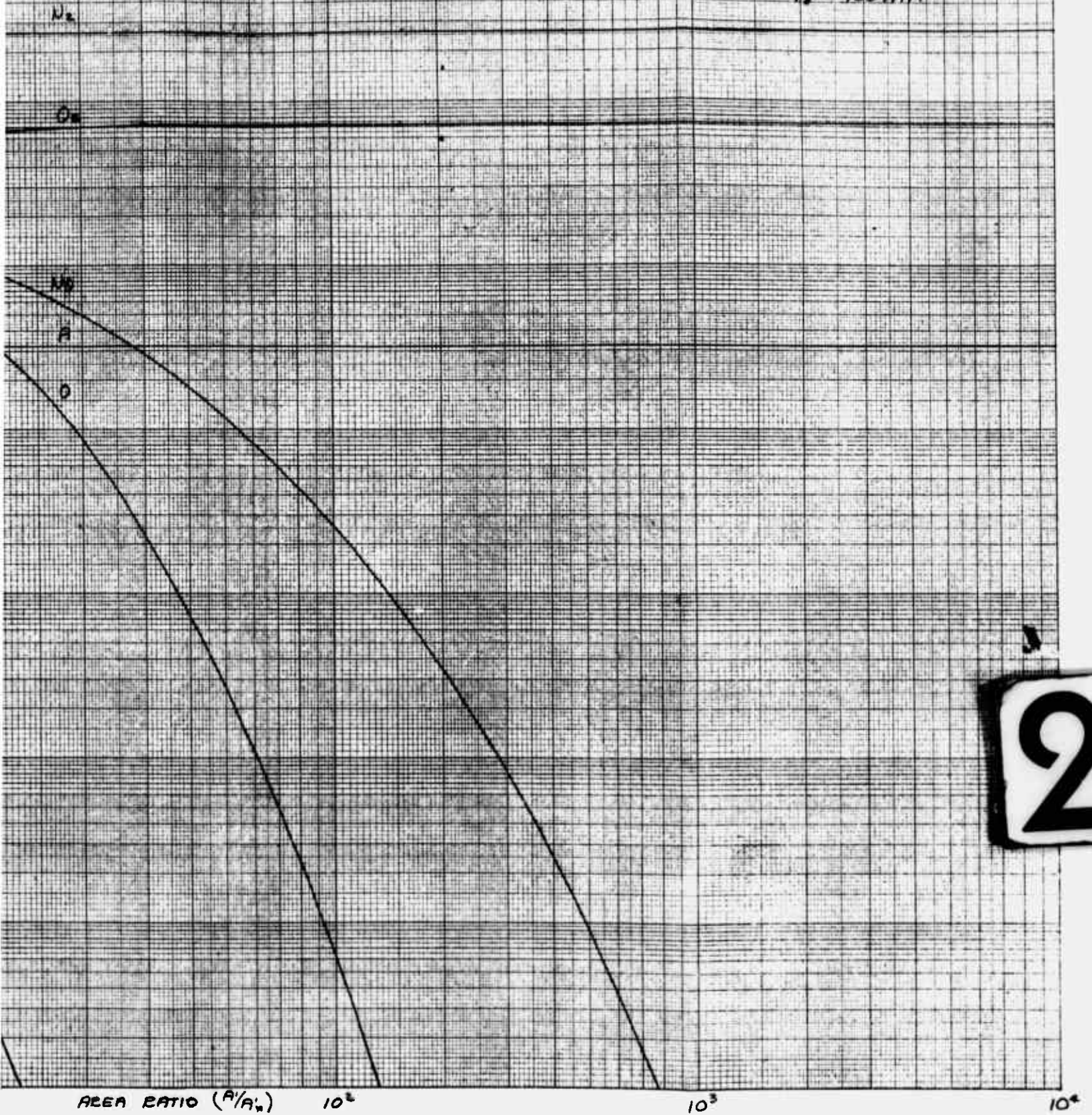


ON VS AREA RATIO IN AN ISENTROPIC EQUILIBRIUM EXPANSION



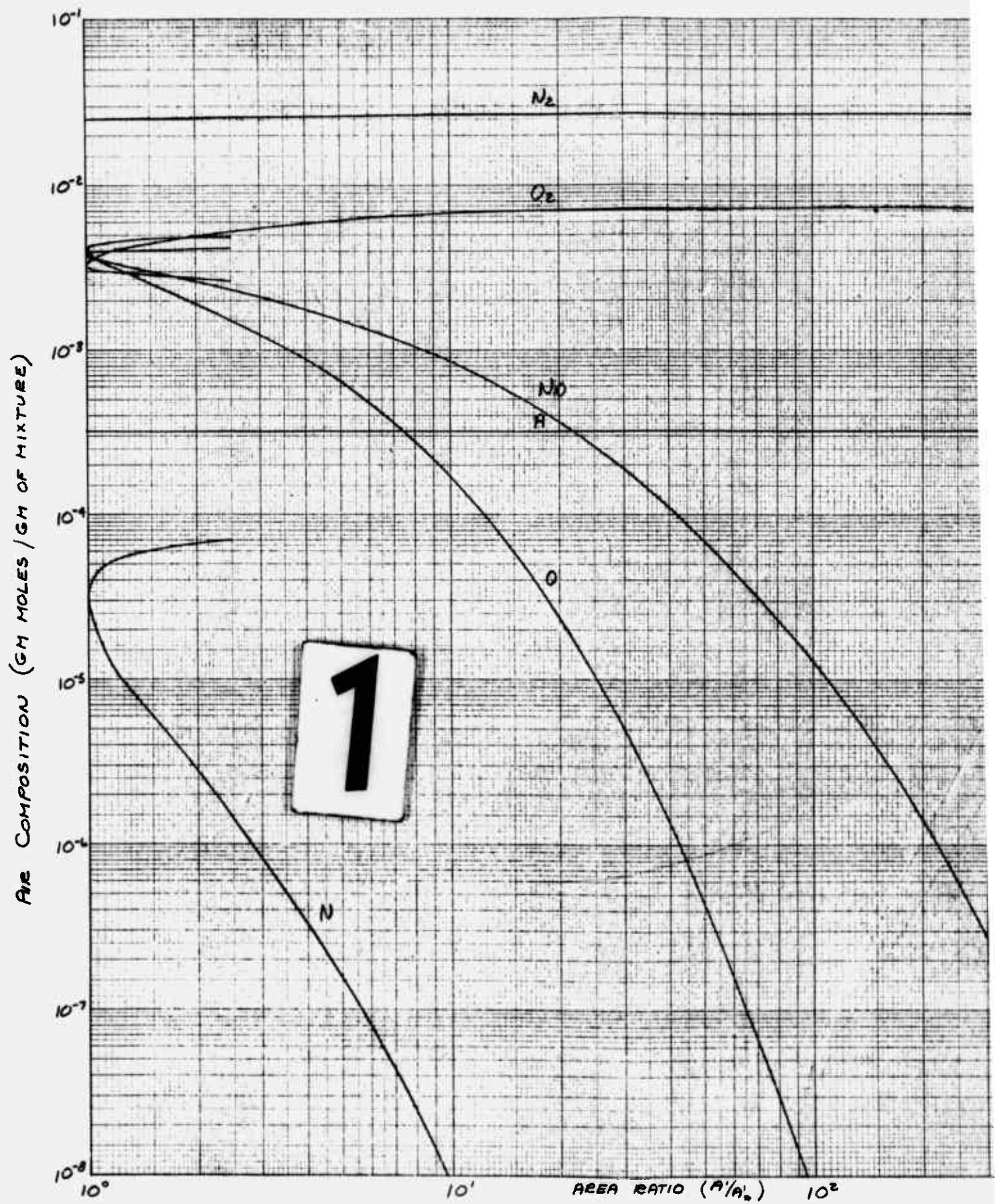
AIR COMPOSITION VS AREA RATIO IN AN ISENTROPIC FLOW

FIGURE NO. 5
 $T_0 = 5000^\circ\text{K}$
 $P_0 = 100\text{ ATM}$



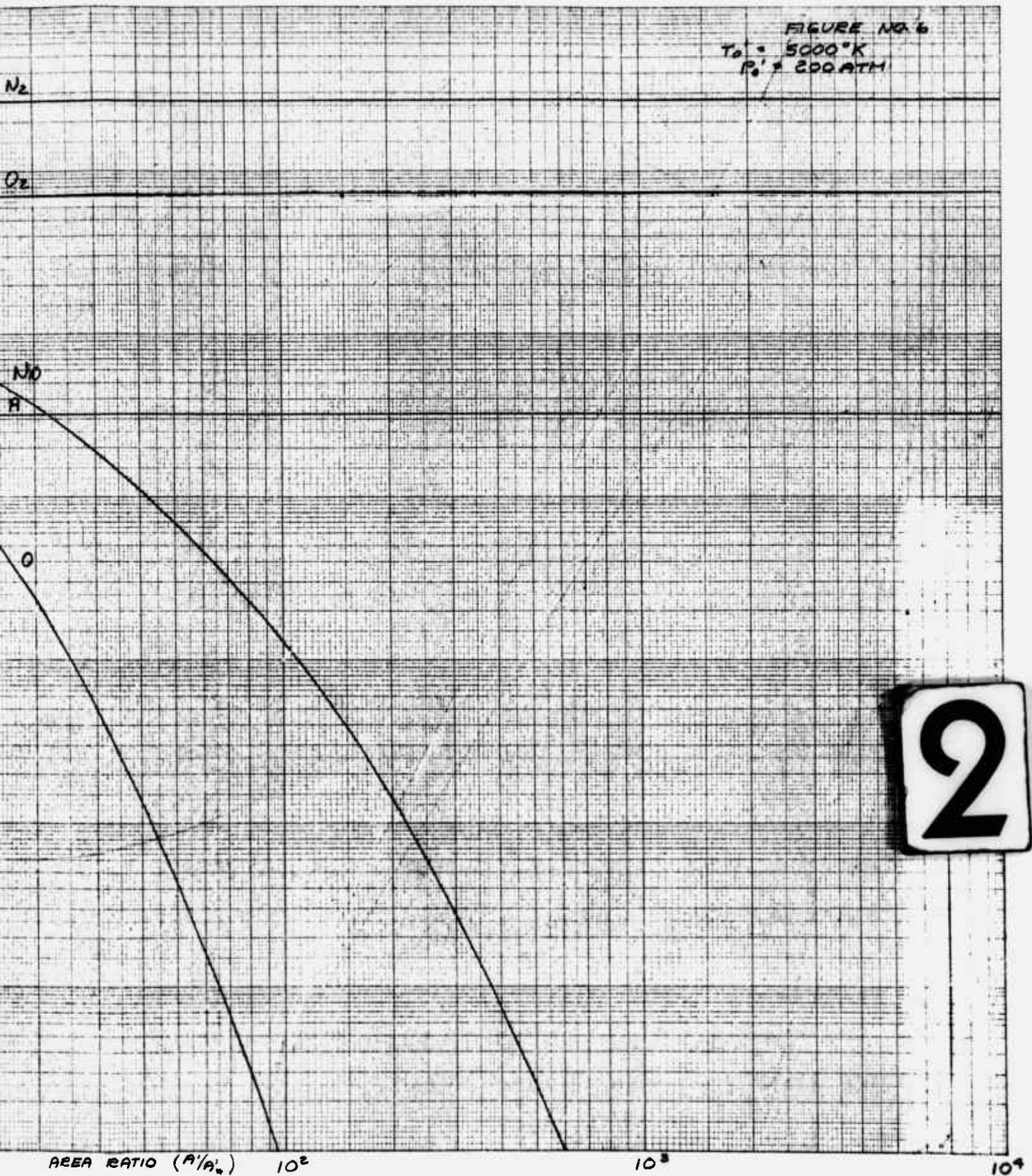
2

VS AREA RATIO IN AN ISENTROPIC EQUILIBRIUM EXPANSION



AIR COMPOSITION VS AREA RATIO IN AN ISENTROPIC FLOW

FIGURE NO. 6
 $T_0 = 5000^\circ\text{K}$
 $P_0 = 200\text{ ATM}$



2

VS AREA RATIO IN AN ISENTROPIC EQUILIBRIUM EXPANSION

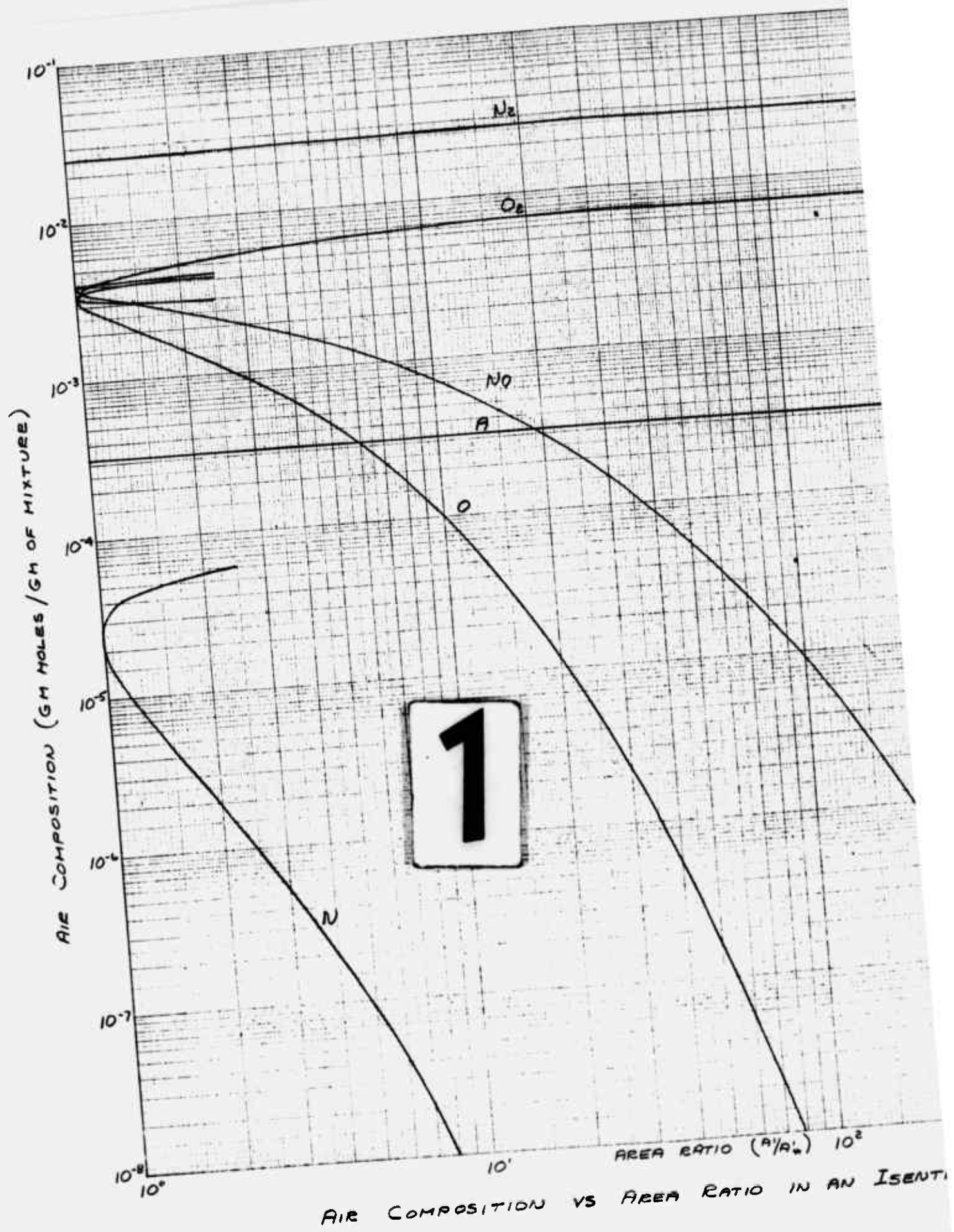


FIGURE NO. 7
 $T_0 = 5000^\circ\text{K}$
 $P_0 = 300\text{ ATM}$

N_2

O_2

10

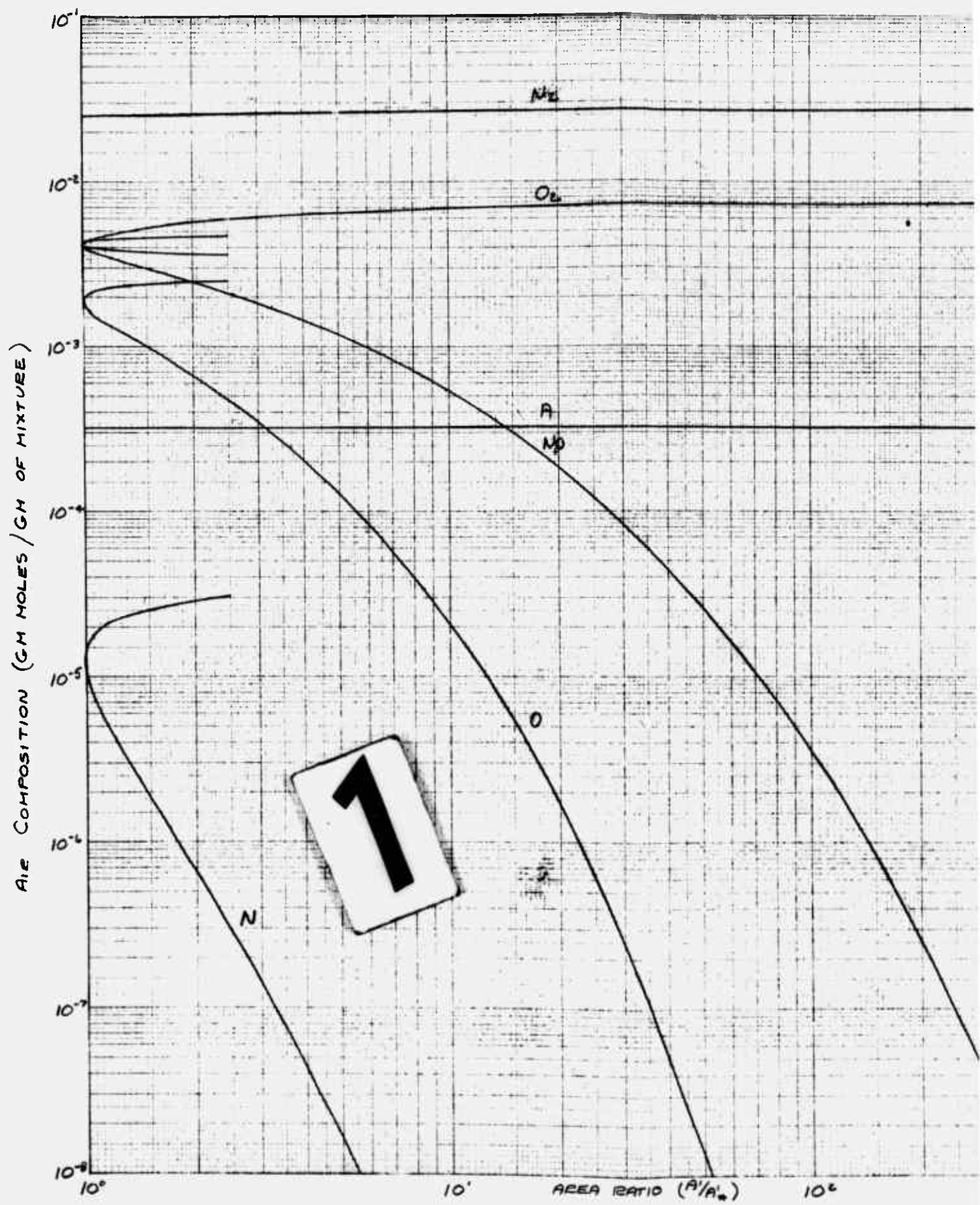
2

AREA RATIO (A/A_*) 10^2

10^3

10^4

VS AREA RATIO IN AN ISENTROPIC EQUILIBRIUM EXPANSION



AIR COMPOSITION VS AREA RATIO IN AN ISENTROPIC

FIGURE NO. 8

$T_0 = 5000^\circ\text{K}$
 $P_0 = 1000 \text{ ATM}$

N₂

O₂

P

30

2

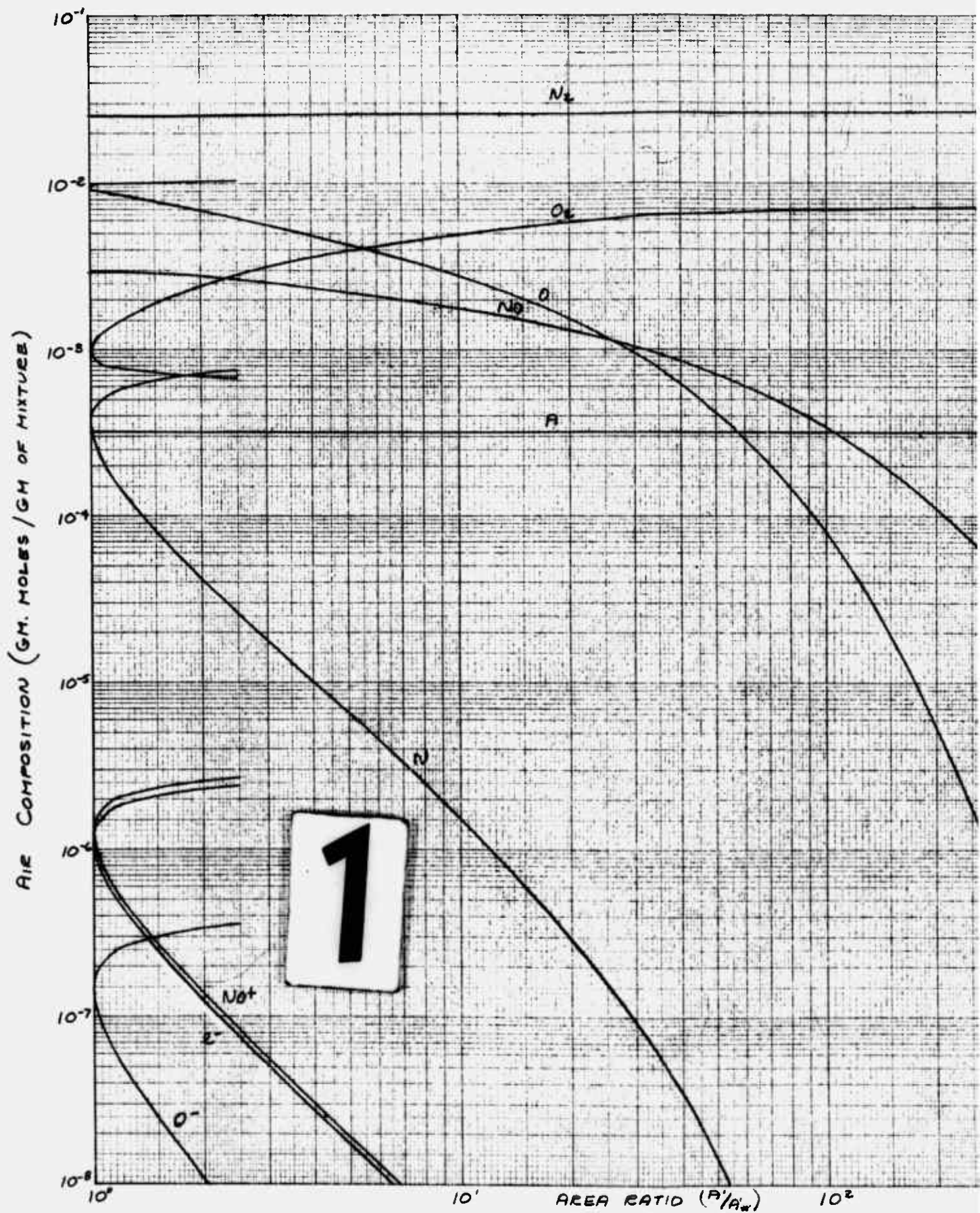
AREA RATIO (A/A_*)

10^2

10^3

10^4

AREA RATIO IN AN ISENTROPIC EQUILIBRIUM EXPANSION



AIR COMPOSITION VS AREA RATIO IN AN ISENTROPIC

FIGURE NO. 9
 $T_0 = 6000^\circ\text{K}$
 $P_0 = 100\text{ ATM}$

N_2

O_2

0

A

2

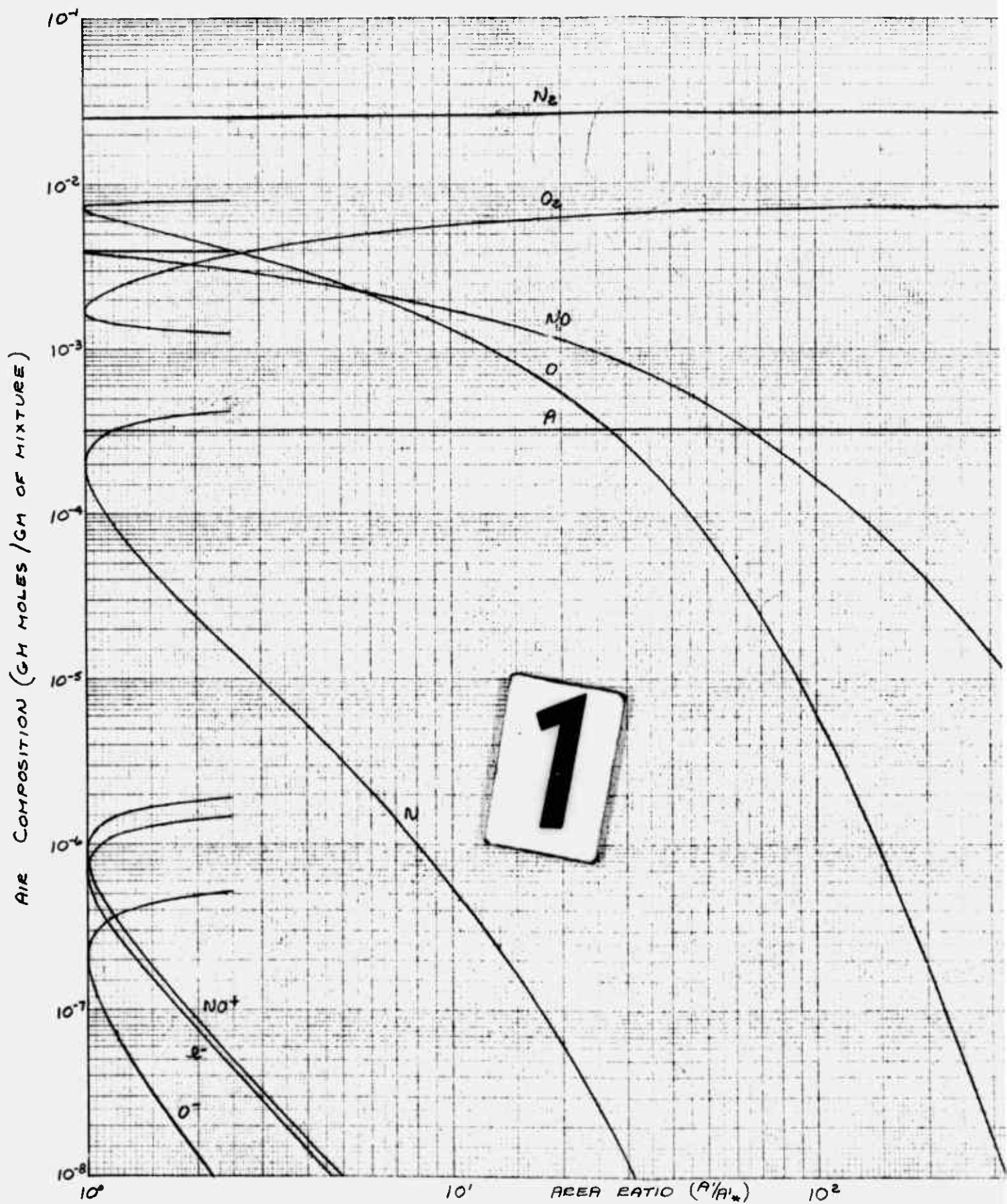
AREA RATIO (A/A_*)

10^2

10^3

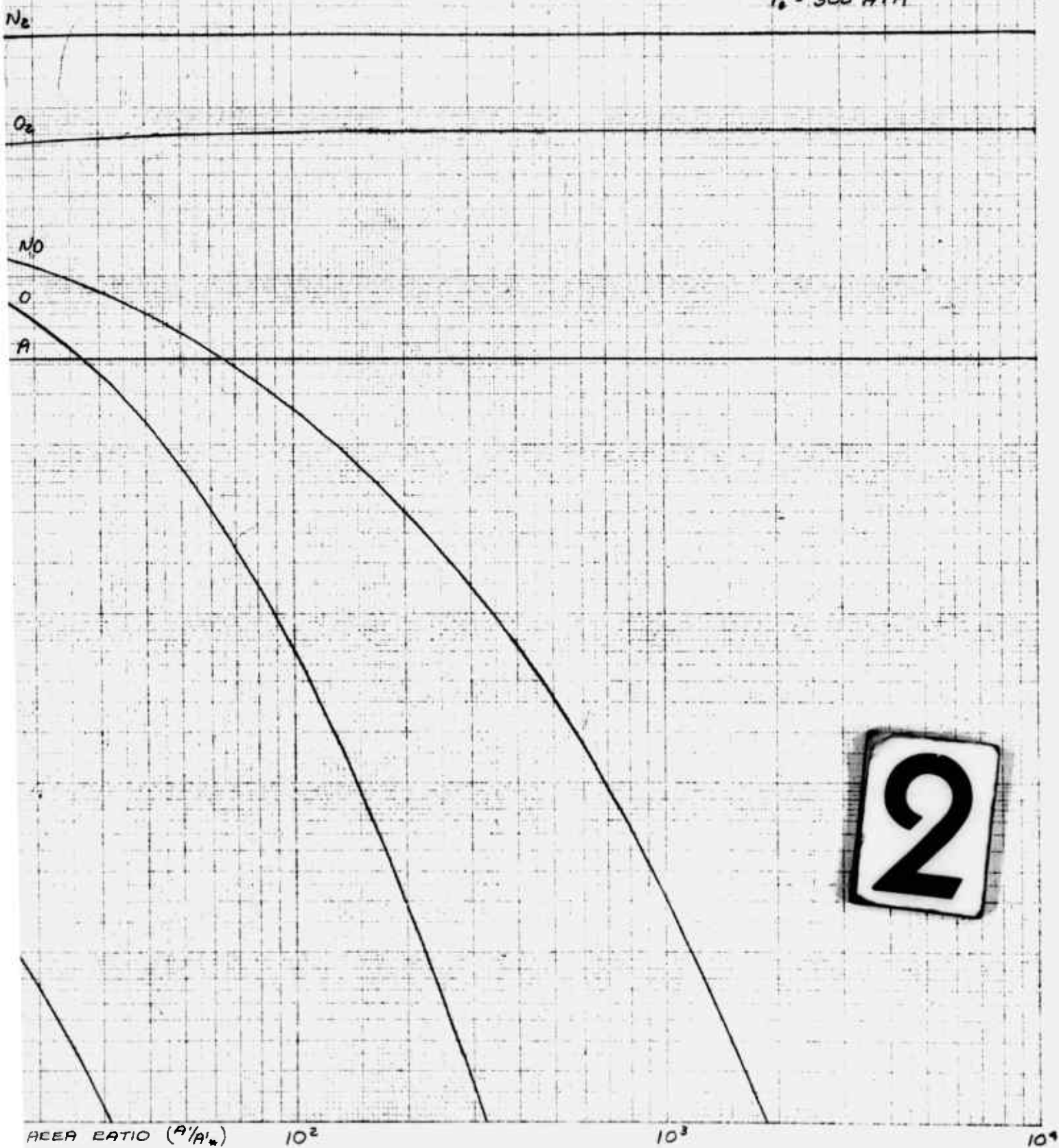
10^4

AREA RATIO IN AN ISENTROPIC EQUILIBRIUM EXPANSION



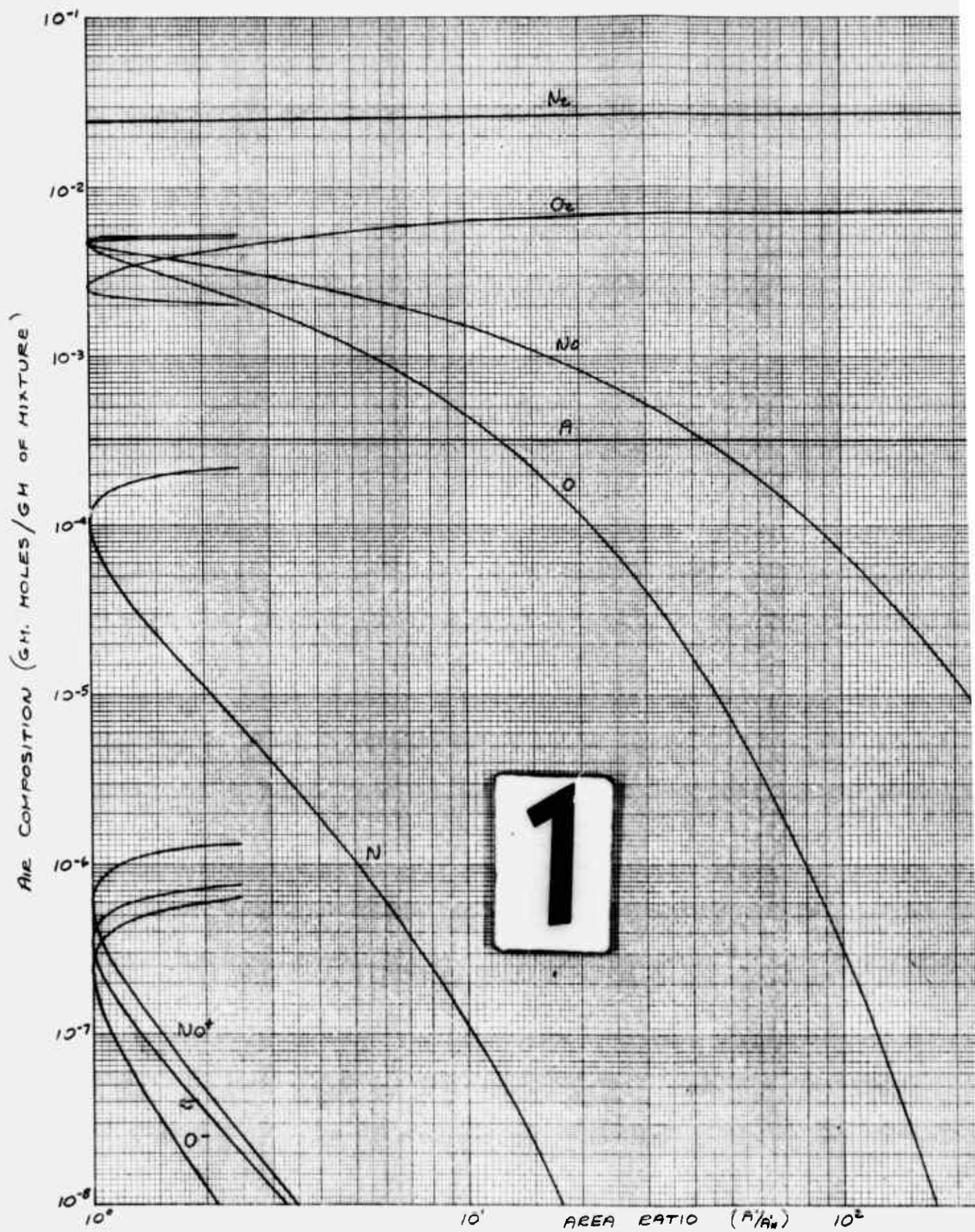
AIR COMPOSITION VS AREA RATIO IN AN ISENTROPIC FLOW

FIGURE NO. 10
 $T_0 = 6000^\circ\text{K}$
 $P_0 = 300\text{ ATM}$



2

AREA RATIO IN AN ISENTROPIC EQUILIBRIUM EXPANSION



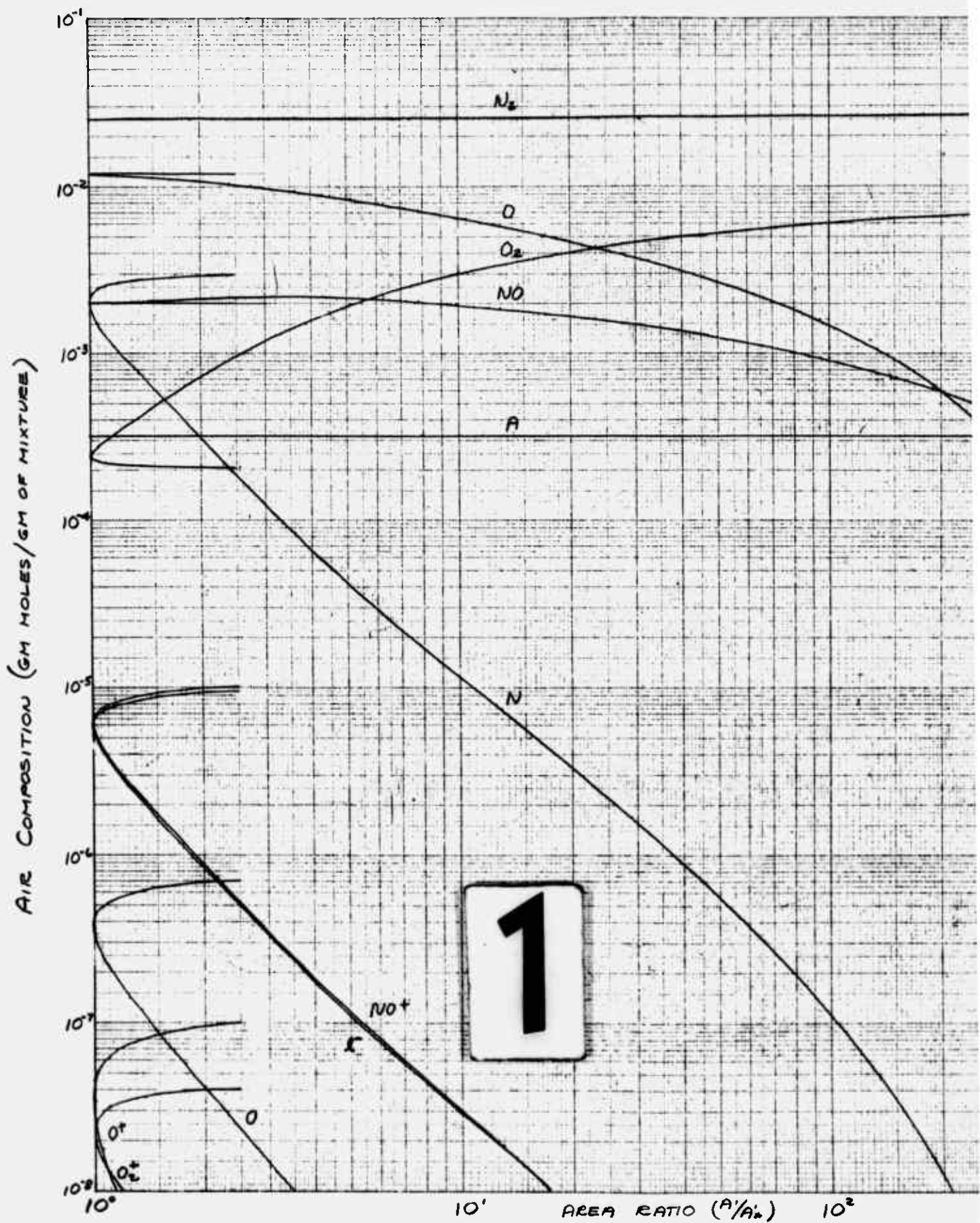
AIR COMPOSITION VS AREA RATIO IN AN ISENTROPIC EQUIL

FIGURE NO. 11
6000°K
1000 ATM

2

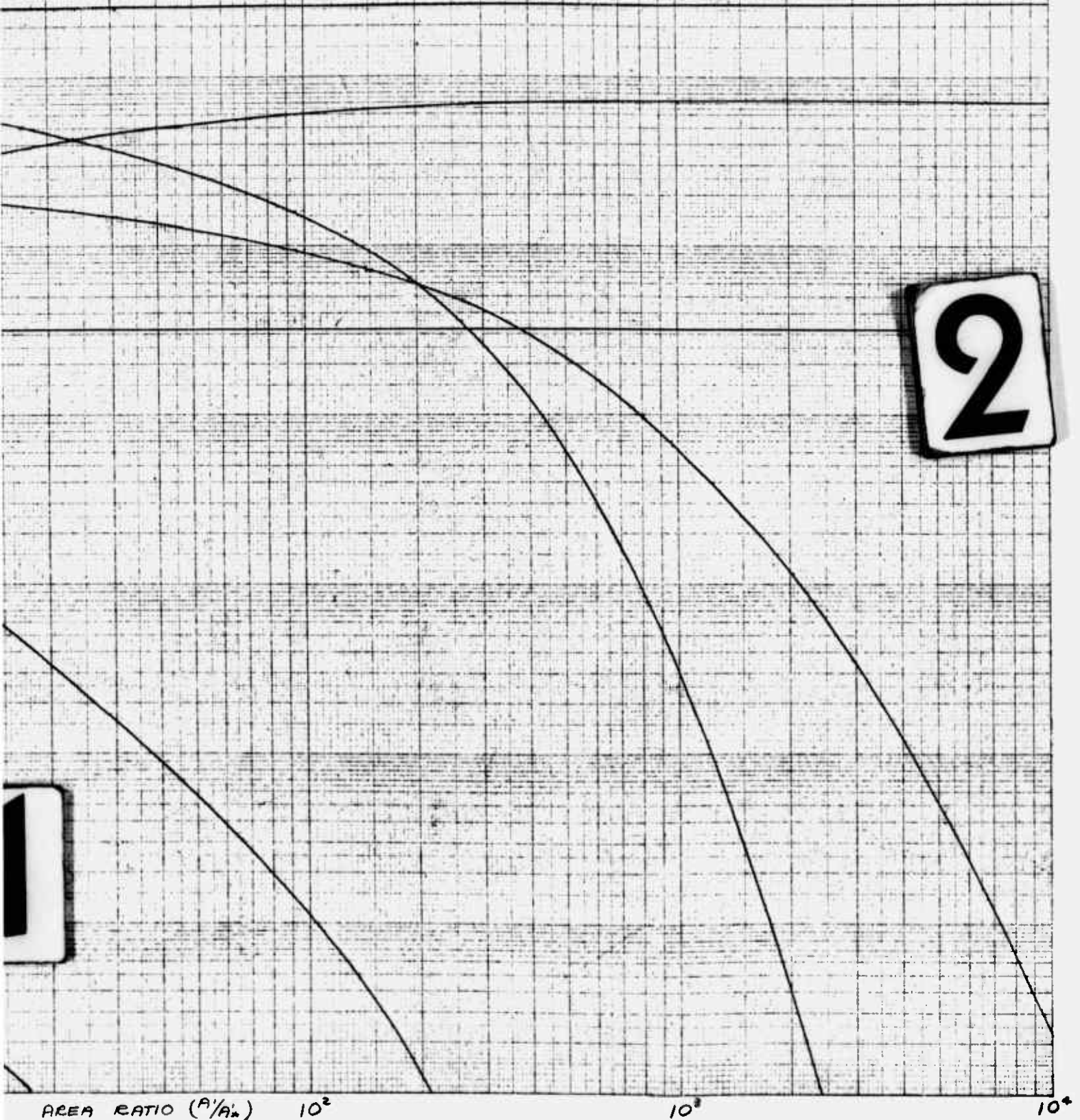
AREA RATIO (A/A_0) 10^2 10^3 10^4

RATIO IN AN ISENTROPIC EQUILIBRIUM EXPANSION

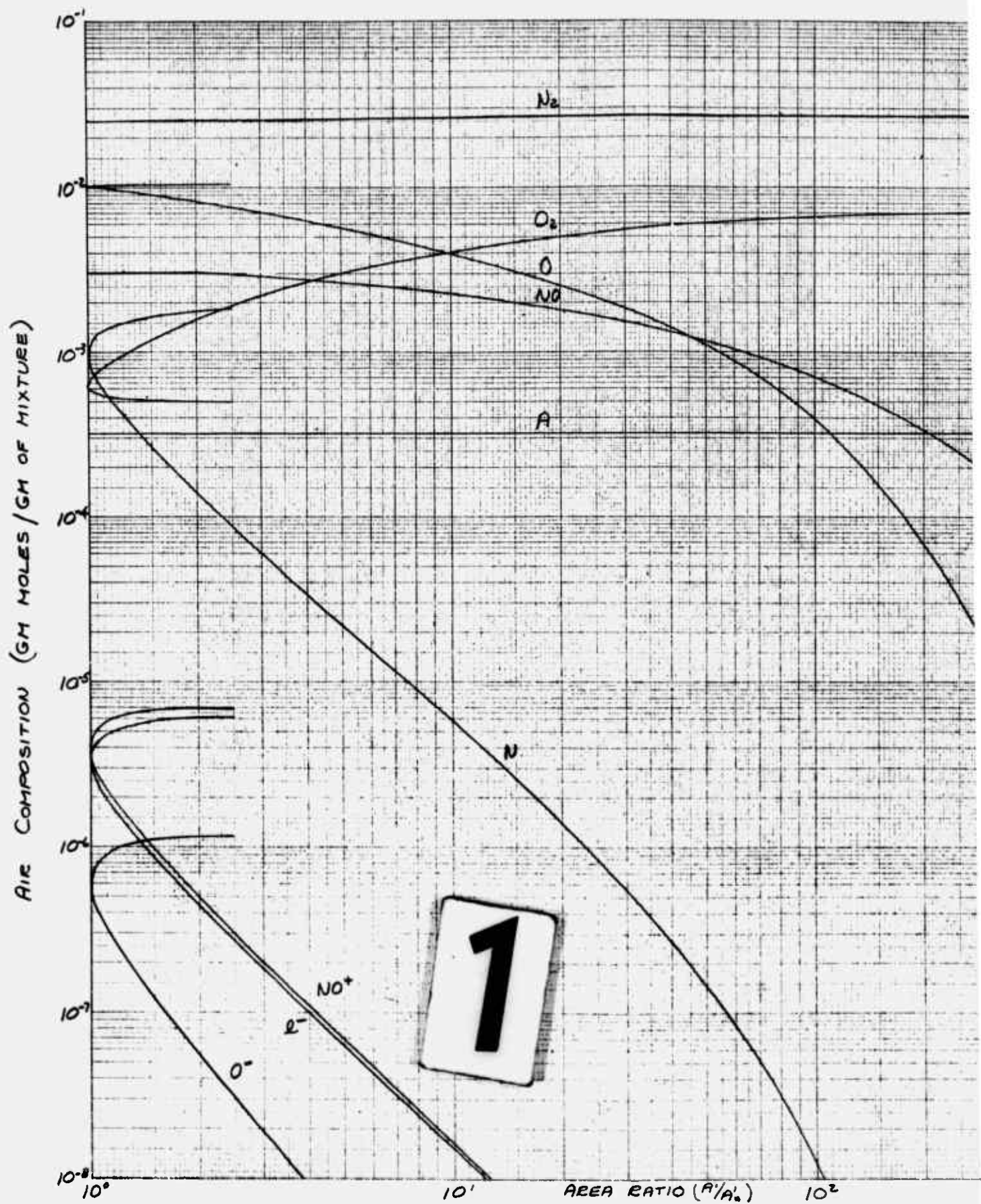


AIR COMPOSITION VS AREA RATIO IN AN ISENTROPIC

FIGURE NO. 12
 $T_0 = 7000^\circ\text{K}$
 $P_0 = 100\text{ ATM}$

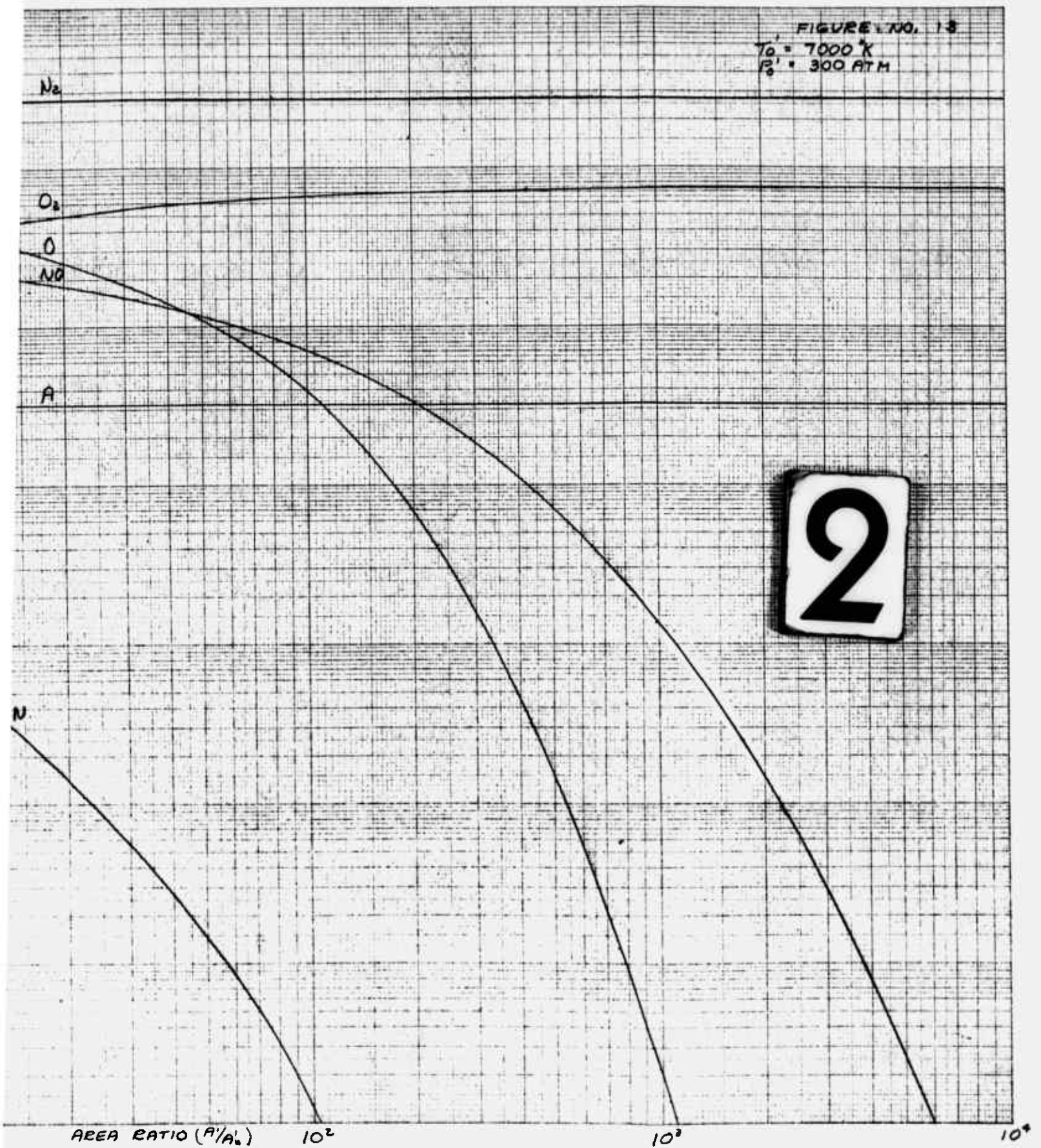


AREA RATIO (A/A^*) 10^2 10^3 10^4
VS AREA RATIO IN AN ISENTROPIC EQUILIBRIUM EXPANSION



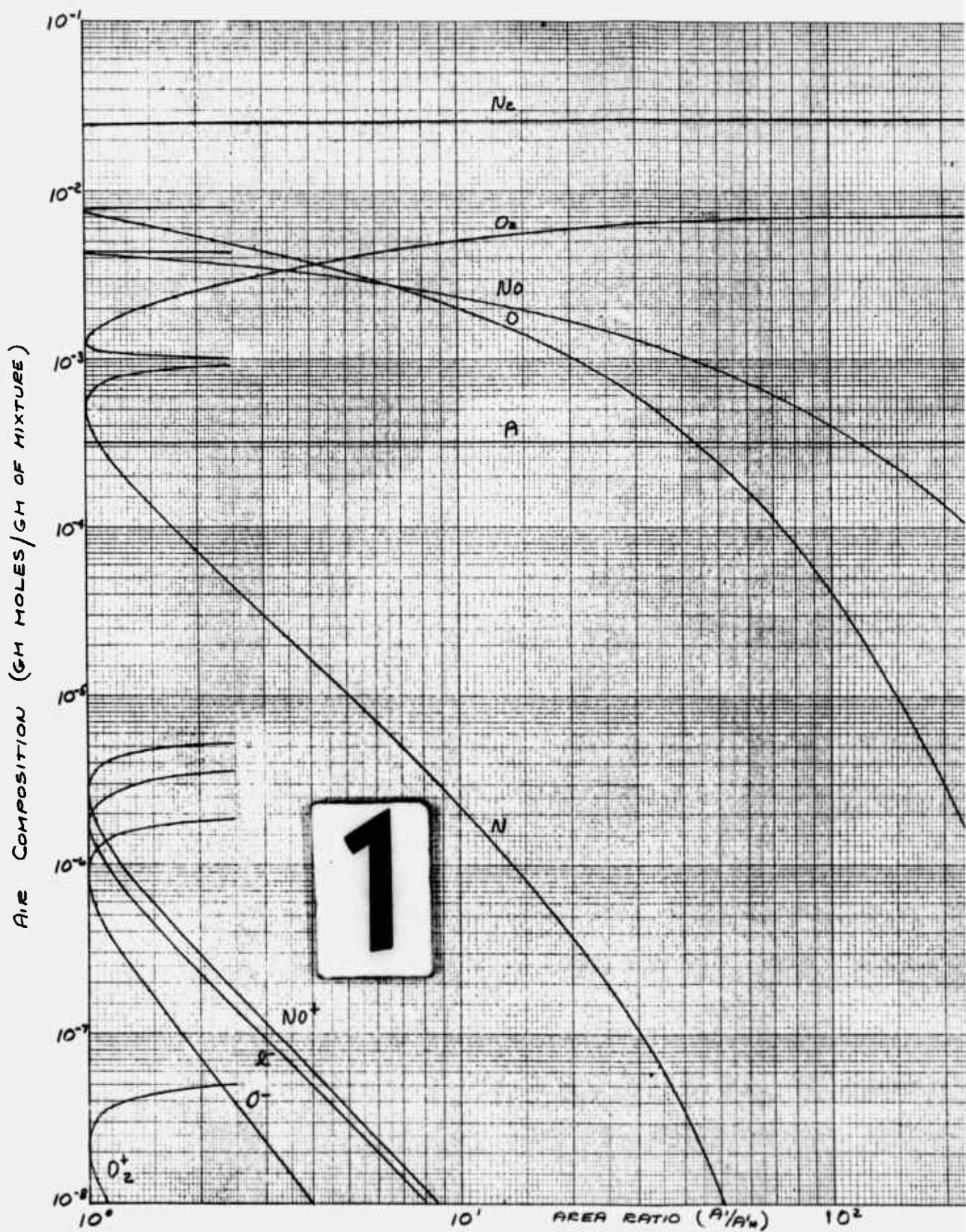
AIR COMPOSITION VS AREA RATIO IN AN ISENTROPIC FLOW

FIGURE NO. 13
 $T_0 = 7000^\circ \text{K}$
 $P_0 = 300 \text{ ATM}$



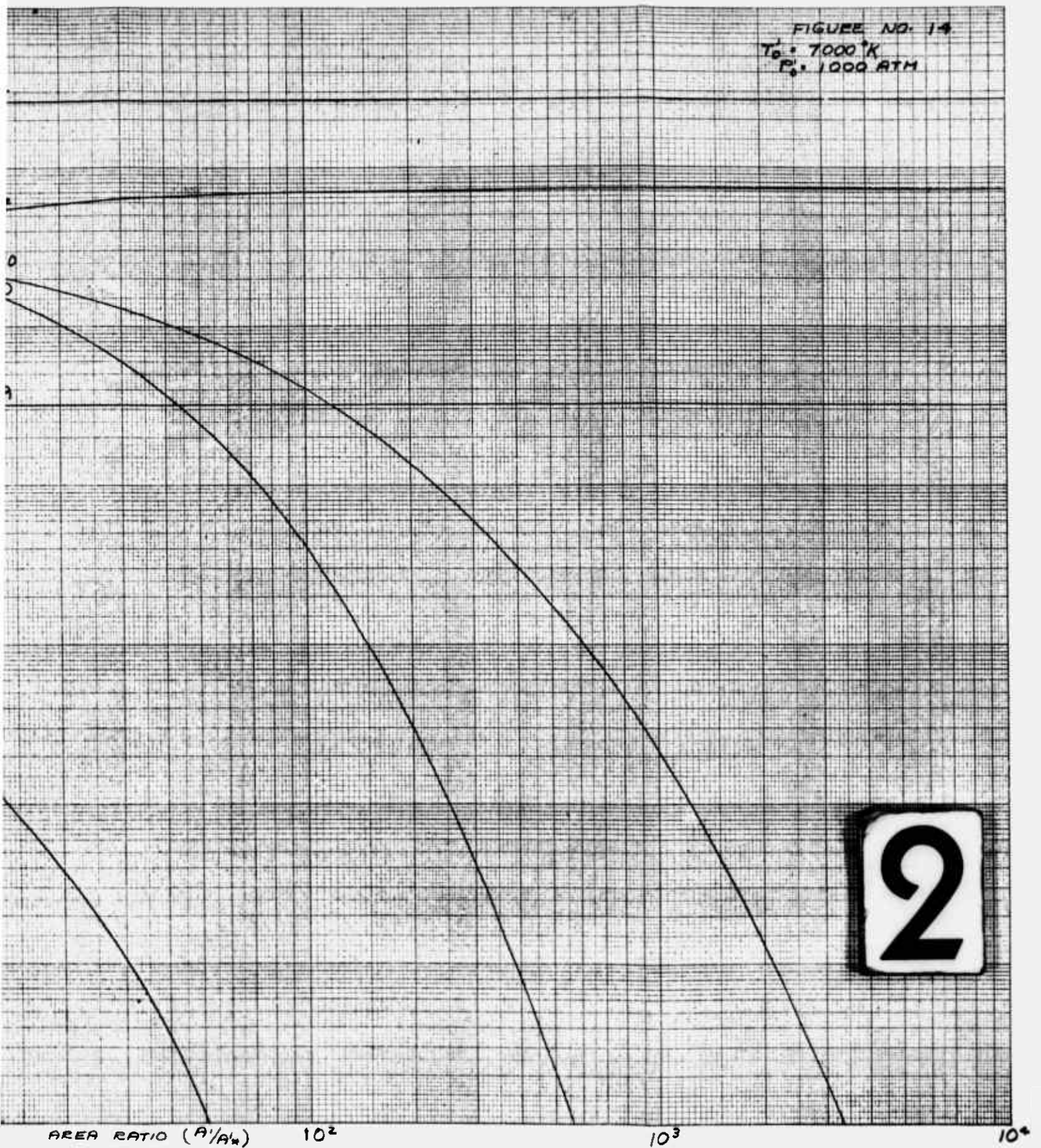
2

ON VS AREA RATIO IN AN ISENTROPIC EQUILIBRIUM EXPANSION



Air Composition vs Area Ratio in an Isentropic

FIGURE NO. 14
 $T_0 = 7000^\circ K$
 $P_0 = 1000 \text{ ATM}$



AREA RATIO (A/A^*)

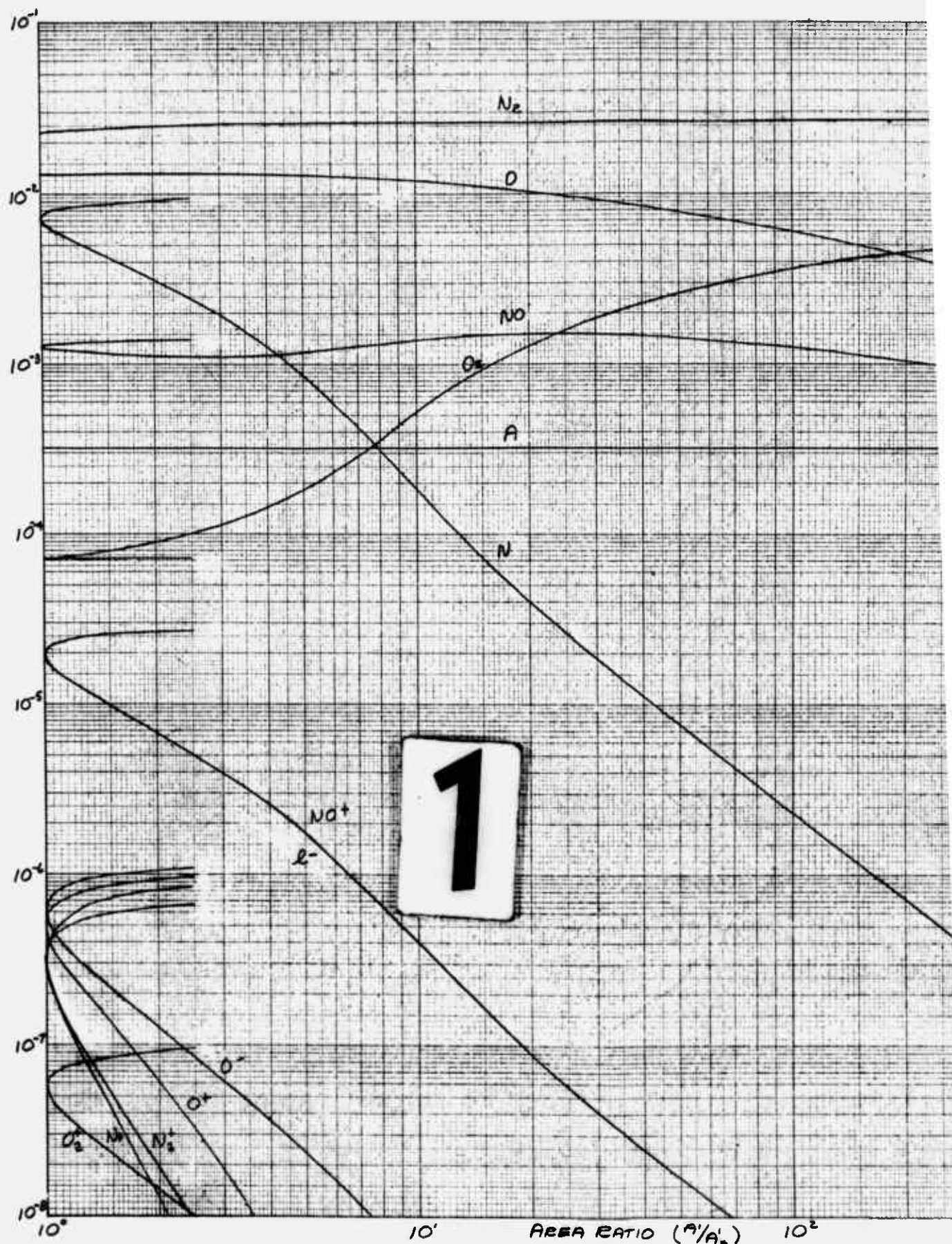
10^2

10^3

10^4

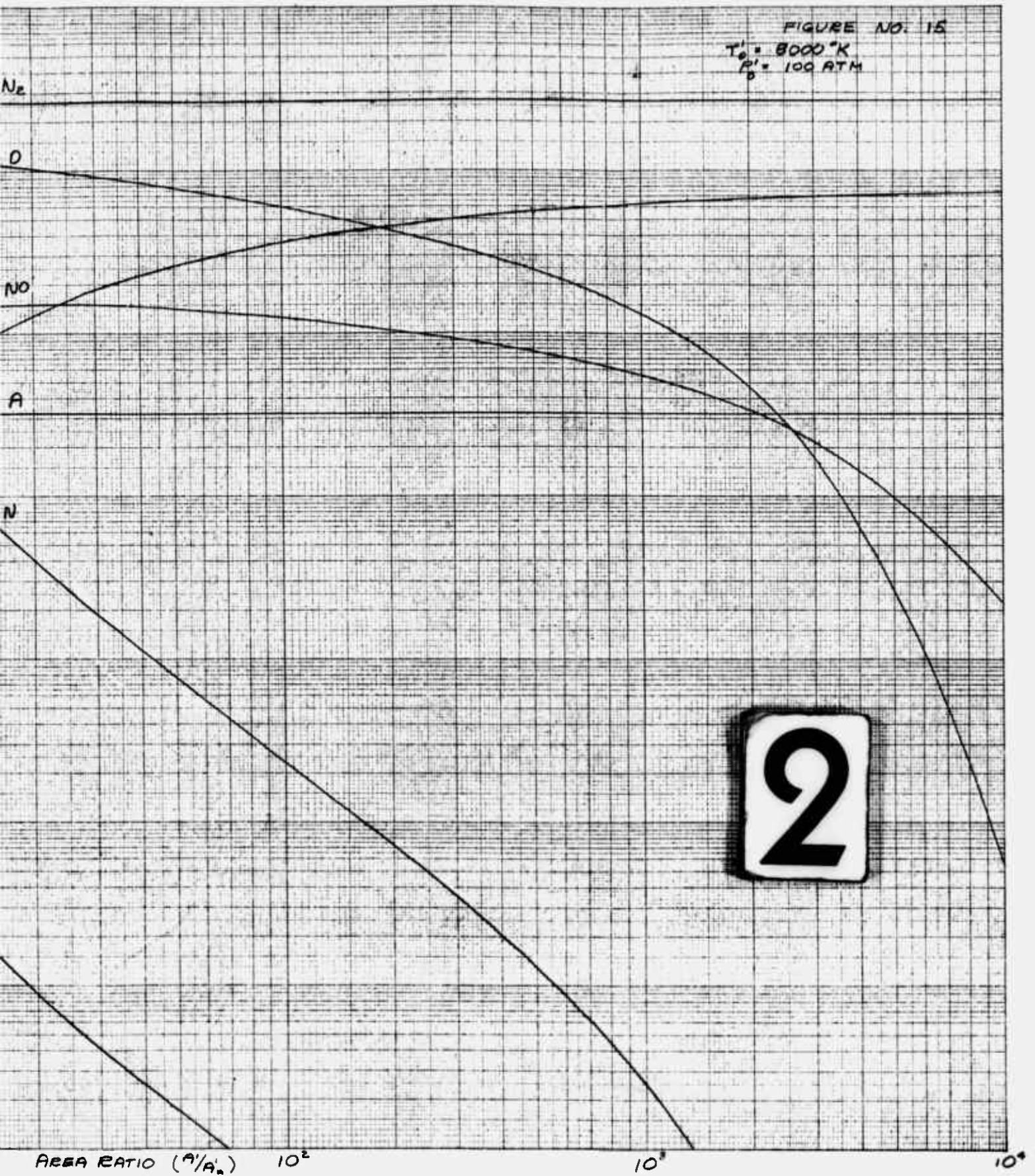
VS AREA RATIO IN AN ISENTROPIC EQUILIBRIUM EXPANSION

AIR COMPOSITION (GM MOLES/GM OF MIXTURE)

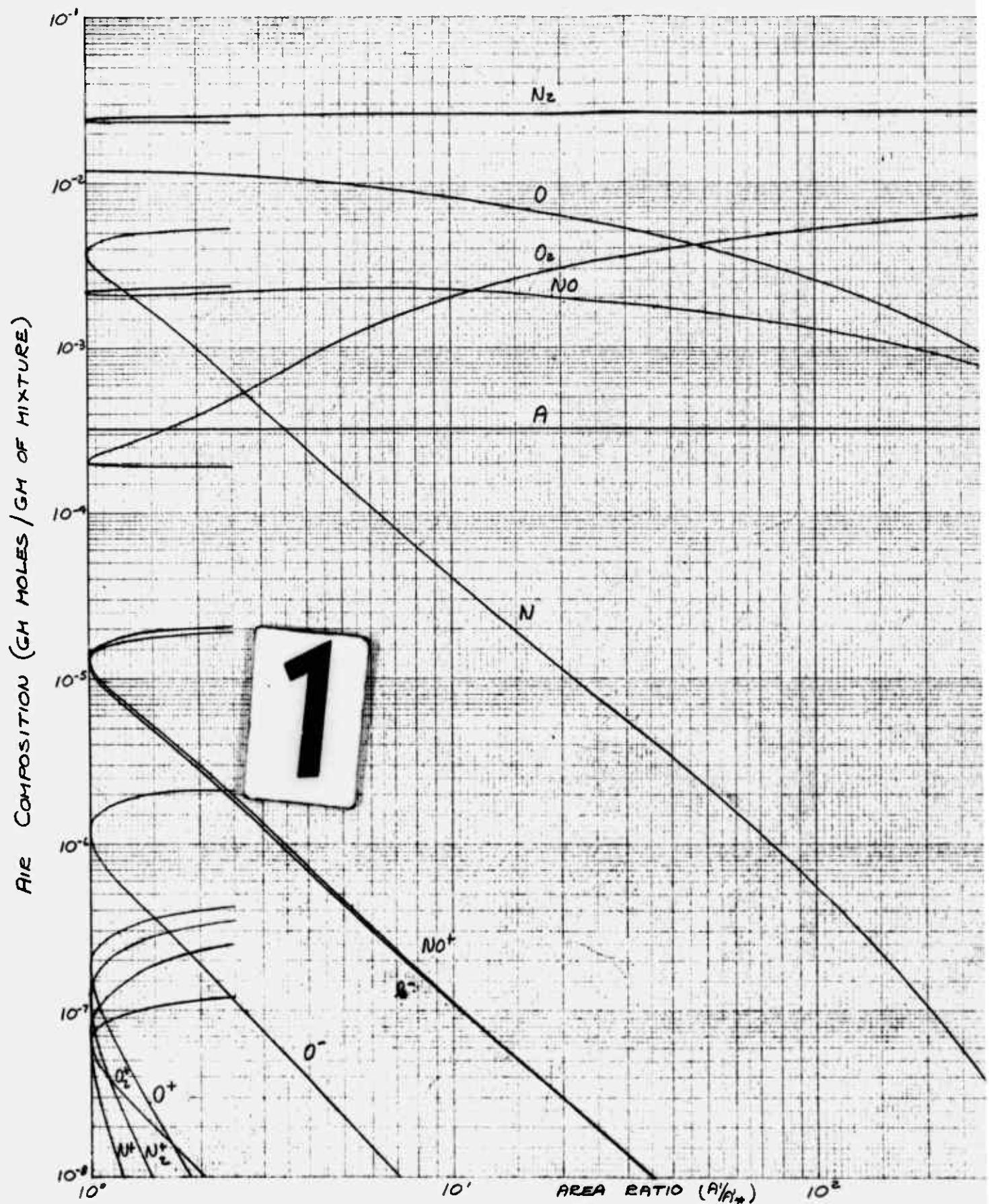


AIR COMPOSITION VS AREA RATIO IN AN ISENTROPIC FLOW

FIGURE NO. 15
 $T_0' = 8000^\circ\text{K}$
 $P_0' = 100\text{ ATM}$



VS AREA RATIO IN AN ISENTROPIC EQUILIBRIUM EXPANSION

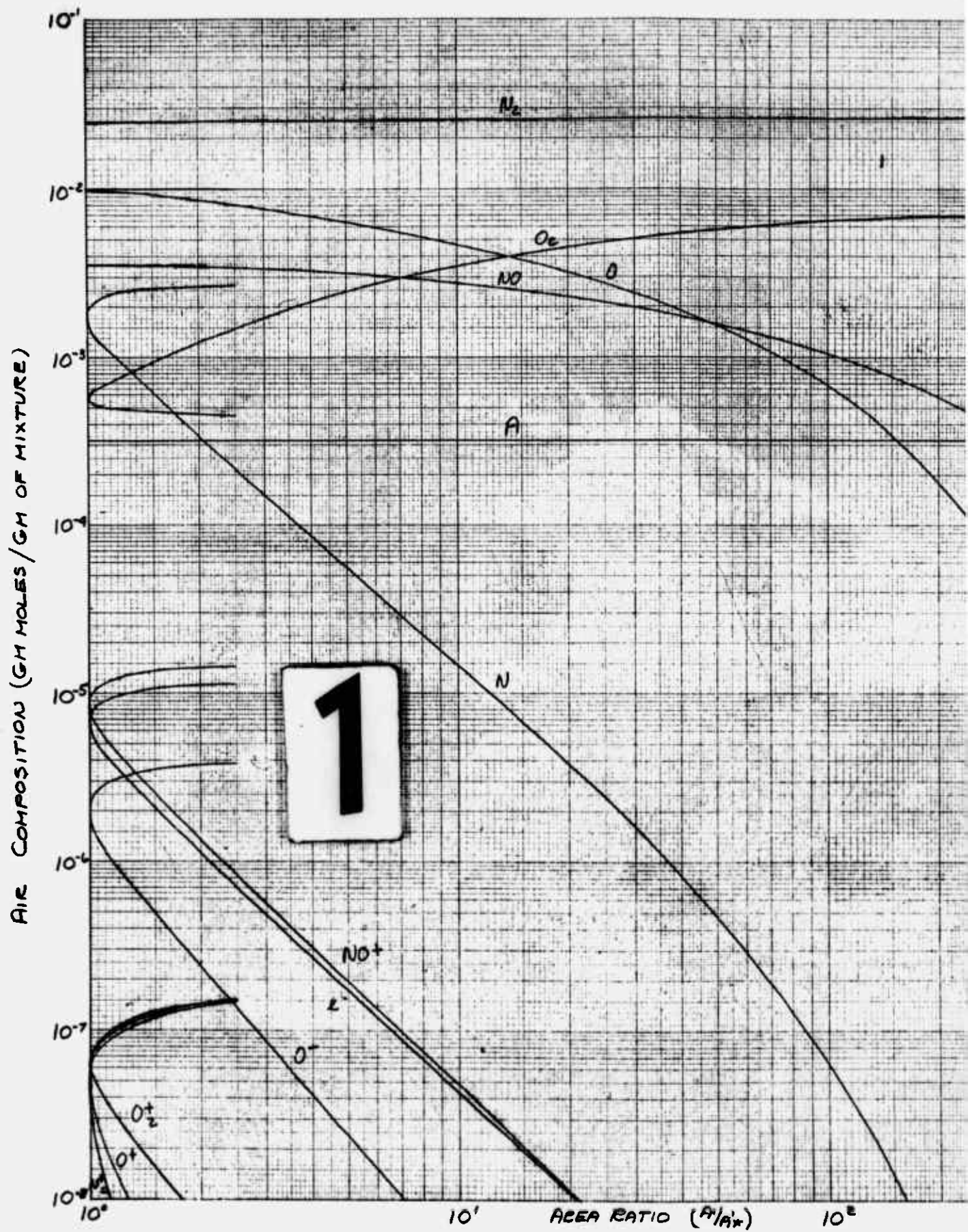


AIR COMPOSITION VS AREA RATIO IN AN ISENTROPIC

2

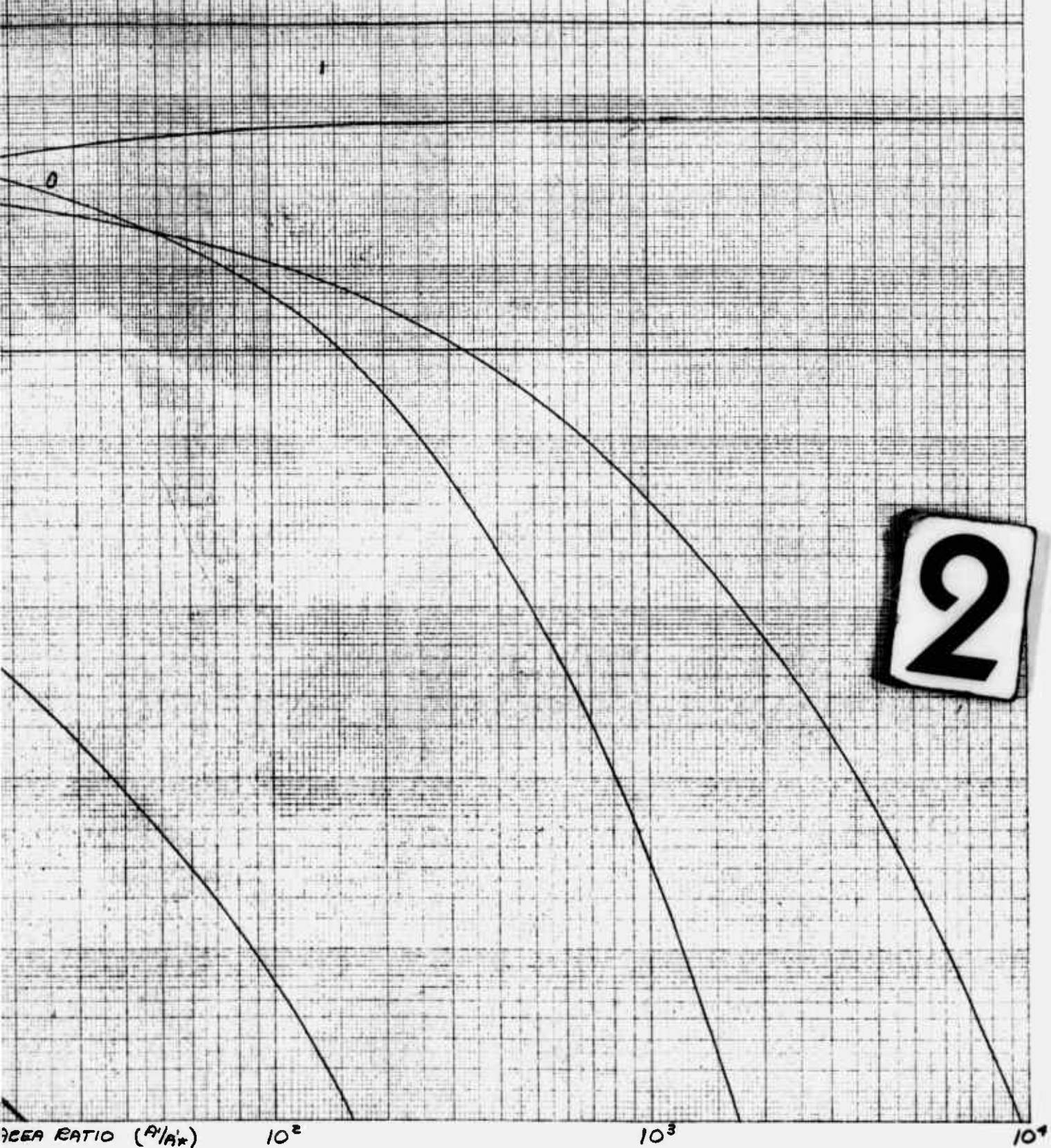


83



AIR COMPOSITION VS AREA RATIO IN AN ISENTROPIC

FIGURE NO. 17
 $T_0 = 8000^\circ\text{K}$
 $P_0 = 1000\text{ ATM}$



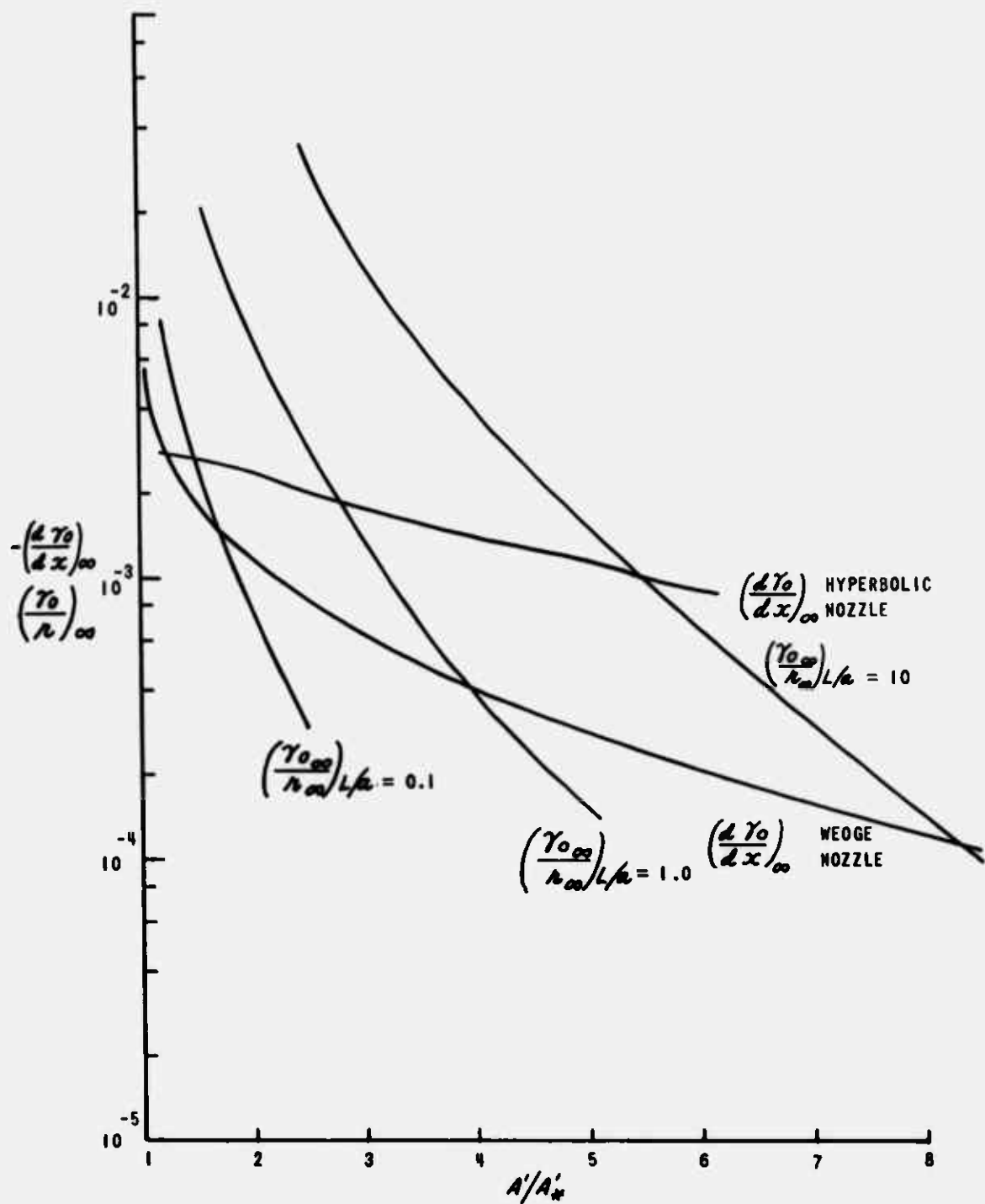


Figure 18 SAMPLE SOLUTION OF EQUATION (36) $T_0' = 5000^\circ \text{ K}$, $P_0' = 100 \text{ ATM}$.

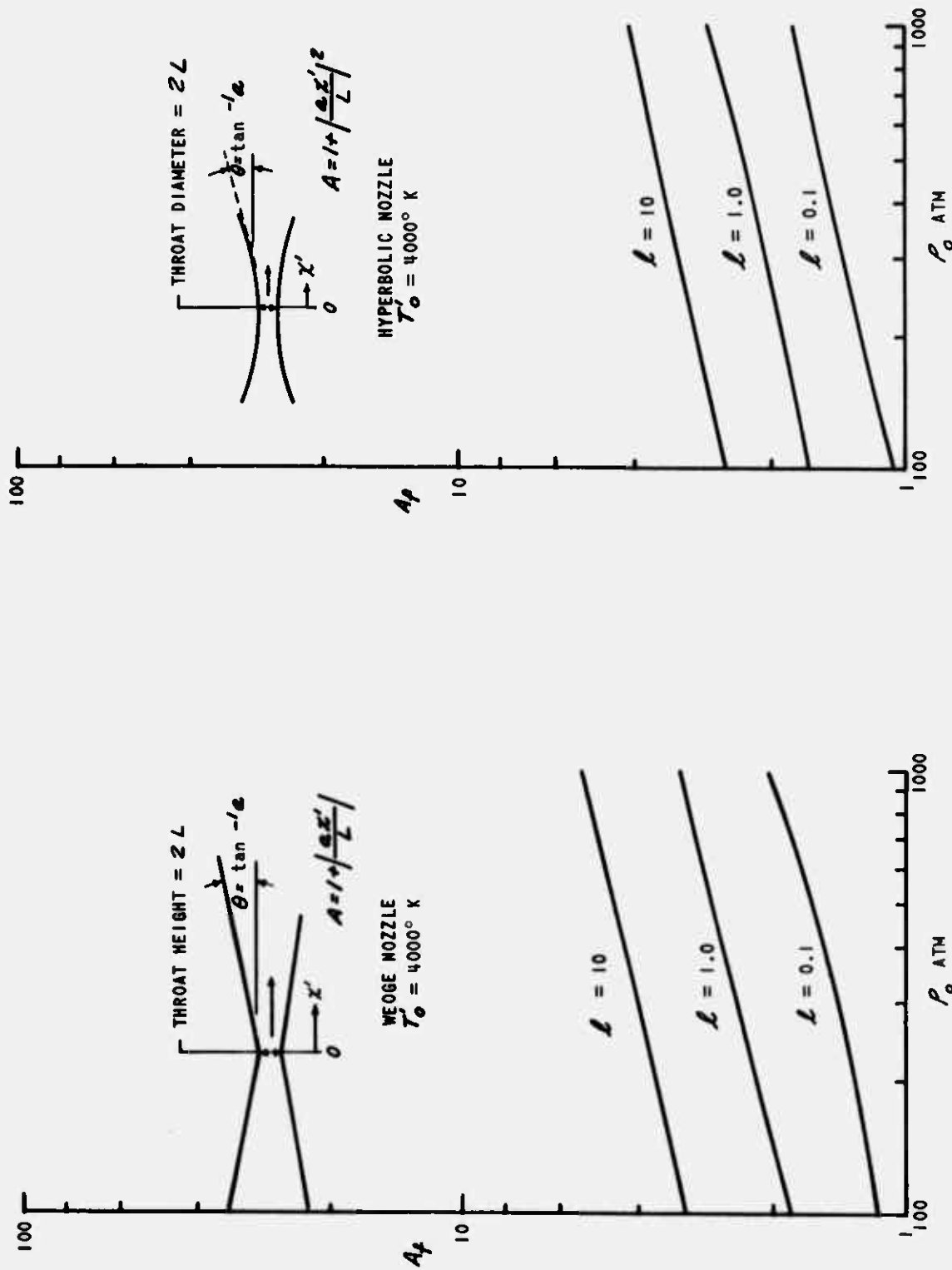


Figure 19 (a) EFFECT OF STAGNATION PRESSURE AND NOZZLE GEOMETRY ON THE AREA RATIO FOR FREEZING FOR A SIMPLIFIED AIR MODEL.

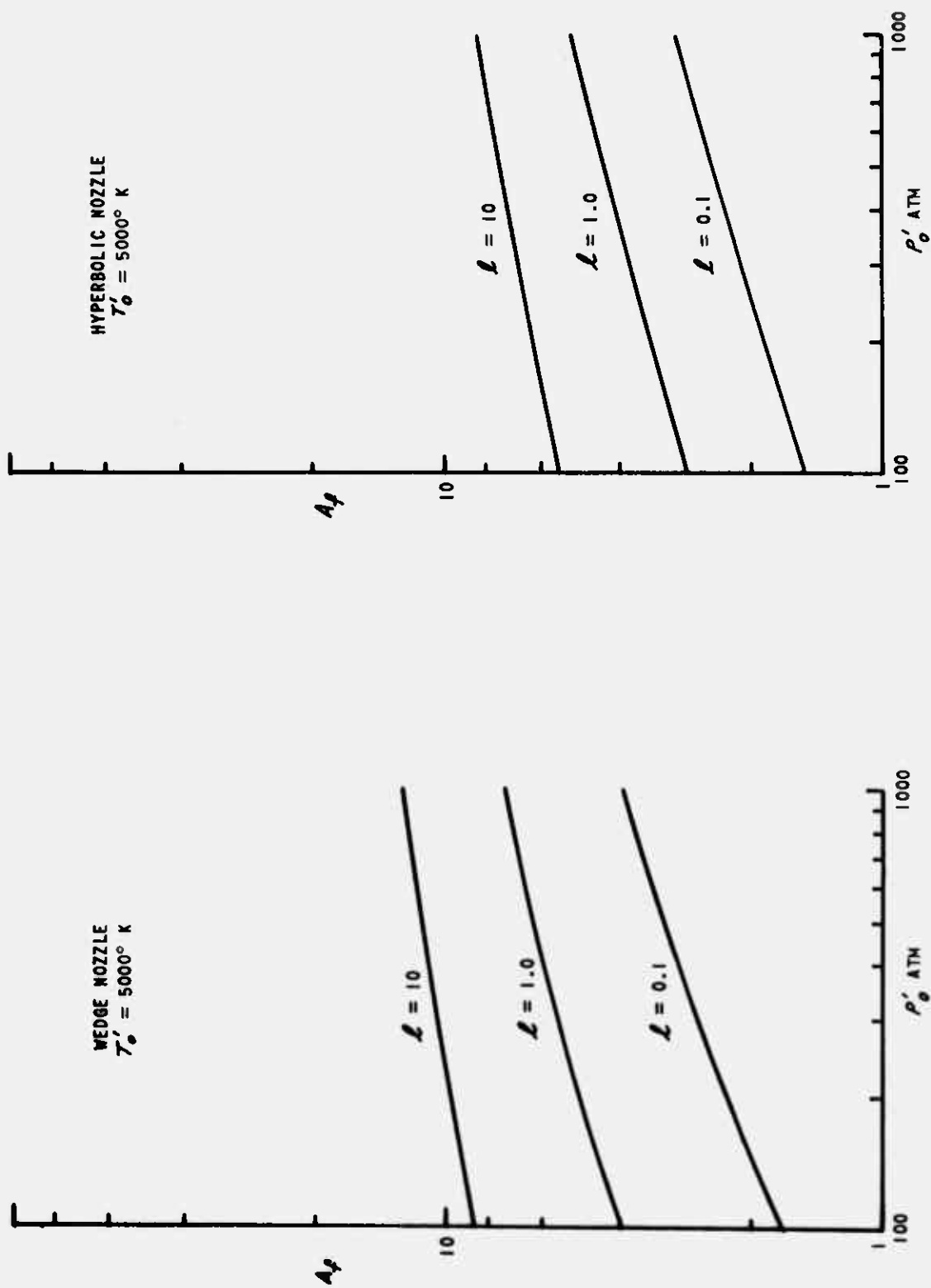


Figure 19 (b) EFFECT OF STAGNATION PRESSURE AND NOZZLE GEOMETRY ON THE AREA RATIO FOR FREEZING FOR A SIMPLIFIED AIR MODEL.

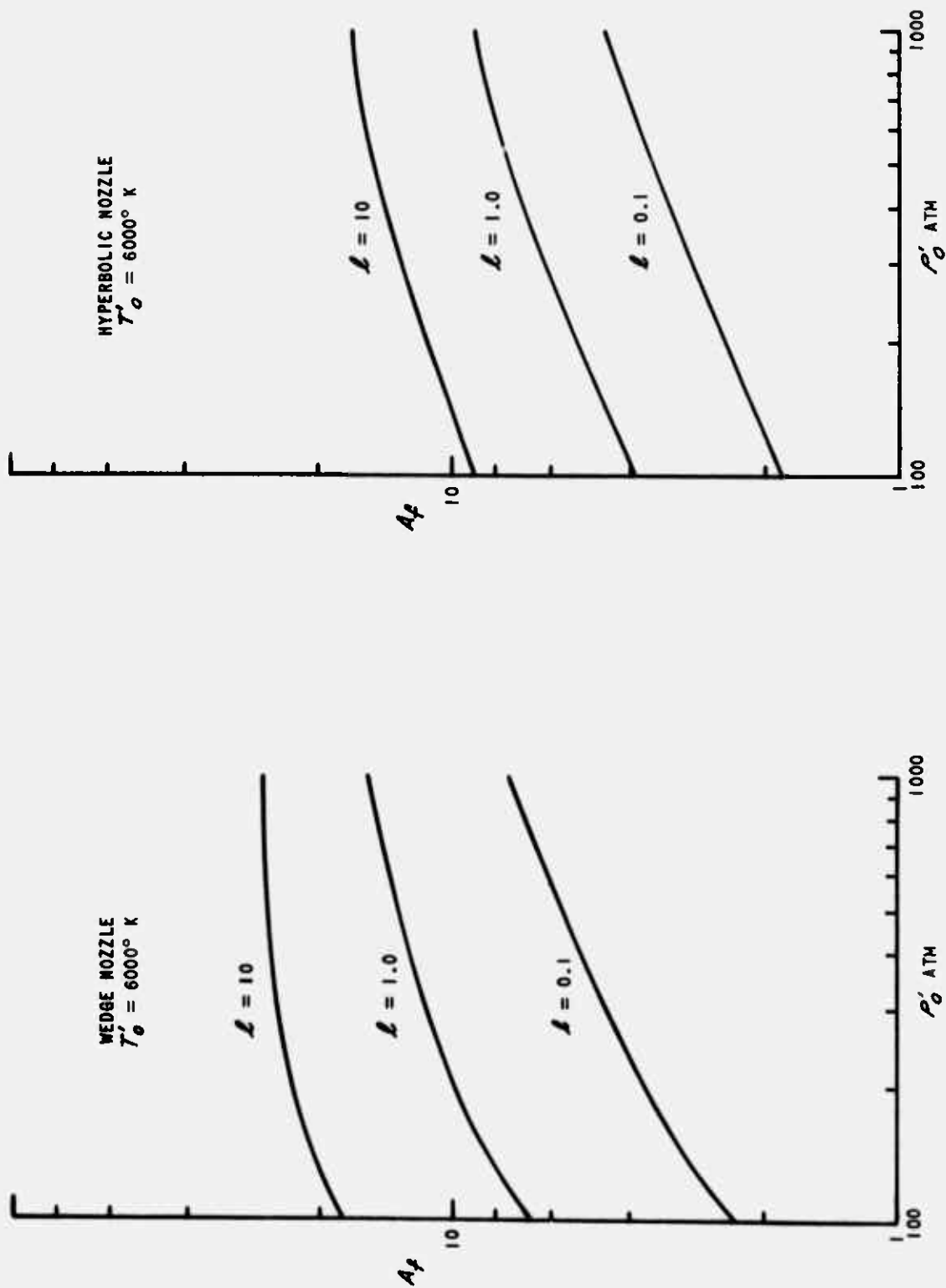


Figure 19 (c) EFFECT OF STAGNATION PRESSURE AND NOZZLE GEOMETRY ON THE AREA RATIO FOR FREEZING FOR A SIMPLIFIED AIR MODEL.

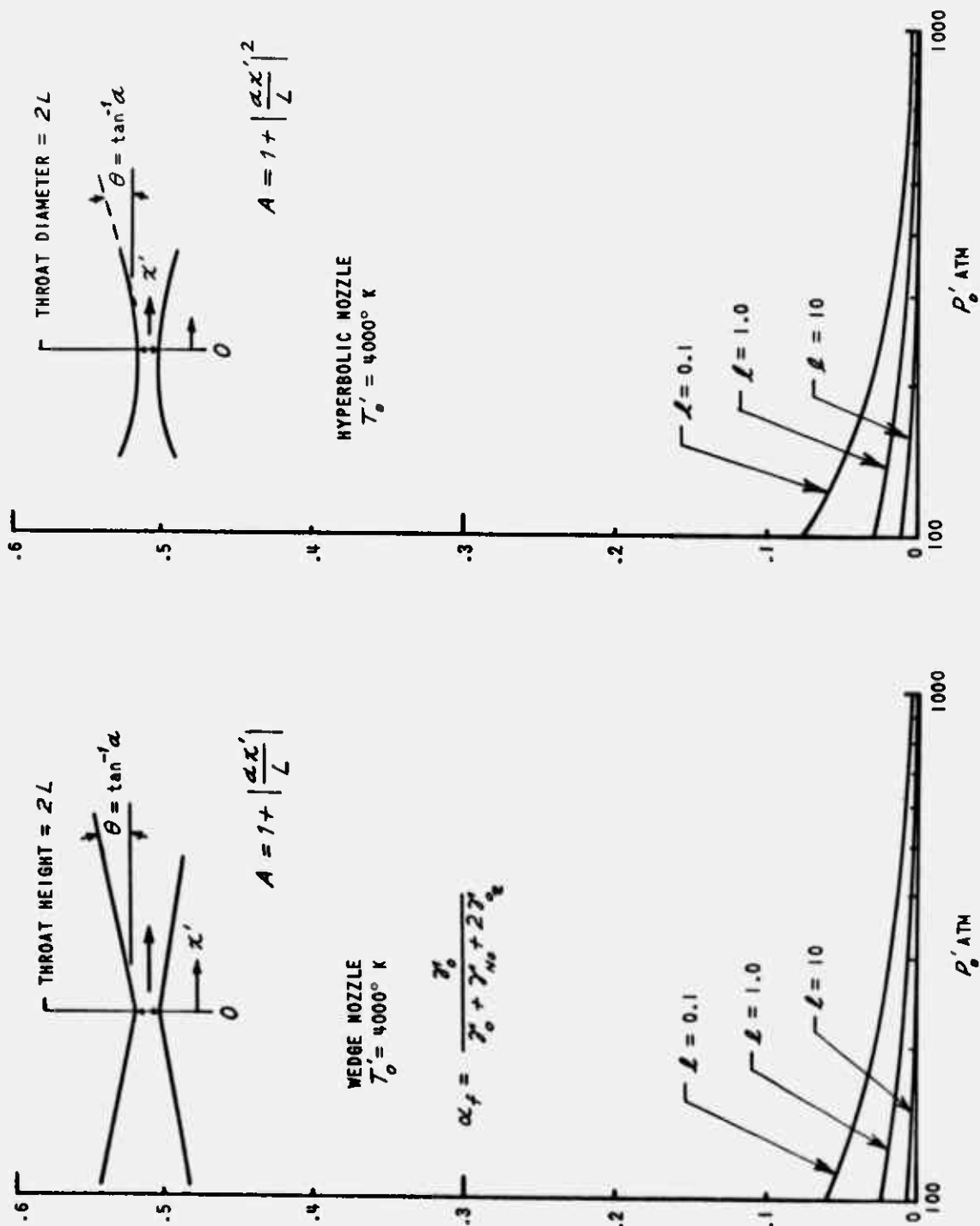


Figure 20 (a) EFFECT OF STAGNATION PRESSURE AND NOZZLE GEOMETRY ON THE FROZEN DEGREE OF OXYGEN DISSOCIATION FOR A SIMPLIFIED AIR MODEL.

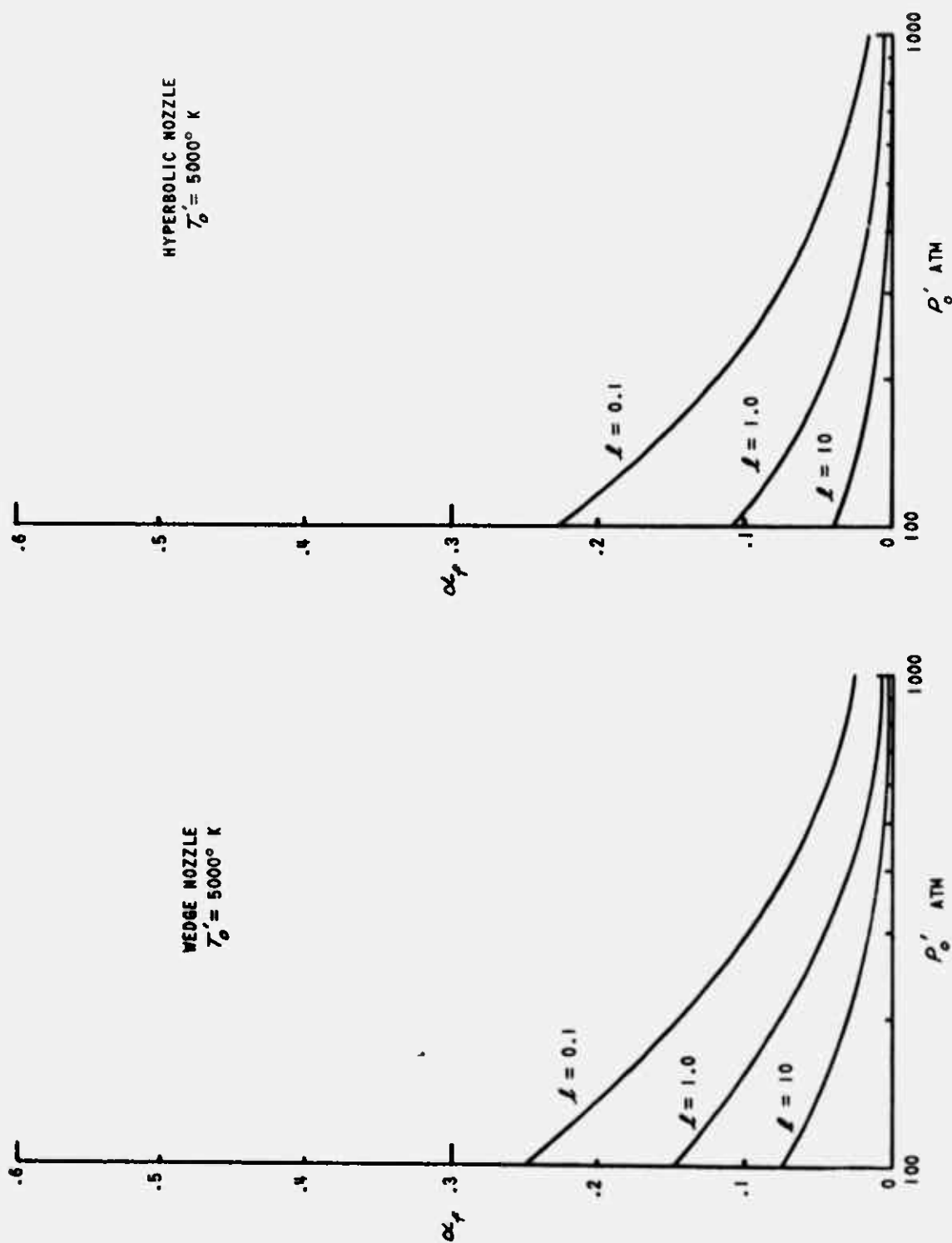


Figure 20 (b) EFFECT OF STAGNATION PRESSURE AND NOZZLE GEOMETRY ON THE FROZEN DEGREE OF OXYGEN DISSOCIATION FOR A SIMPLIFIED AIR MODEL.

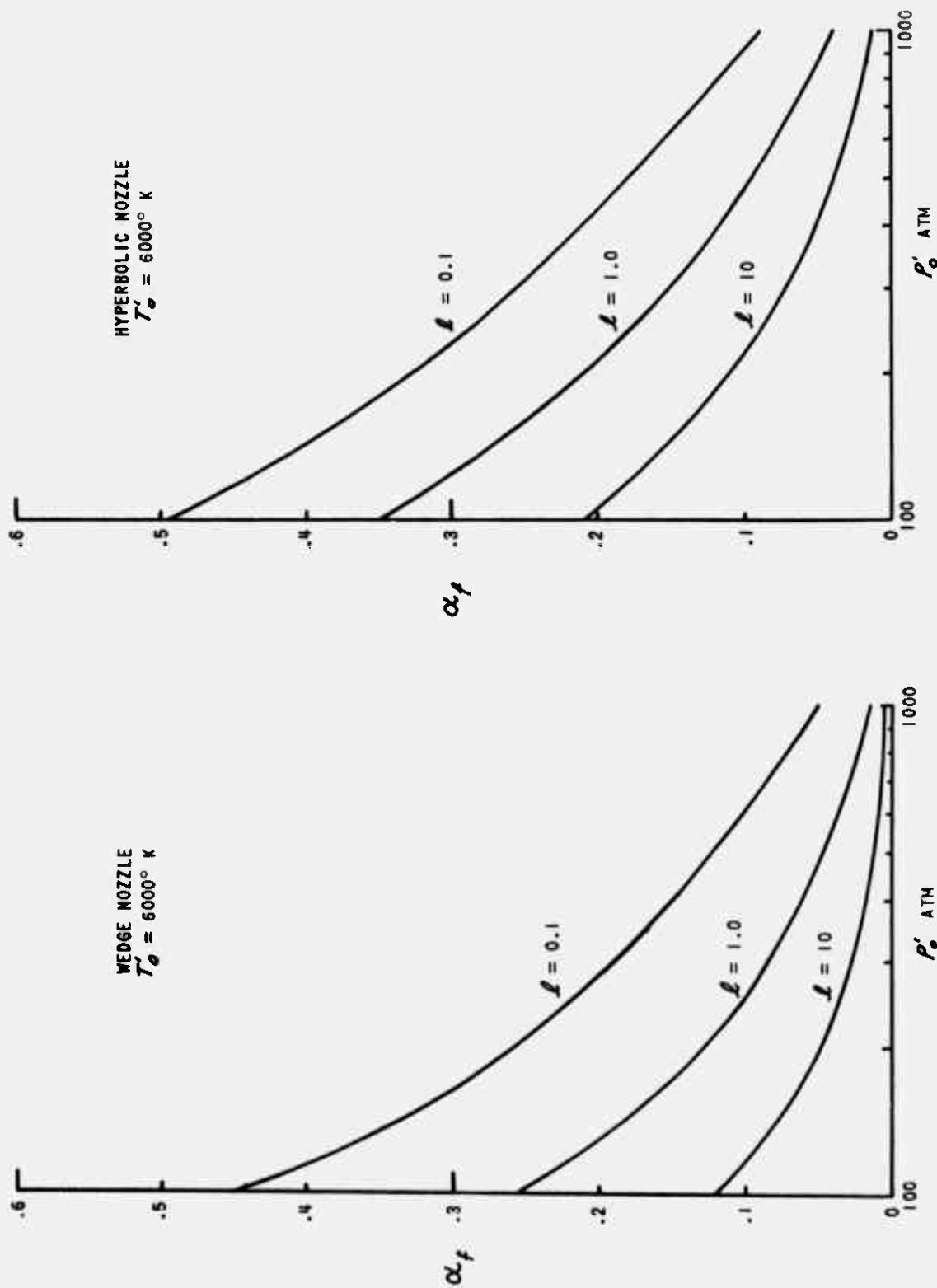


Figure 20 (c) EFFECT OF STAGNATION PRESSURE AND NOZZLE GEOMETRY ON THE FROZEN DEGREE OF OXYGEN DISSOCIATION FOR A SIMPLIFIED AIR MODEL.

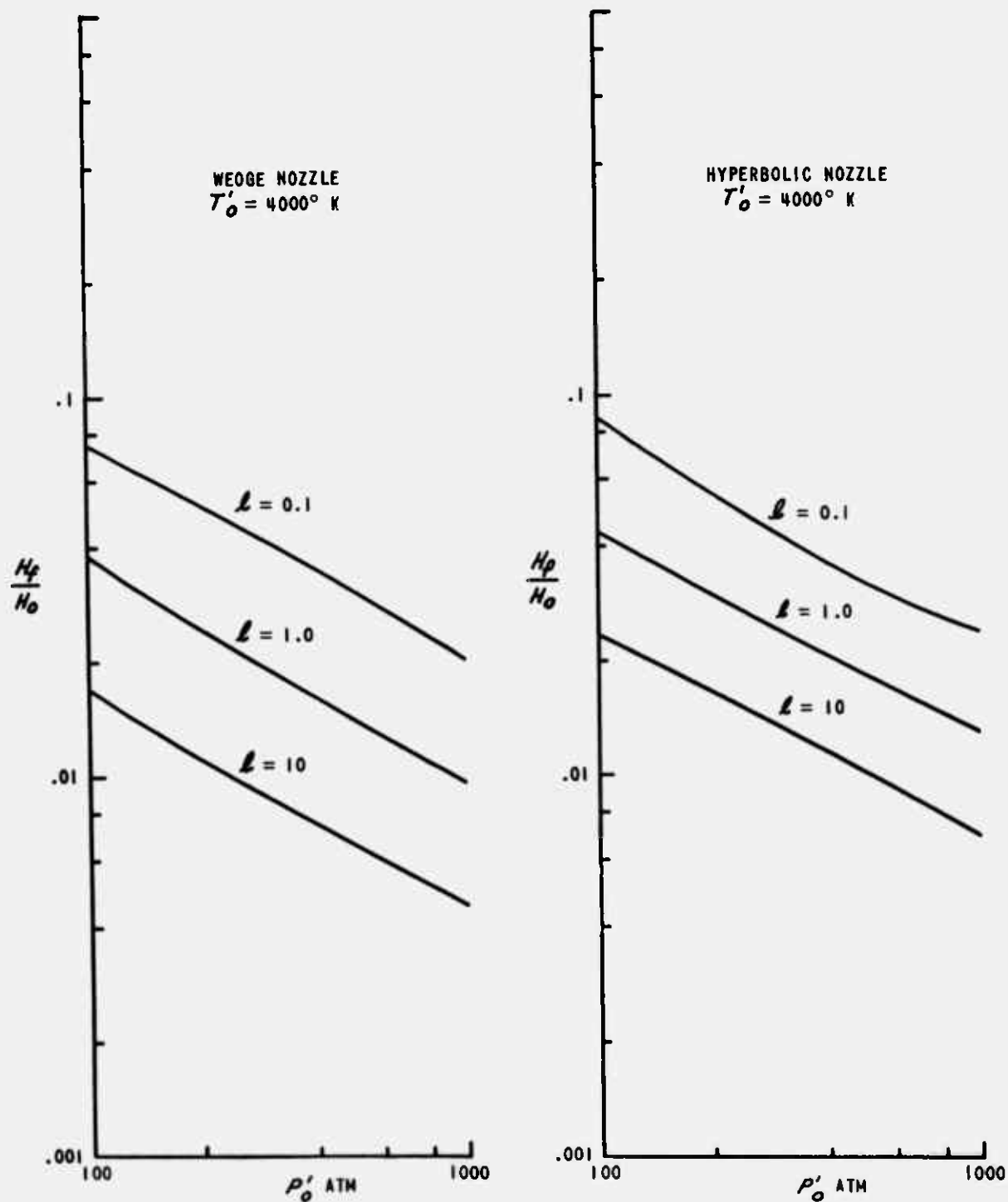


Figure 21 (a) EFFECT OF STAGNATION PRESSURE AND NOZZLE GEOMETRY ON THE FROZEN CHEMICAL ENERGY FOR A SIMPLIFIED AIR MODEL

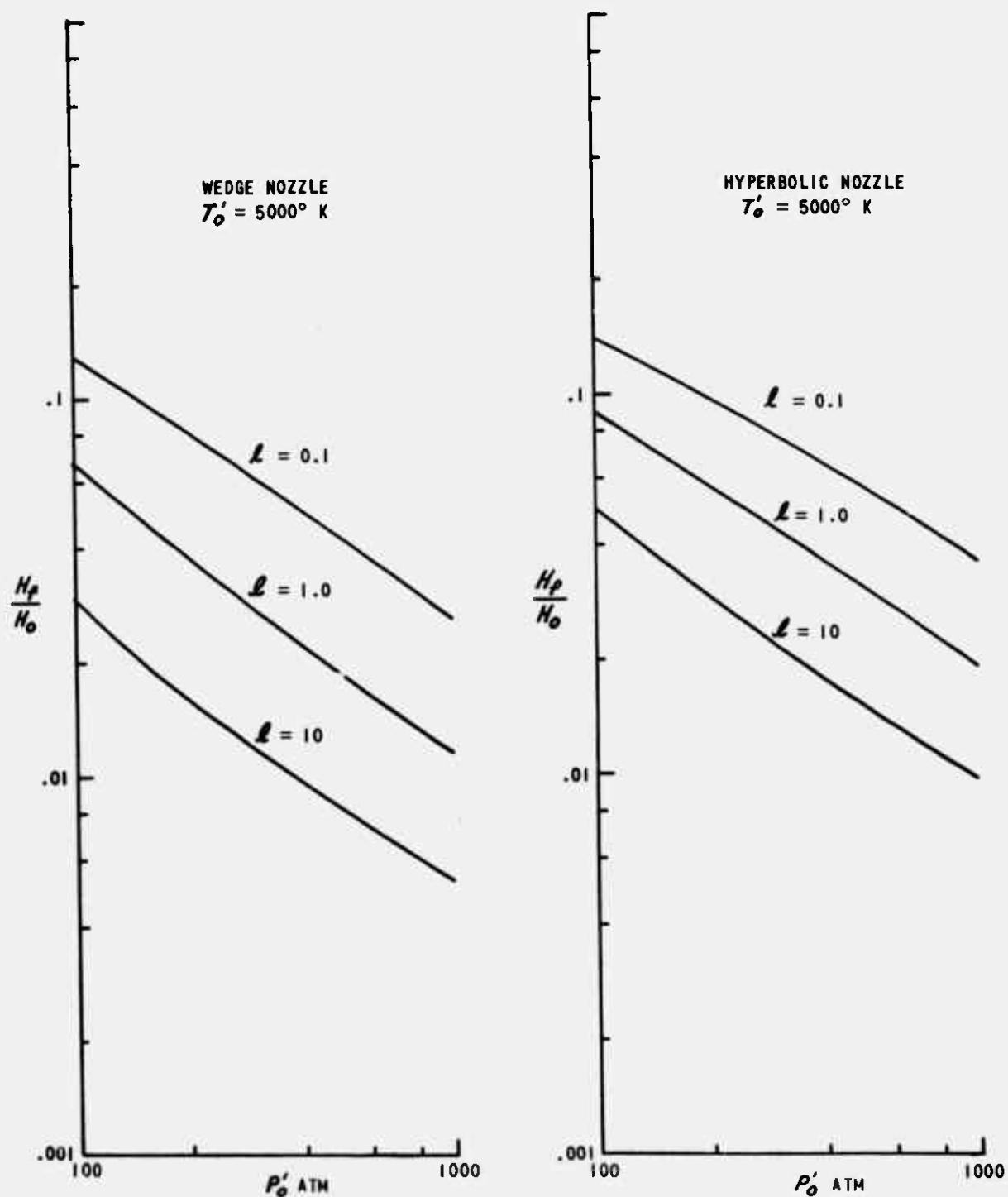


Figure 21 (b) EFFECT OF STAGNATION PRESSURE AND NOZZLE GEOMETRY ON THE FROZEN CHEMICAL ENERGY FOR A SIMPLIFIED AIR MODEL

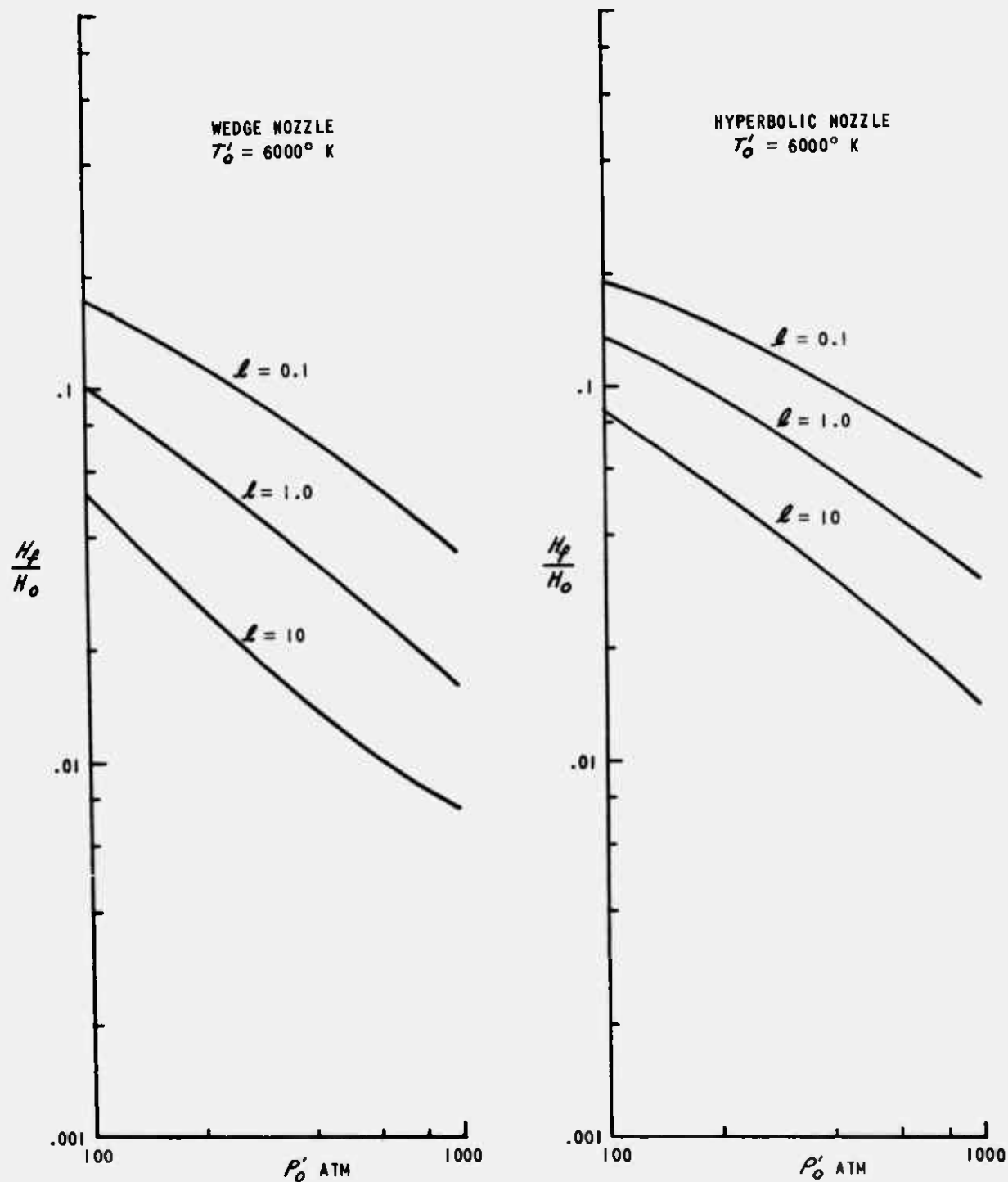


Figure 21 (c) EFFECT OF STAGNATION PRESSURE AND NOZZLE GEOMETRY ON THE FROZEN CHEMICAL ENERGY FOR A SIMPLIFIED AIR MODEL

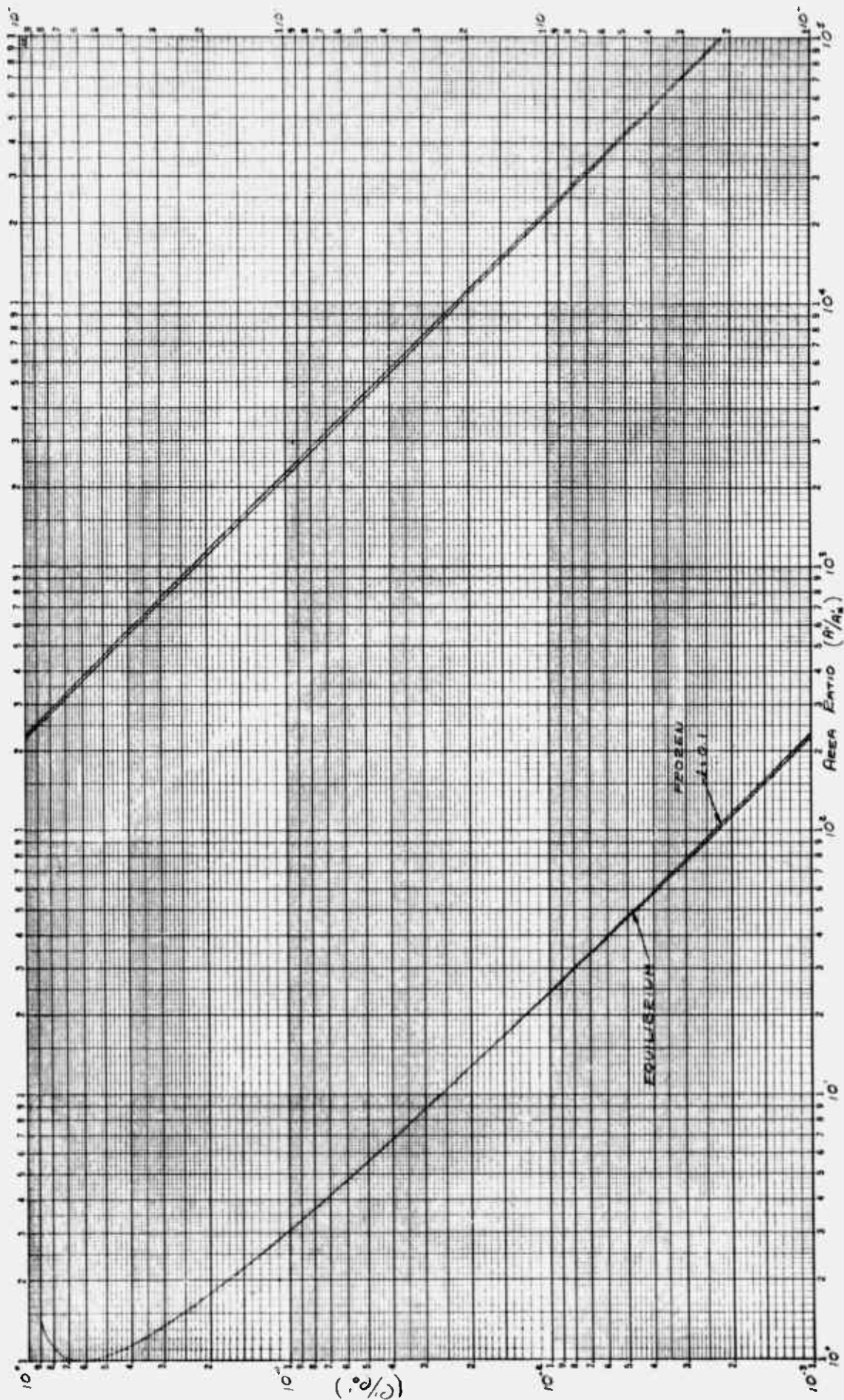


FIGURE NO 20.2 HYPERBOLIC NOZZLE $T_0' = 4000^\circ K$ $P_0' = 100 \text{ ATM}$

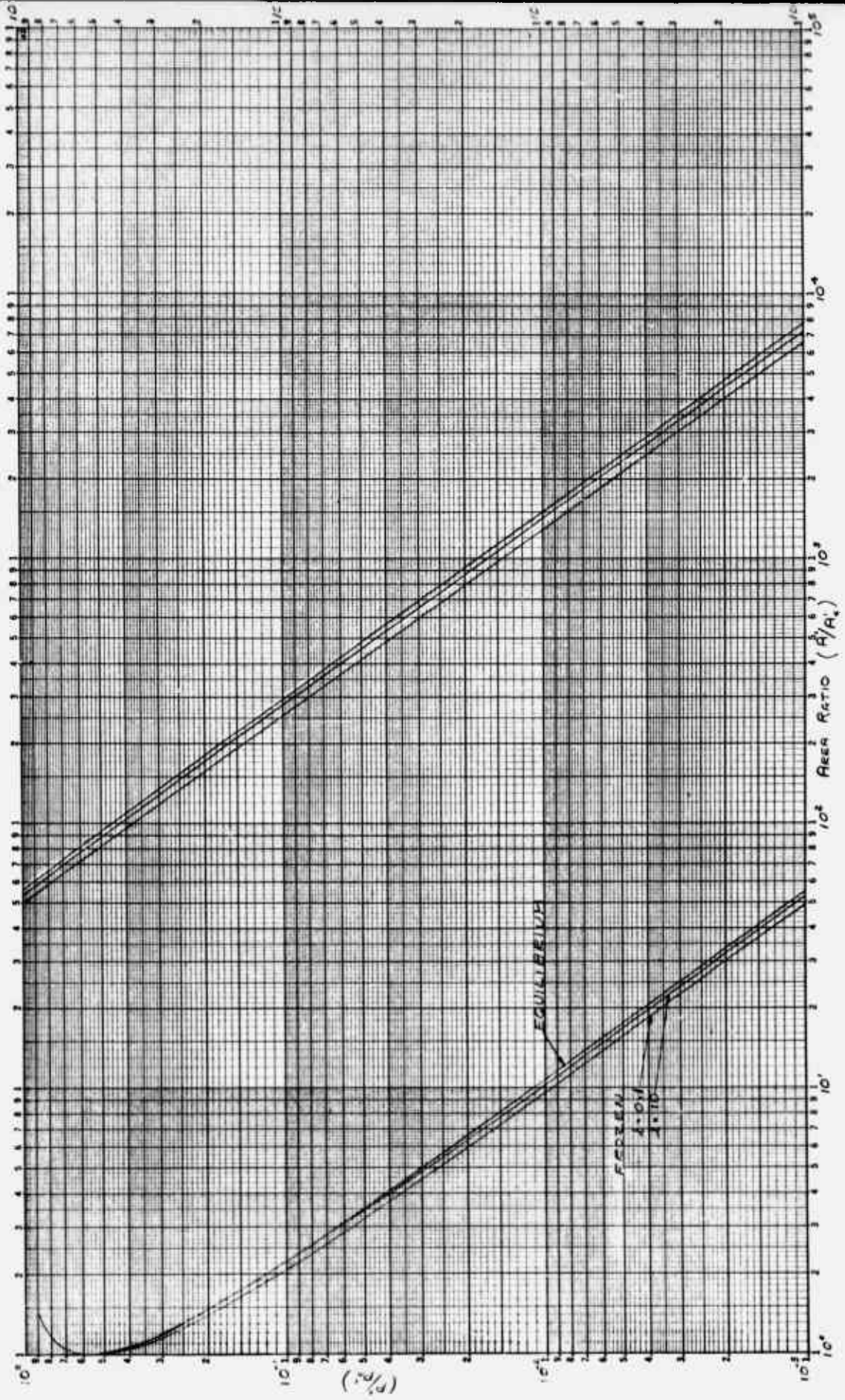


FIGURE NO. 22.2 HYPERBOLIC NOZZLE $T_0' = 4000^\circ \text{K}$ $P_0' = 100 \text{ A-M}$

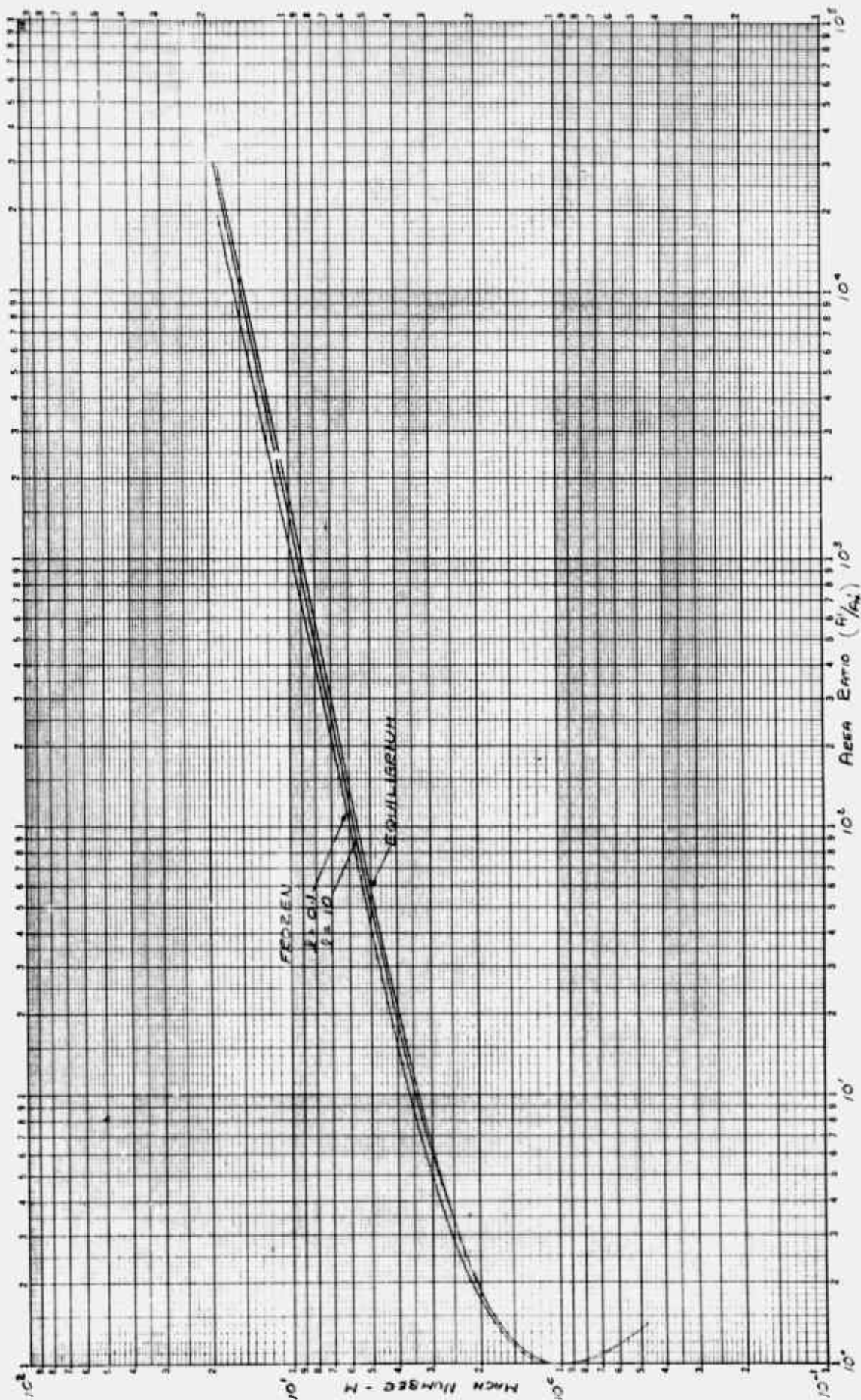


FIGURE 100 IS: HYPERBOLIC NOZZLE $T_0' = 4000^\circ K$ $P_0' = 100 \text{ ATM}$

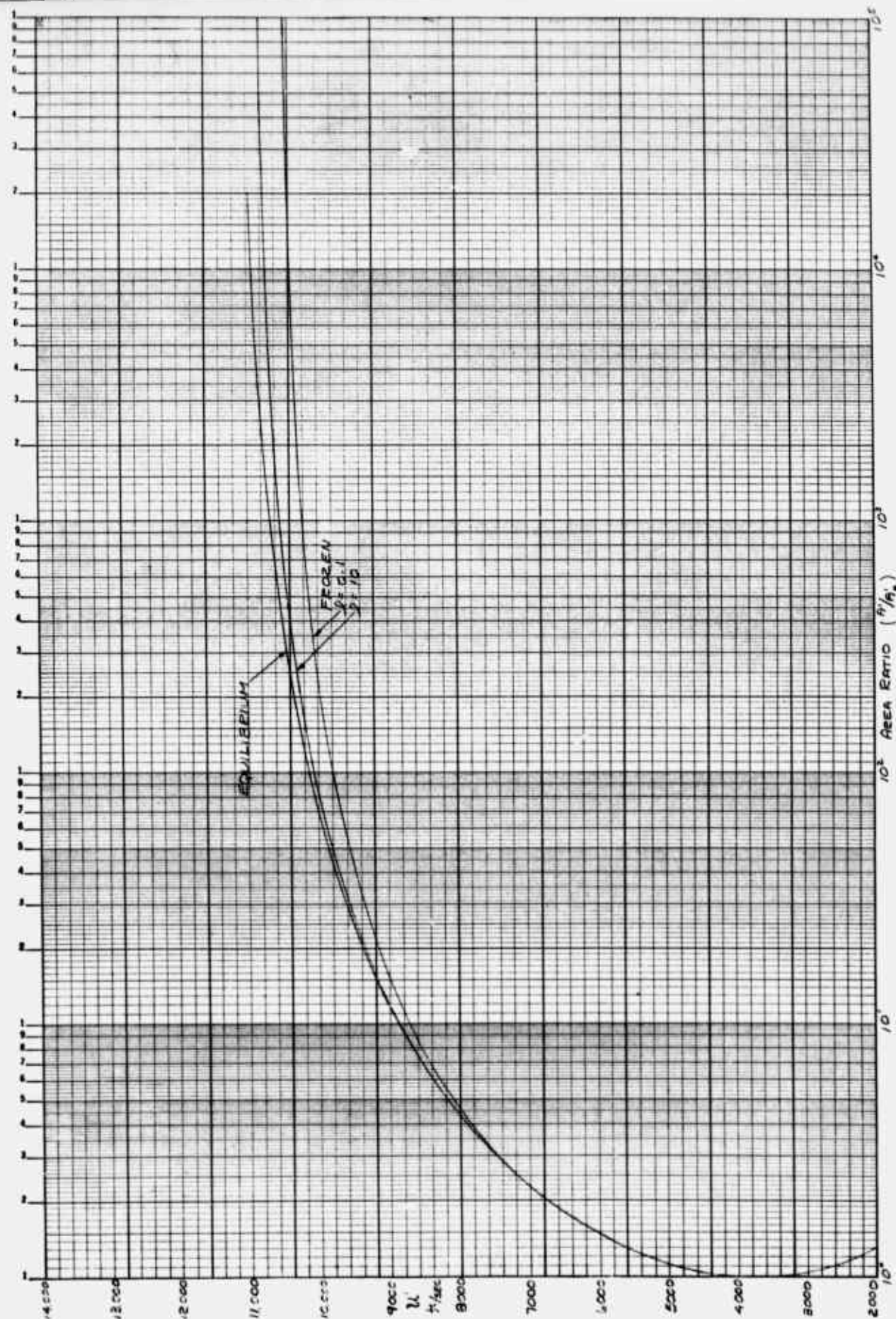


FIGURE NO. 22d HYPERBOLIC NOZZLE $T_0 = 4000^\circ K$ $P_0 = 100 \text{ ATM}$

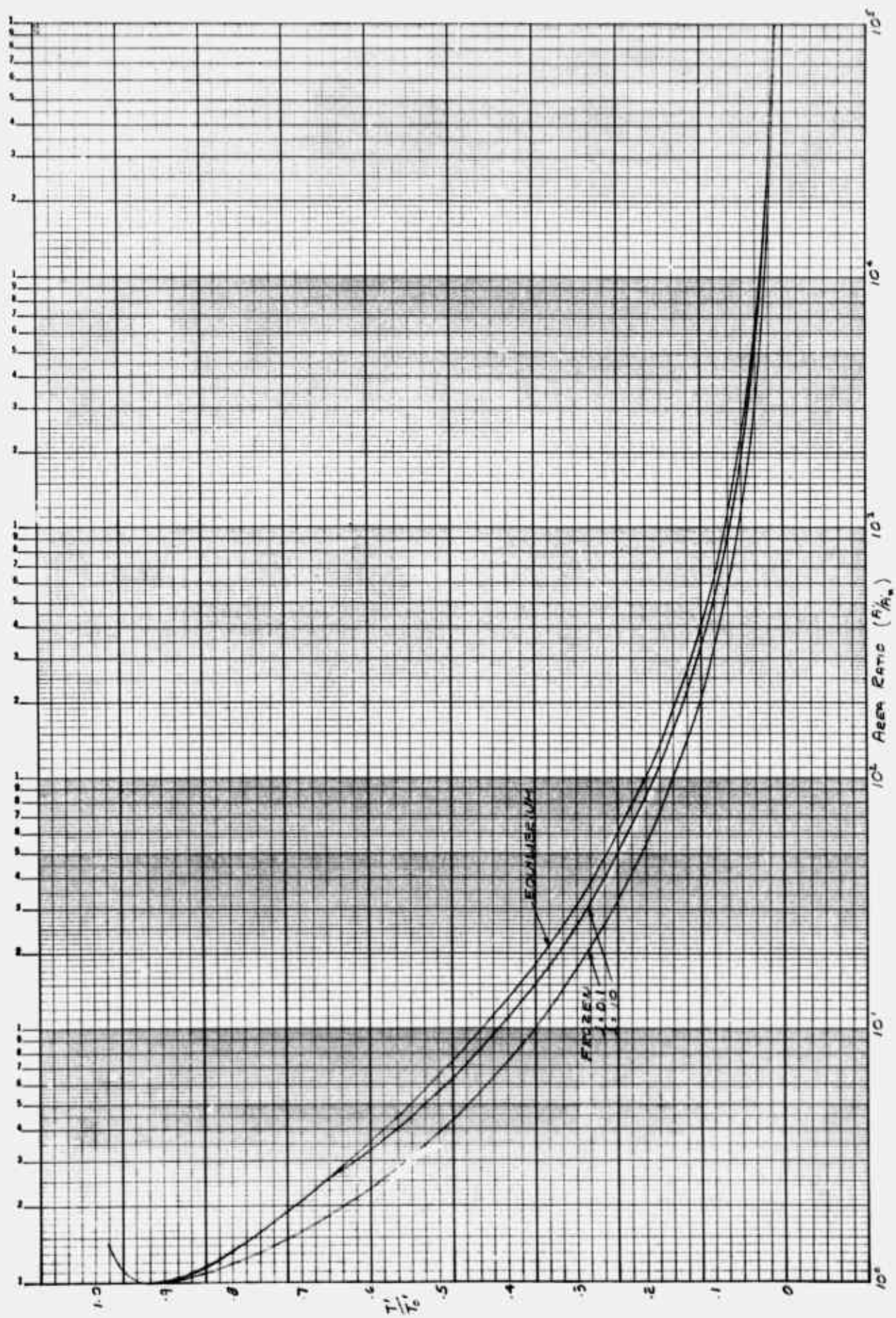


FIGURE NO. 22c HYPERBOLIC NOZZLE $T_0 = 4000^\circ K$ $P_0 = 100 \text{ ATM}$

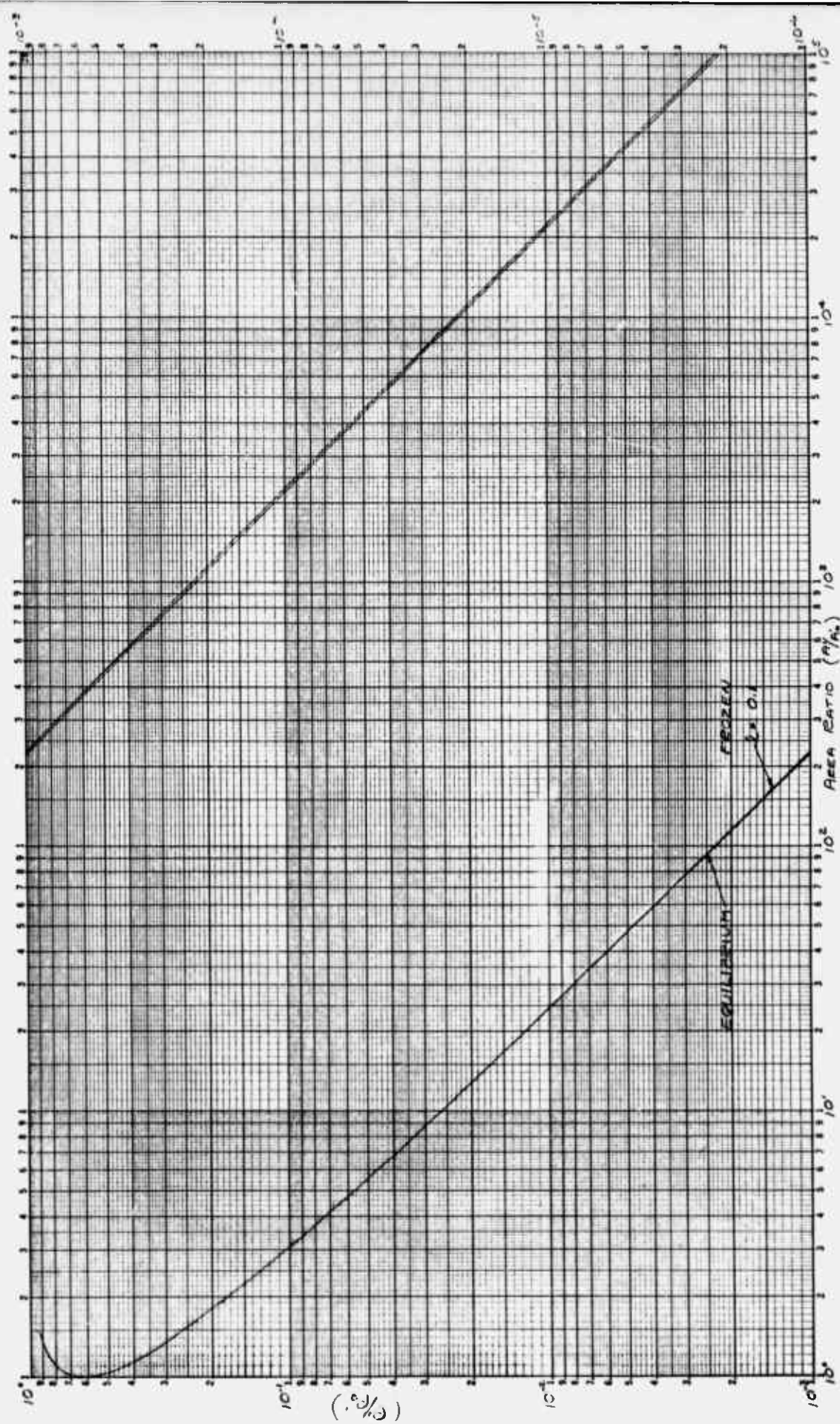


FIGURE NO 292 HYPERBOLIC NOZZLE $T_0 = 4000^\circ\text{K}$ $P_0 = 300\text{ATM}$

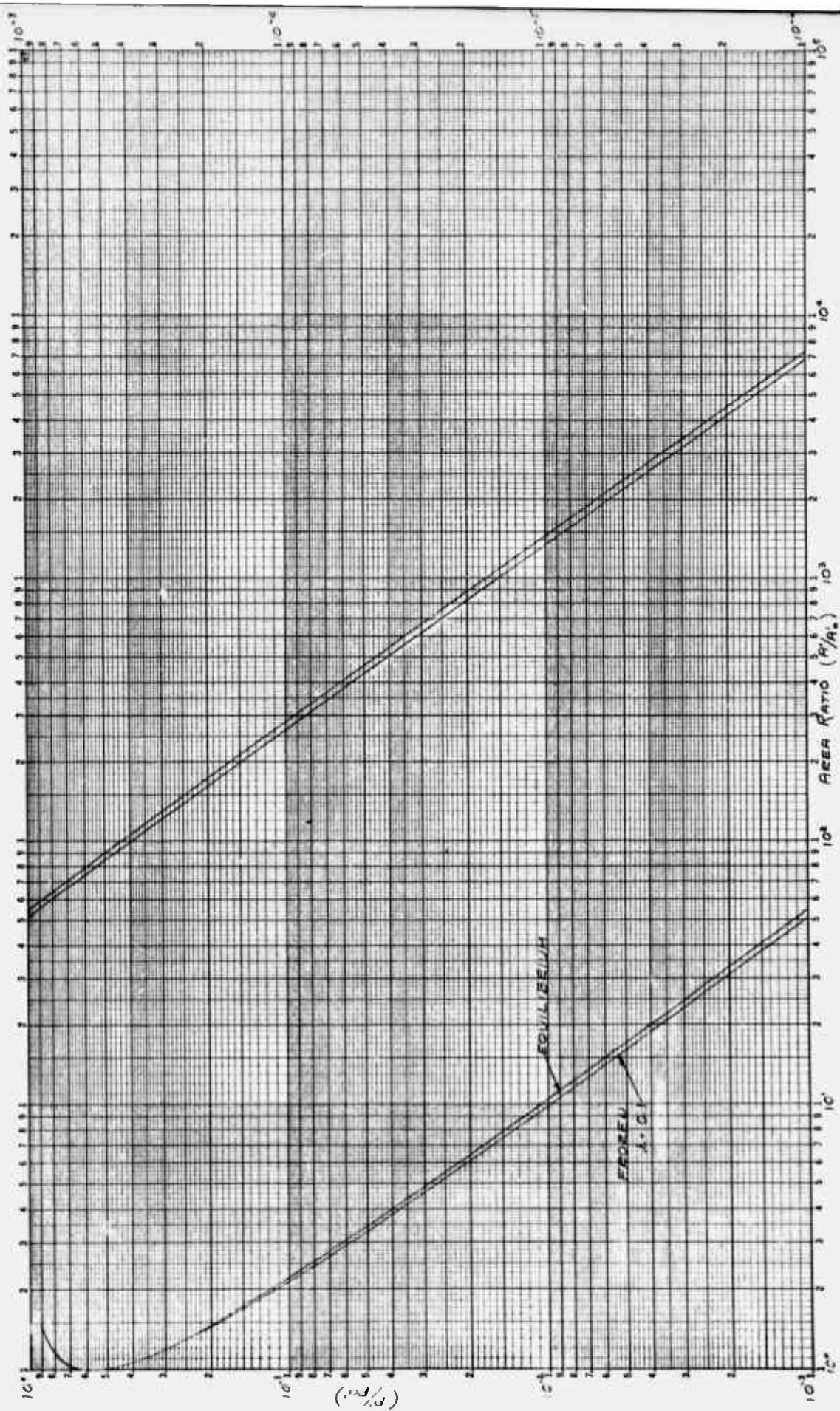


FIGURE NO. 23b HYPERBOLIC NOZZLE $T_i = 4000^\circ K$ $P_0 = 300$ ATM

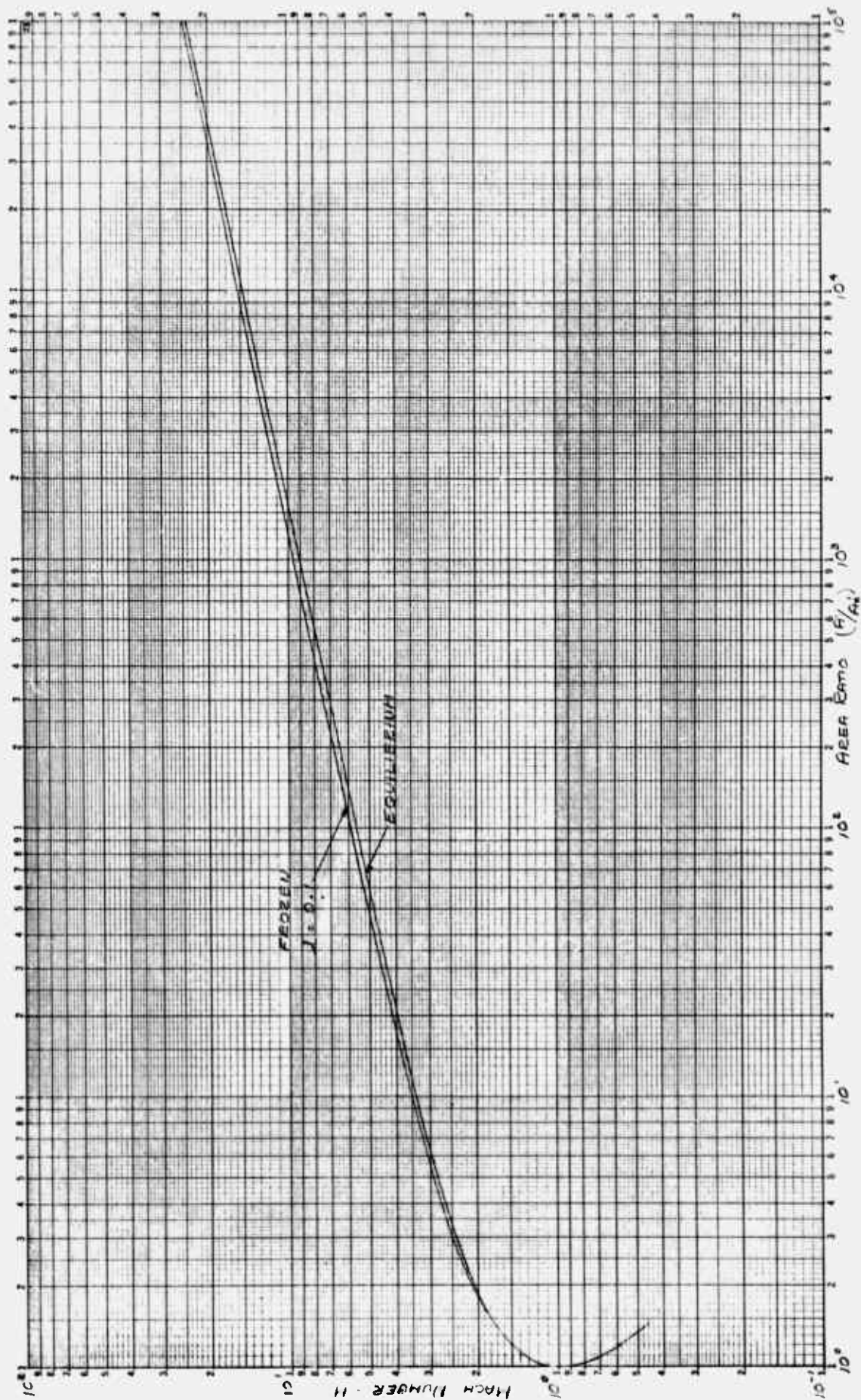


FIGURE 10. 25c HYPERBOLIC NOZZLE $T_0' = 4000^\circ K$ $P_0' = 300 \text{ ATM}$

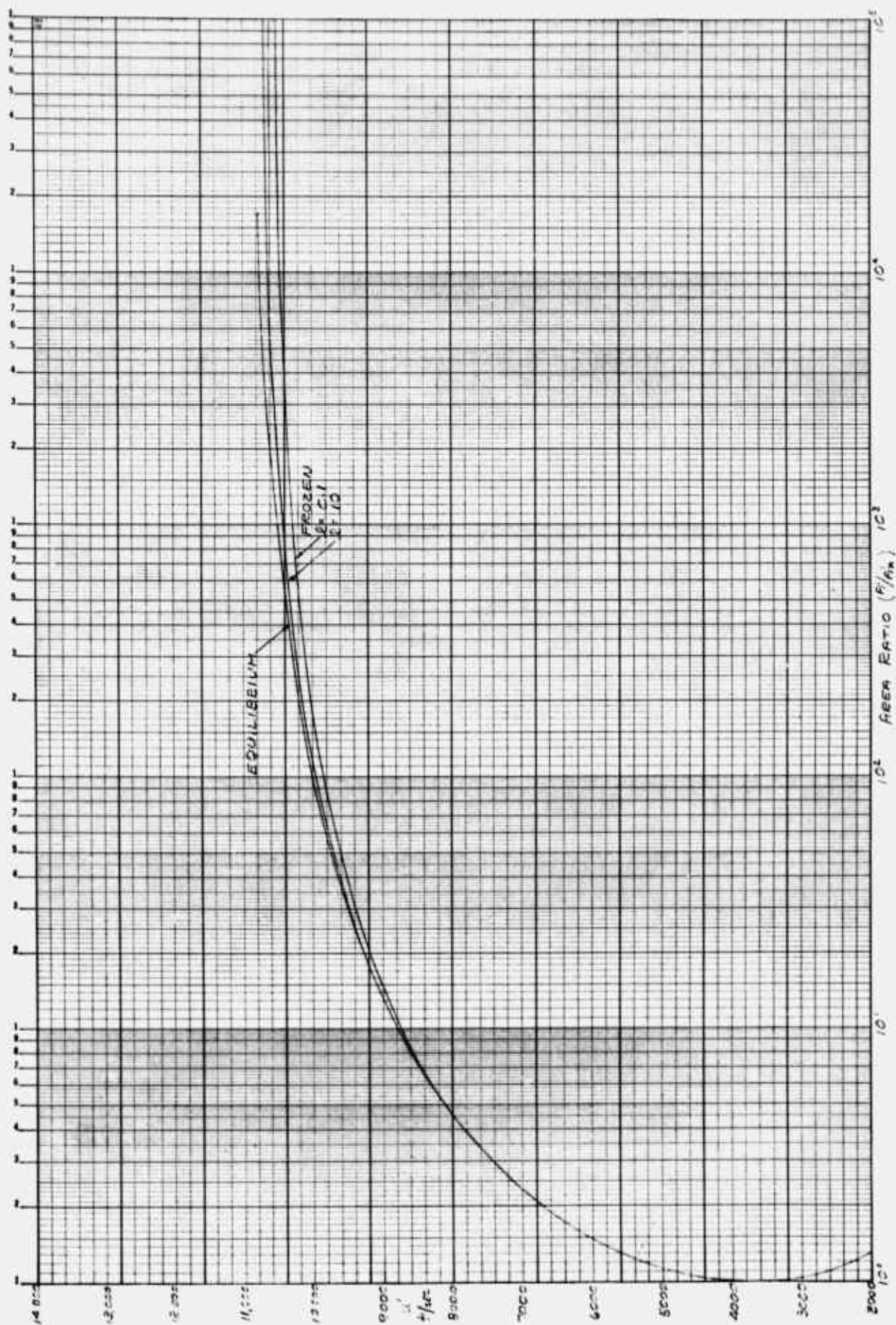


FIGURE NO. 23J NOZZLE $T_0 = 4000^\circ K$ $P_0 = 300 \text{ ATM}$

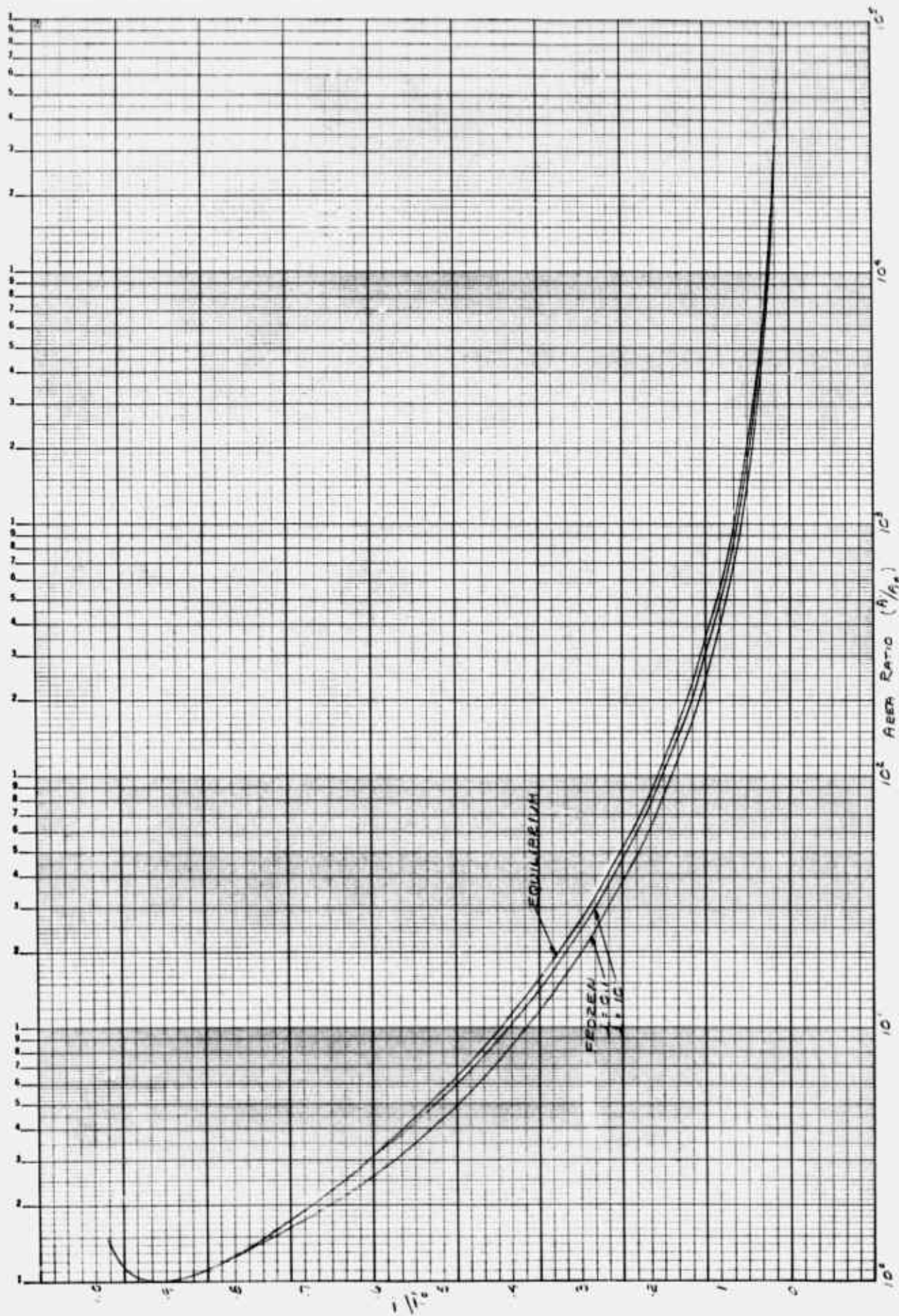


FIGURE NO. 23e HYPERBOLIC NOZZLE $T_0 = 4000^\circ K$ $P_0' = 300 \text{ ATM}$

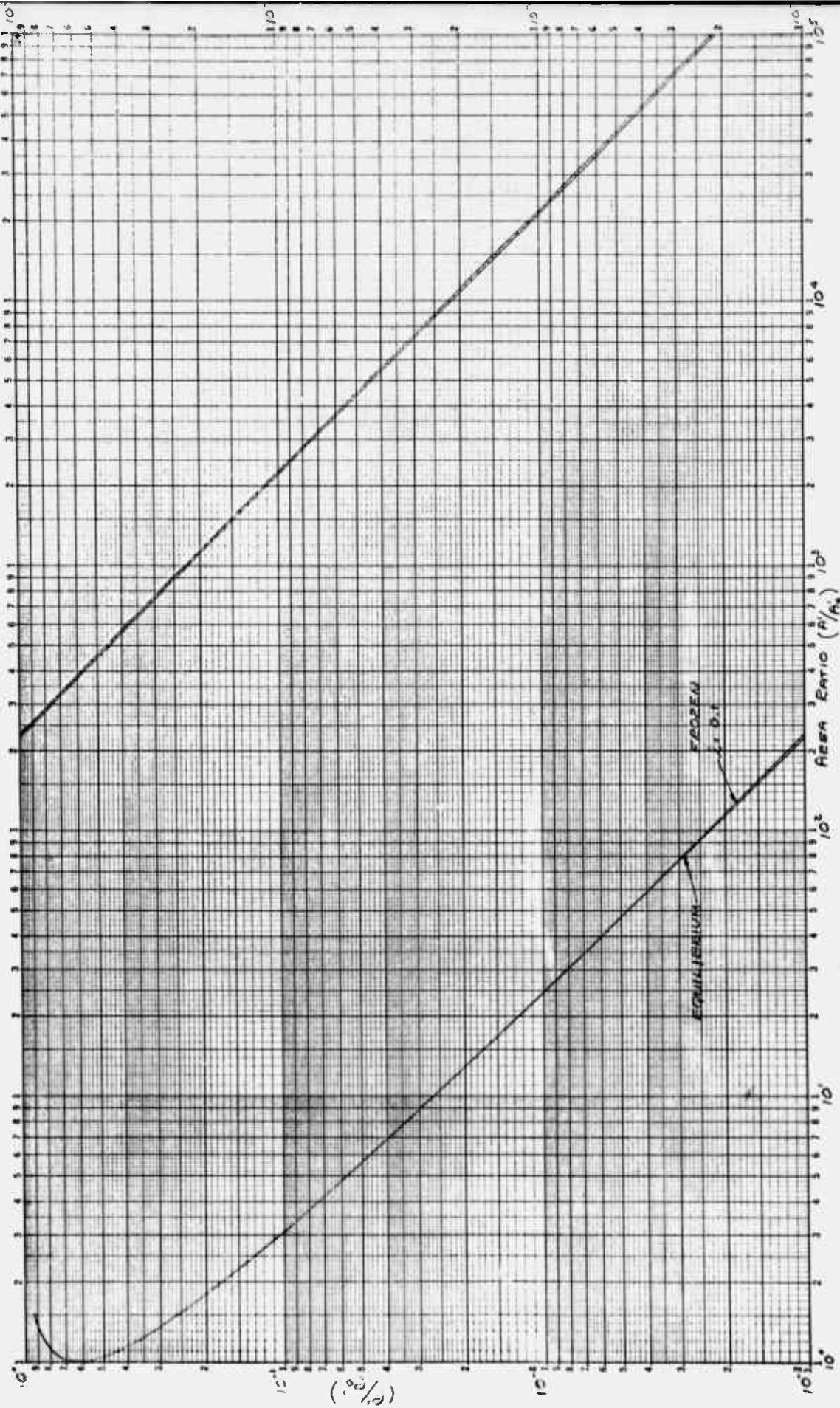


FIGURE NO. 24.2. HYPERBOLIC NOZZLE $T_0' = 4000^\circ K$ $P_0' = 1000$ ATM

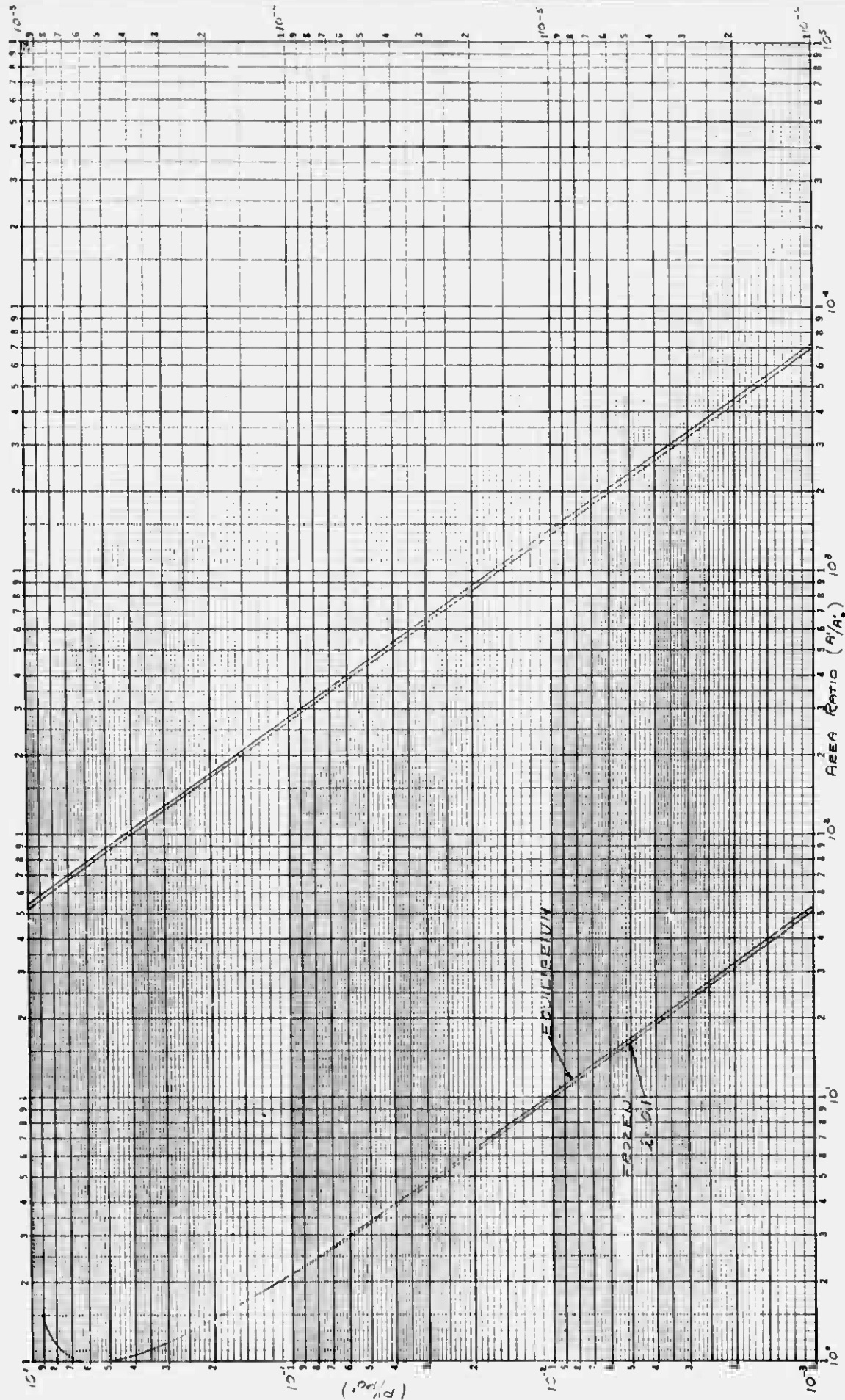
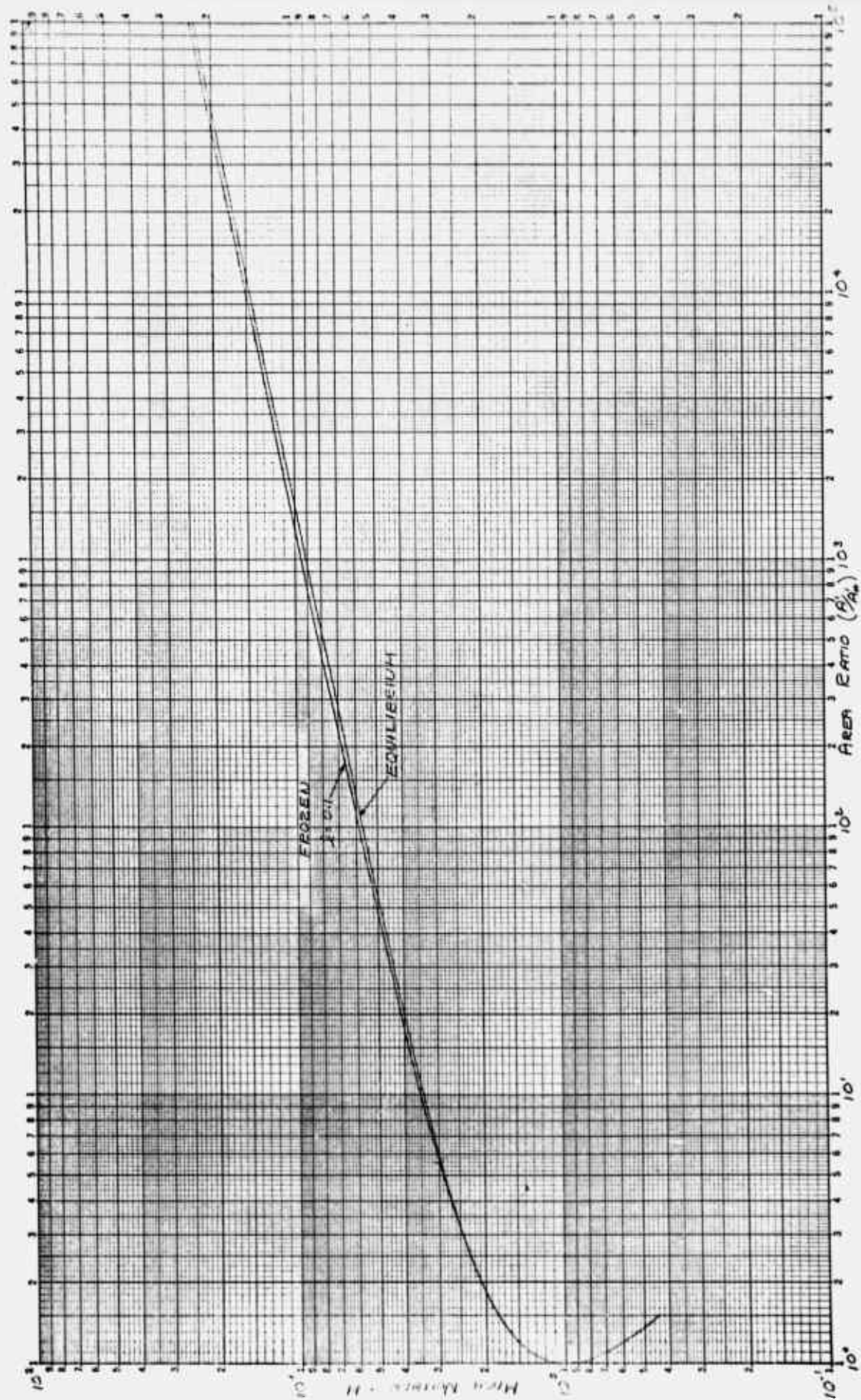


FIGURE NO. 246 HYPERBOLIC NOZZLE $T'_0 = 4000^\circ K$ $P'_0 = 1000 \text{ ATM}$



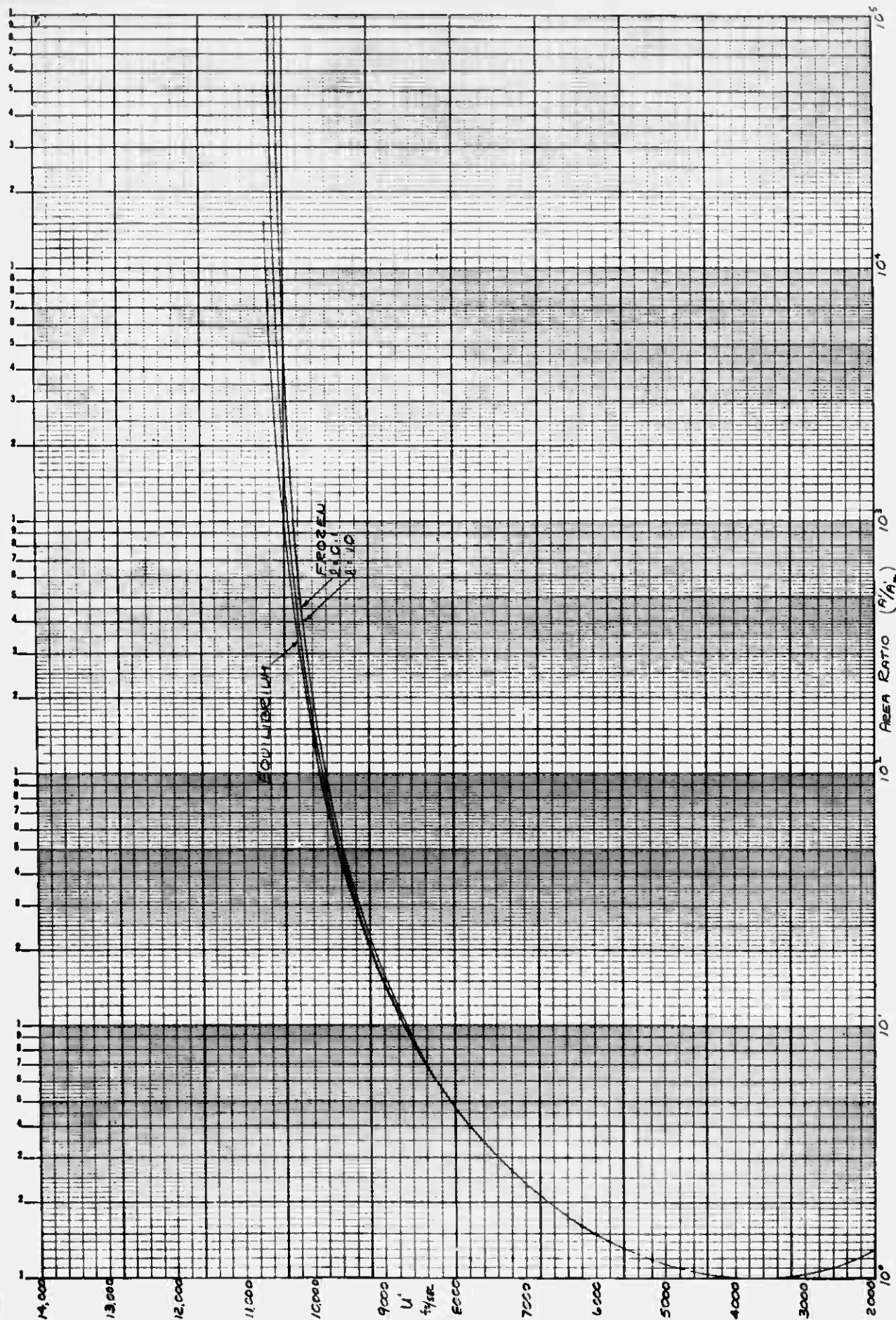


FIGURE NO. 24d HYPERBOLIC NOZZLE $T_0' = 4000^\circ K$ $P_0' = 1000 \text{ ATM}$

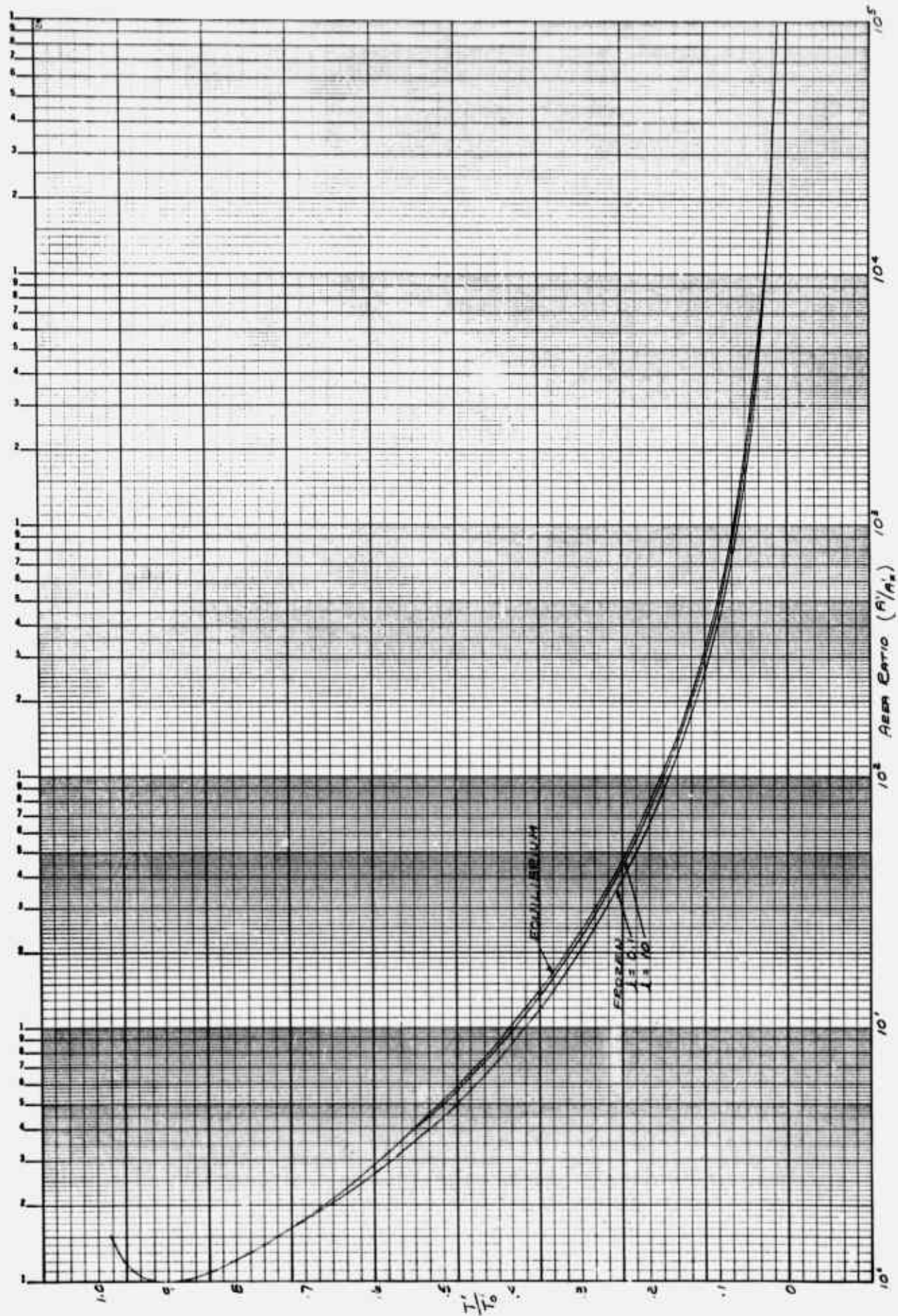


FIGURE NO. 24c HYPERSONIC NOZZLE $T'_0 = 4000^\circ K$ $P'_0 = 1000 \text{ ATM}$

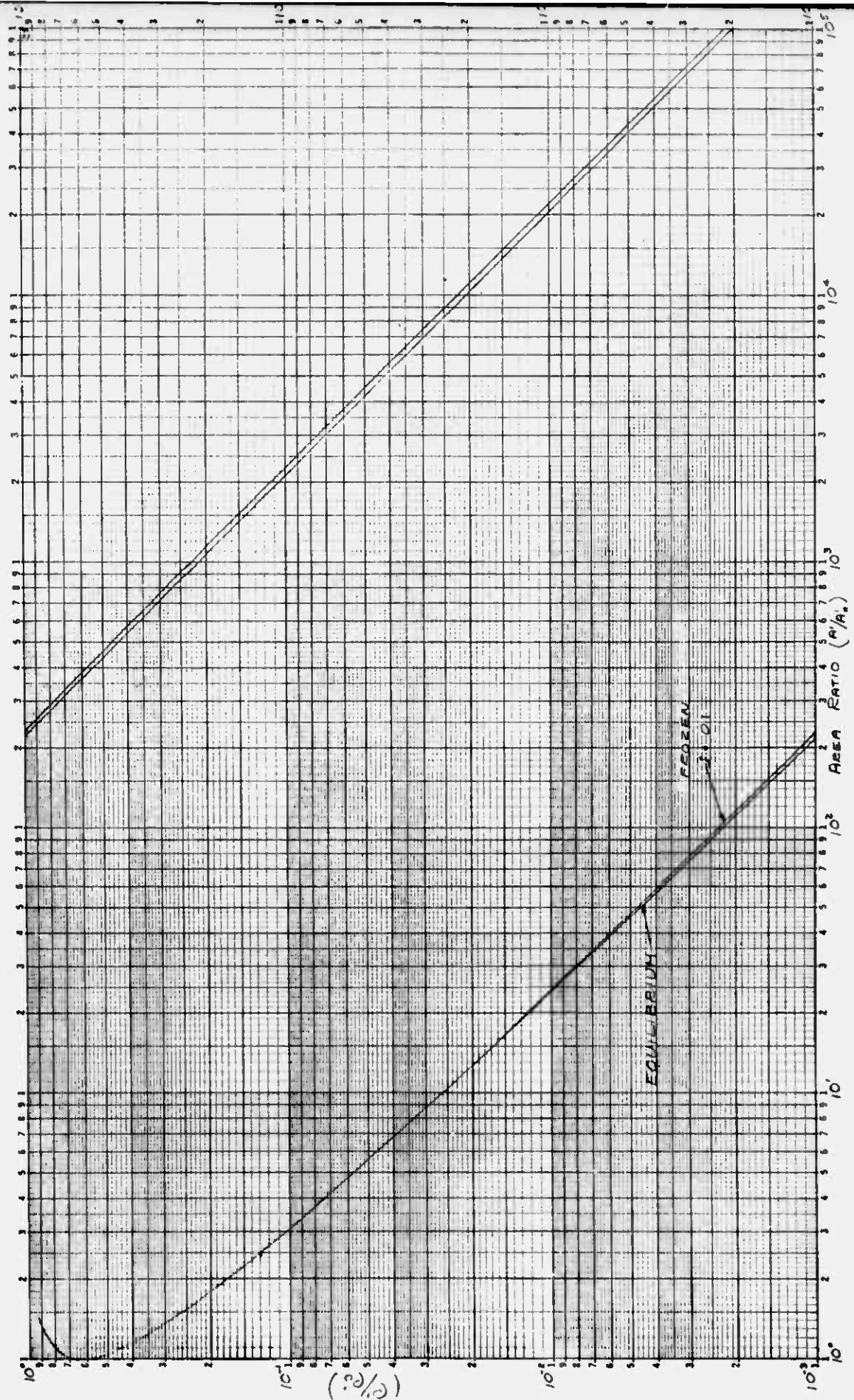


FIGURE NO. 25. HYPERBOLIC NOZZLE $T_0 = 5000^\circ\text{K}$ $P_0 = 100 \text{ ATM}$

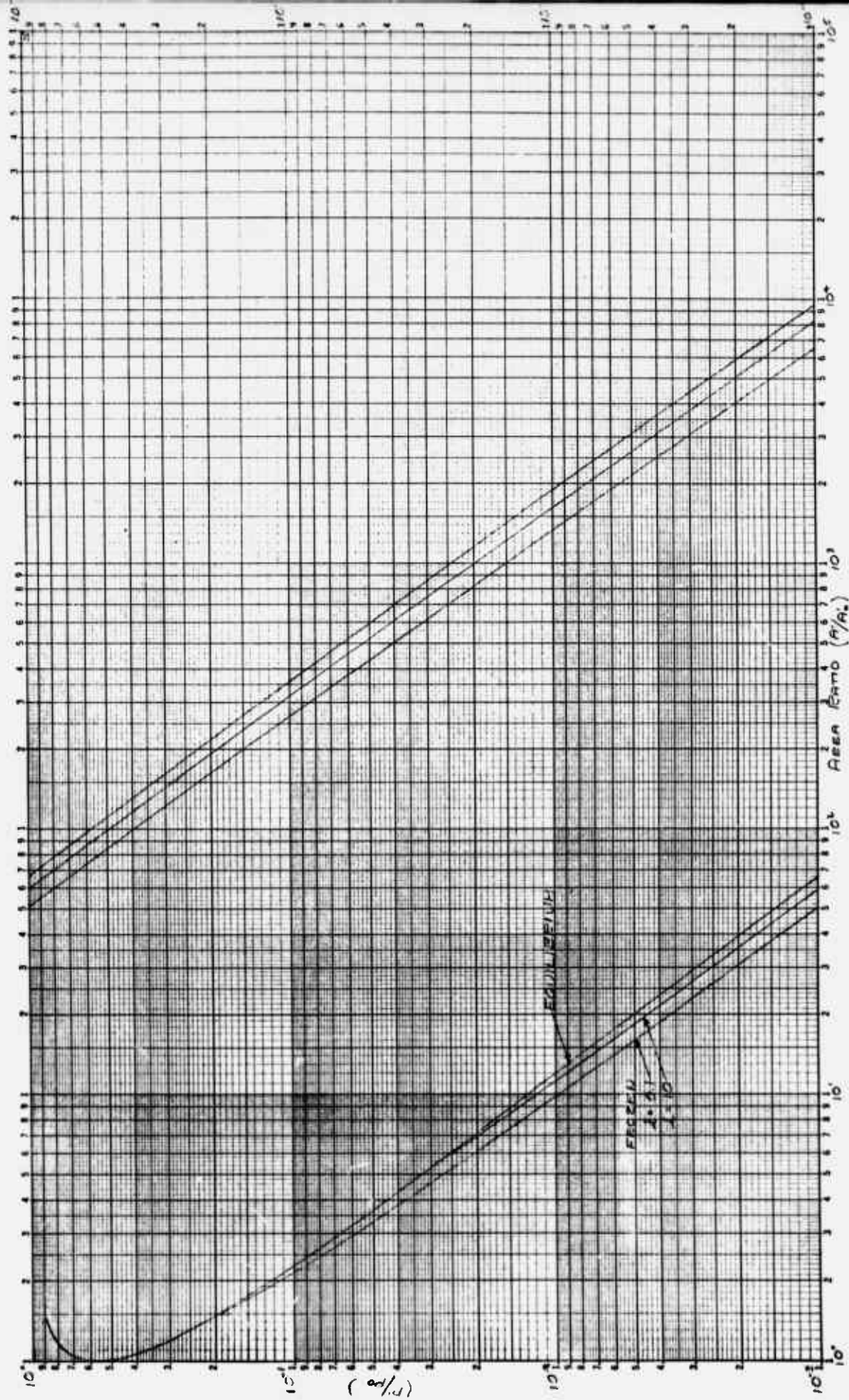


FIGURE NO. 25b HYPERBOLIC NOZZLE $T'_0 = 5000^\circ K$ $P'_0 = 100 \text{ ATM}$

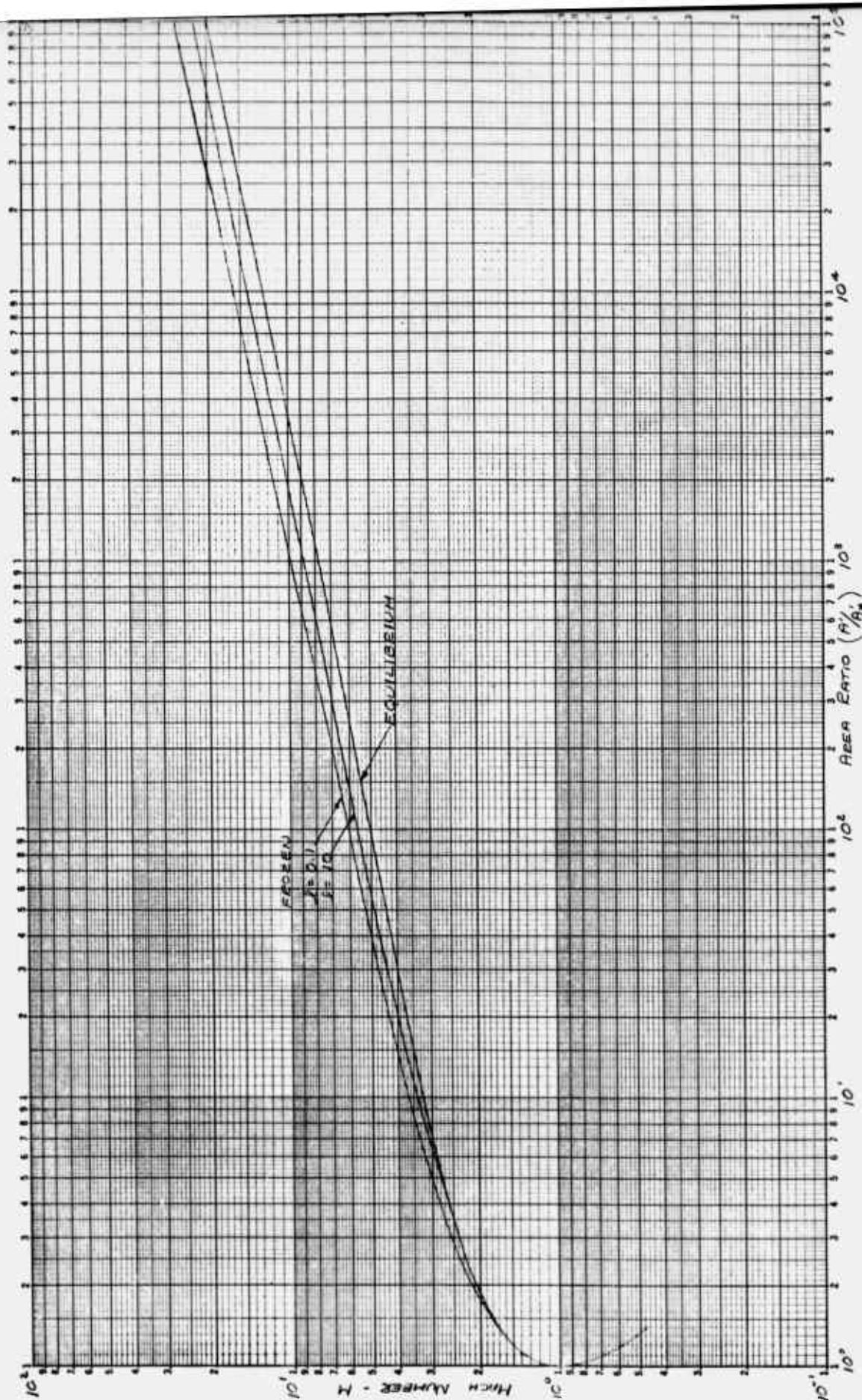


FIGURE NO. 25c HYPERBOLIC NOZZLE $T_0' = 5000^\circ K$ $P_0' = 100 \text{ ATM}$

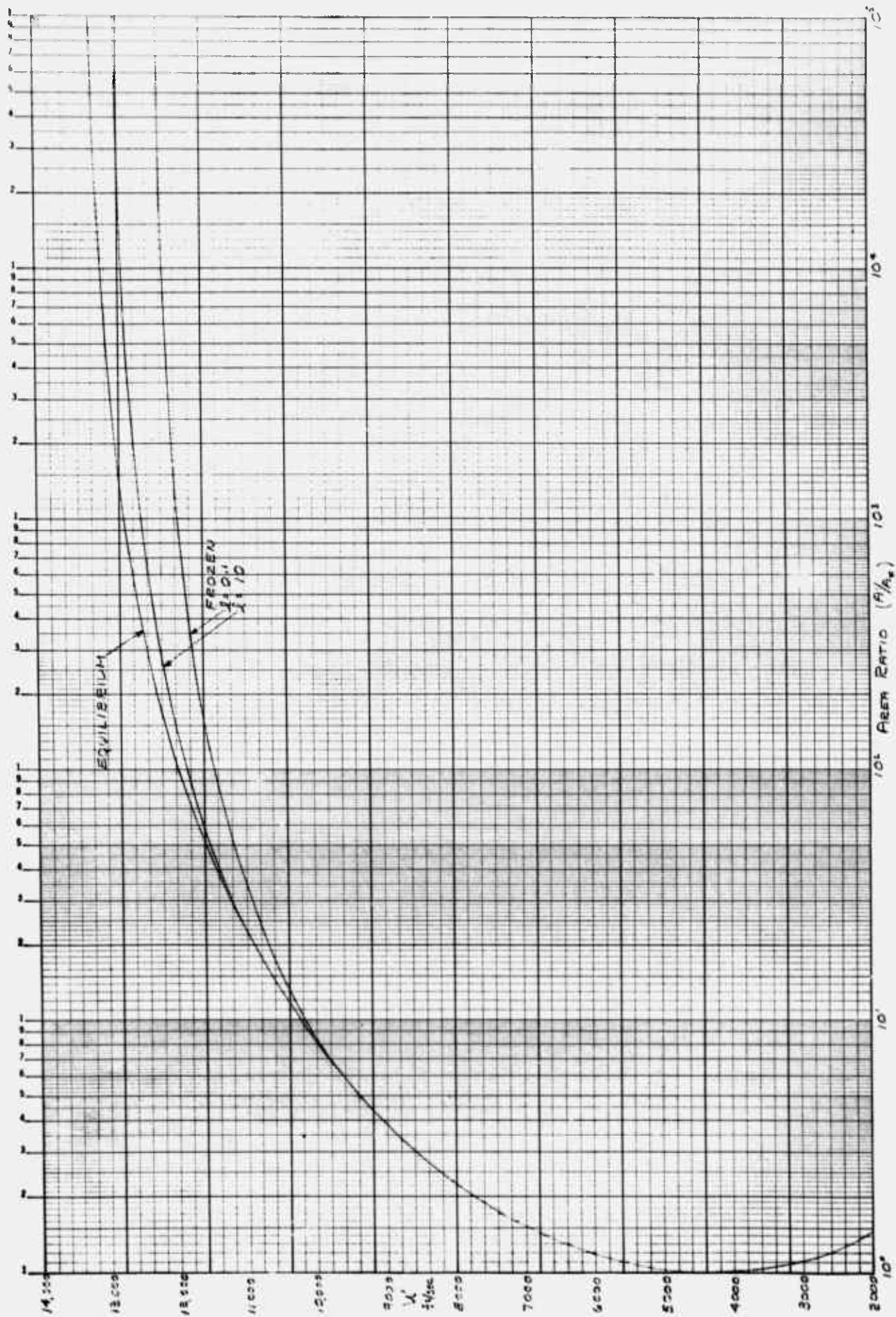


FIGURE NO. 25d HYPERBOLIC LOG $T_0 = 5000 K$

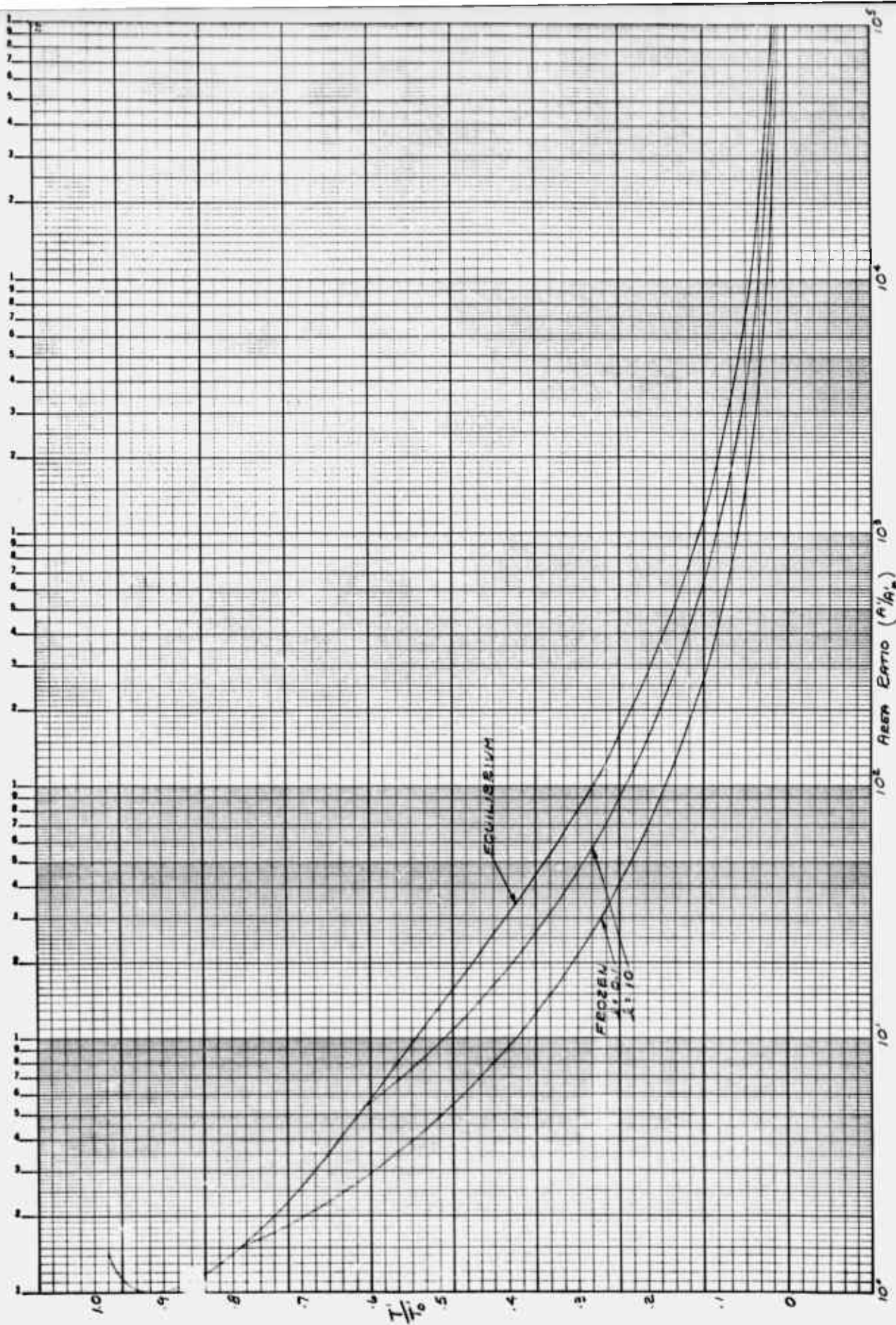


FIGURE NO. 25e HYPERBOLIC NOZZLE $T_0 = 5000^\circ K$ $P_0 = 100 \text{ ATM}$

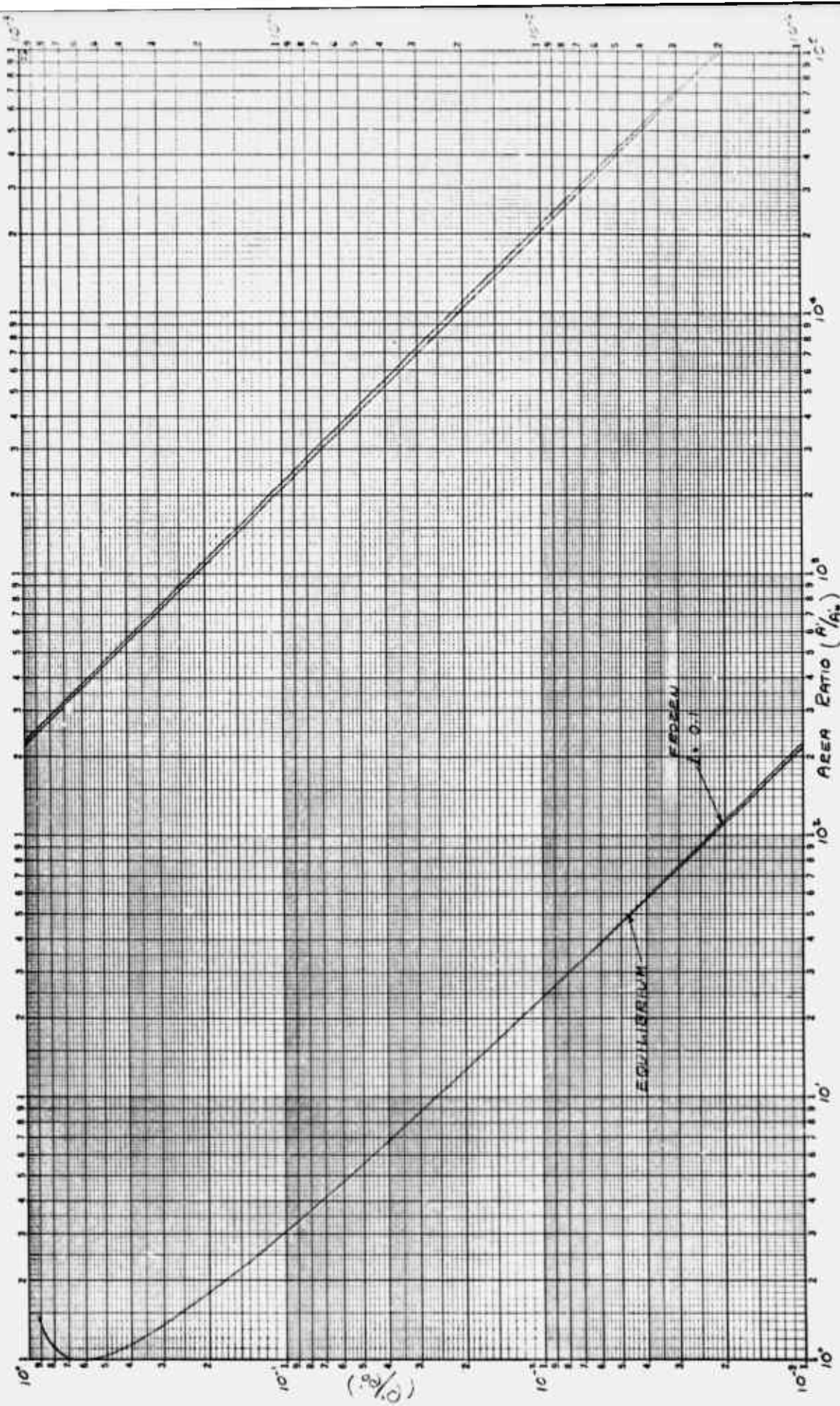


FIGURE NO. 26a HYPERSONIC NOZZLE $T_1 = 5000^\circ K$ $P_1 = 200 \text{ ATM}$

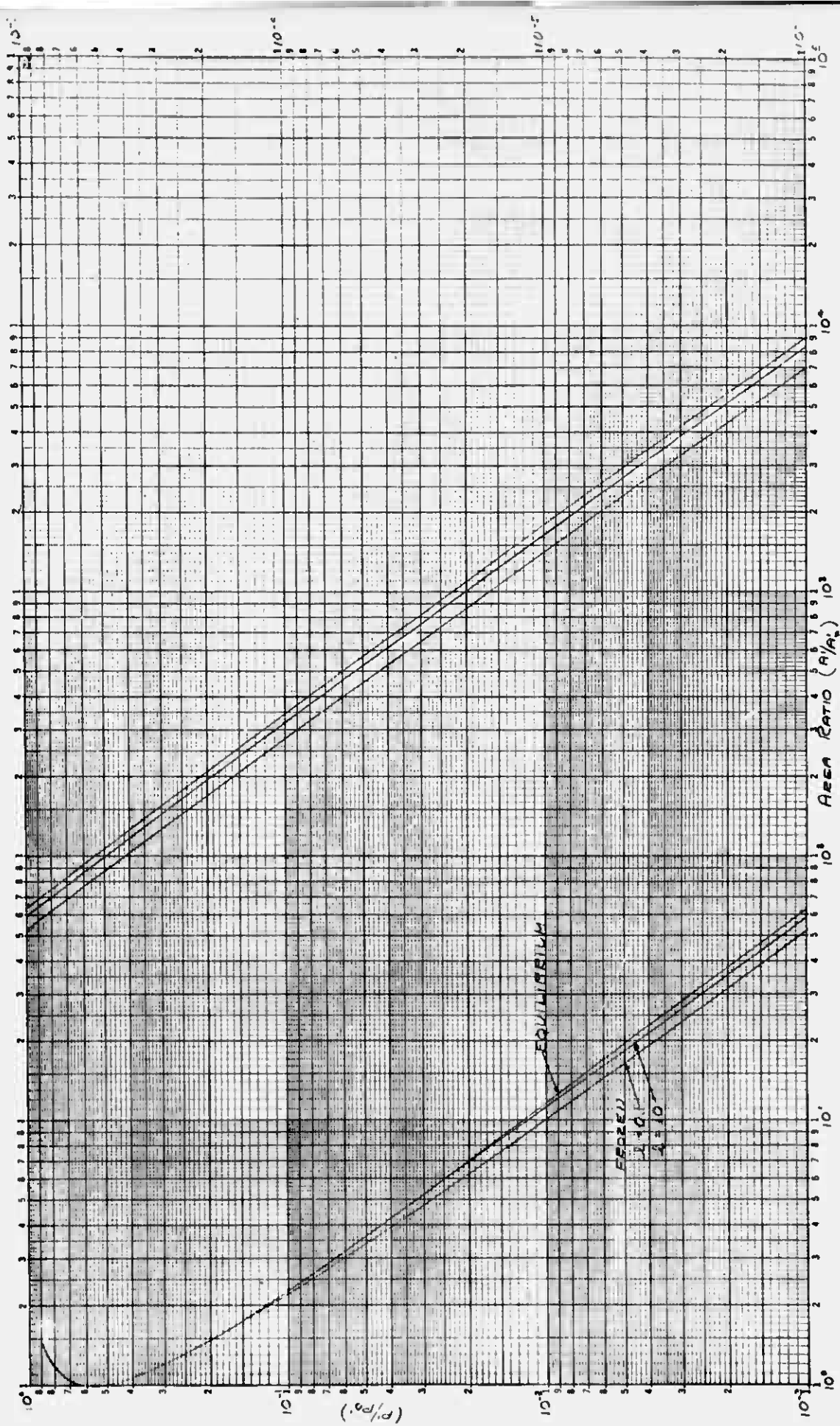


FIGURE NO. 266 HYPERBOLIC NOZZLE $T'_i = 5000^\circ K$ $P'_0 = 200 \text{ ATM}$

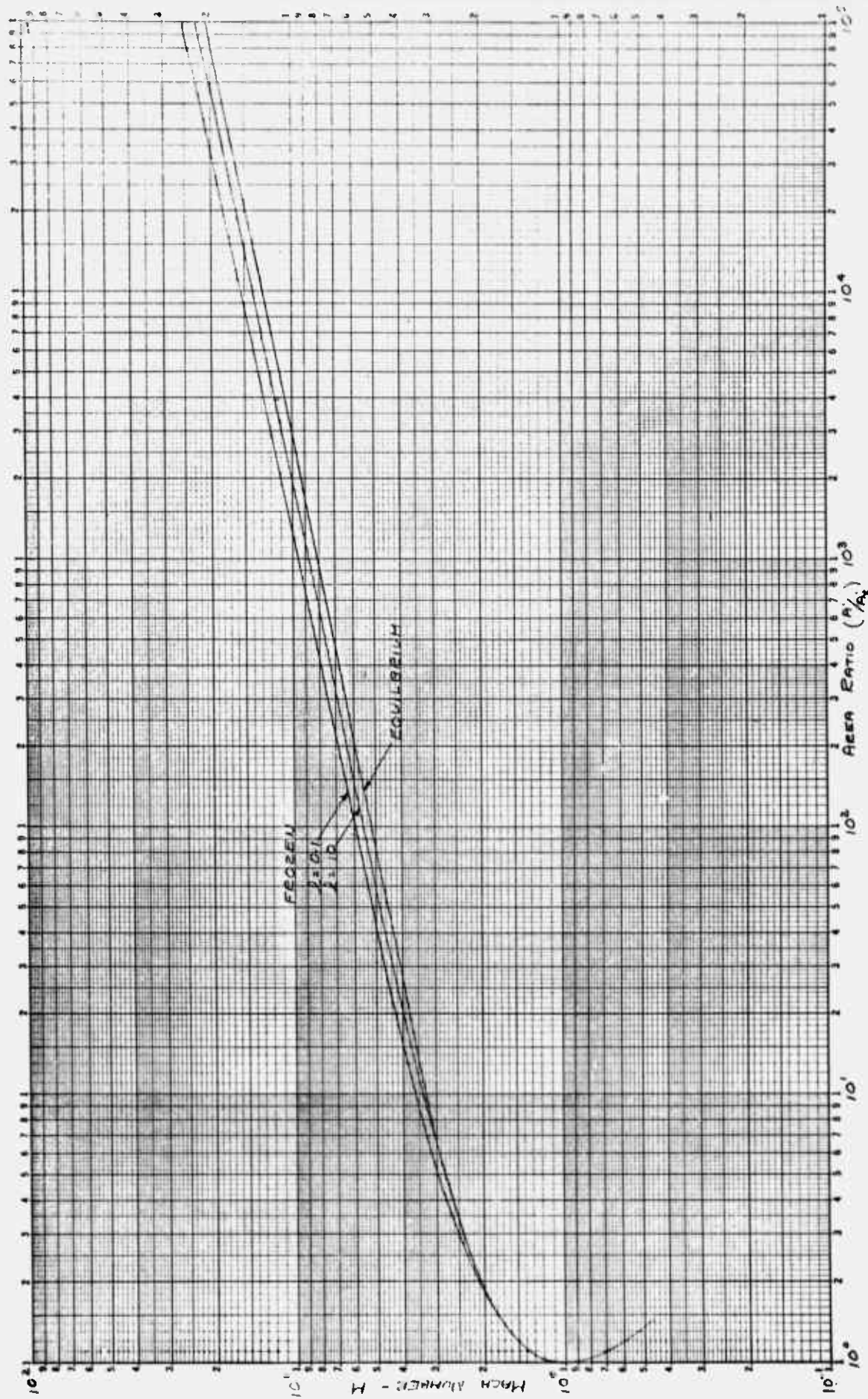


FIGURE NO. 266 $P_2/P_1 = 0.1$ $T_2/T_1 = 1.0$ $T_1 = 5000^\circ R$ $P_1 = 200 \text{ atm}$

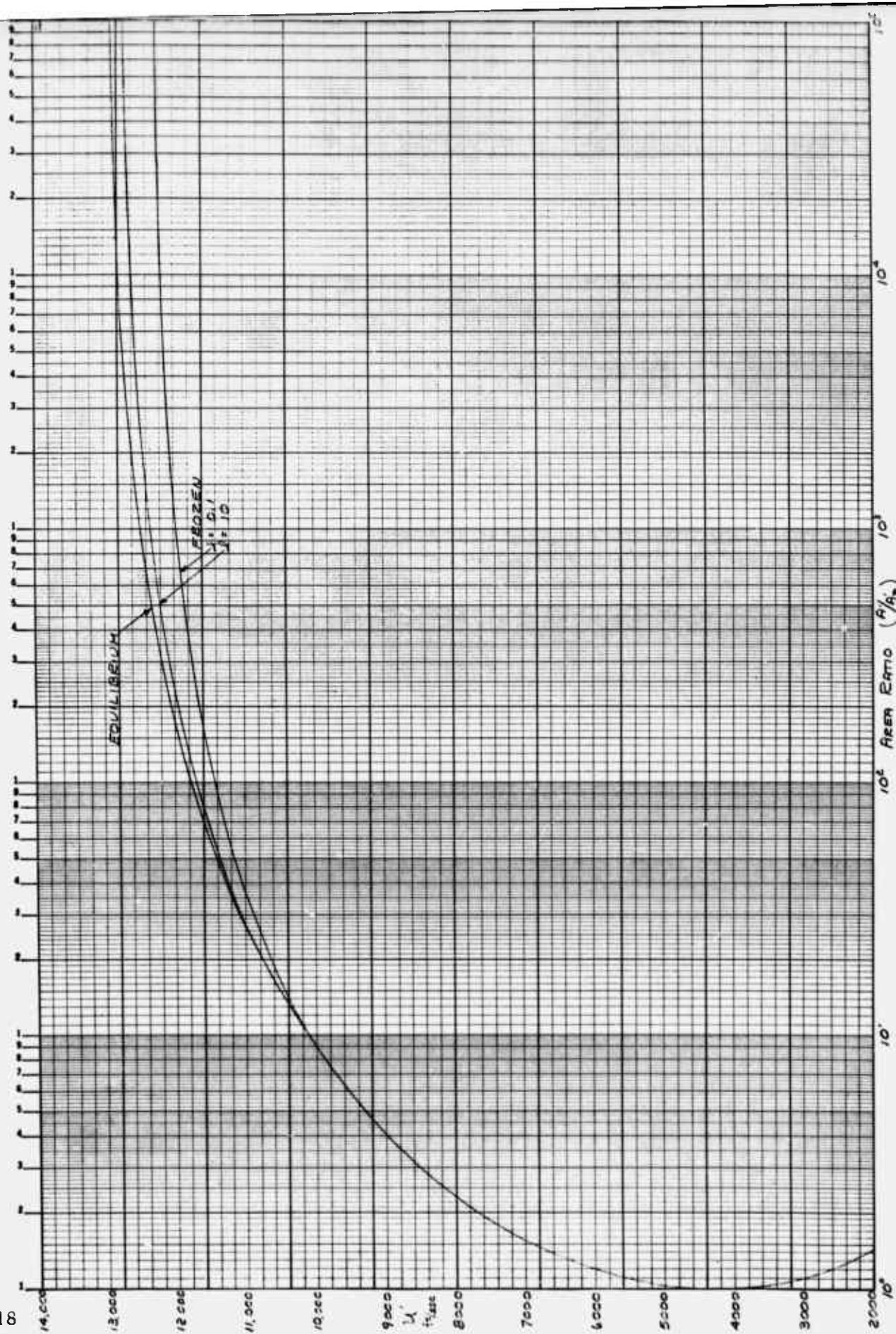


FIGURE NO. 2/3 HYPHEN 20. C NO. 22 LE $T_0 = 5000^\circ \text{K}$ $P_0 = 200 \text{ ATM}$

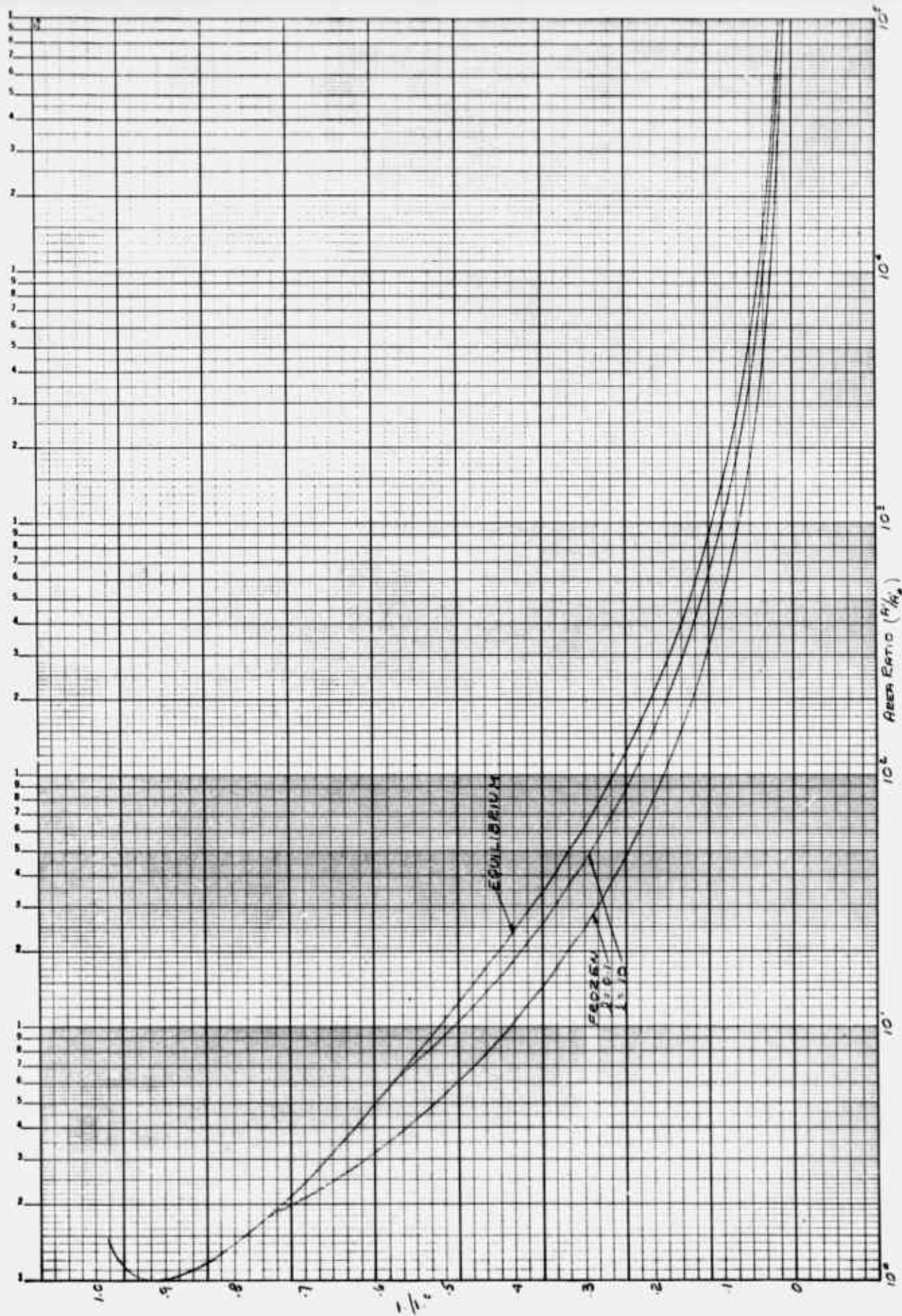


FIGURE NO 262 HYPERBOLIC NOZZLE $T_i = 5000^\circ K$ $P_i = 200 \text{ ATM}$

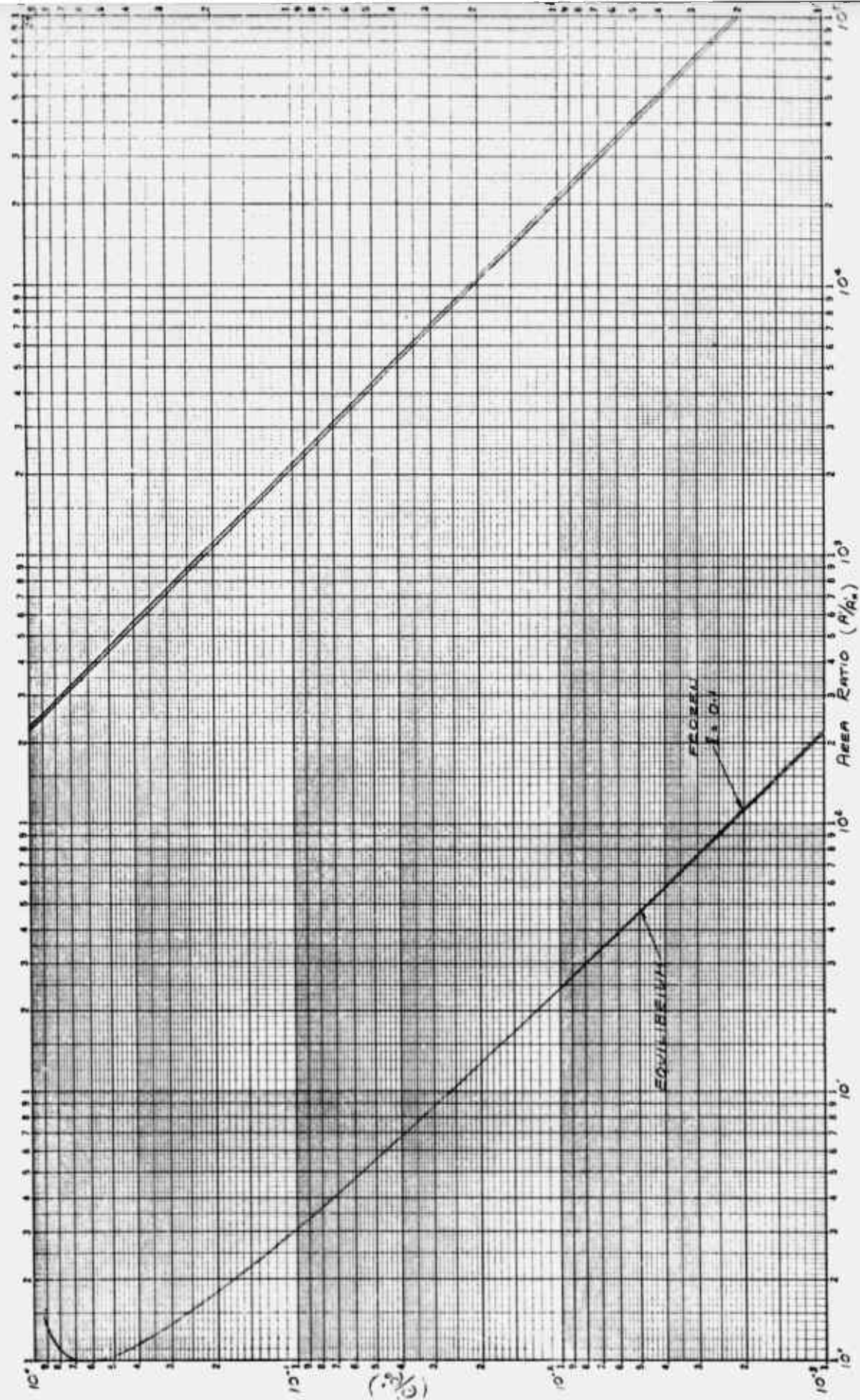


FIGURE NO. 27a NOZZLE $T_0 = 5000^\circ K$ $P_0 = 300 \text{ ATM}$

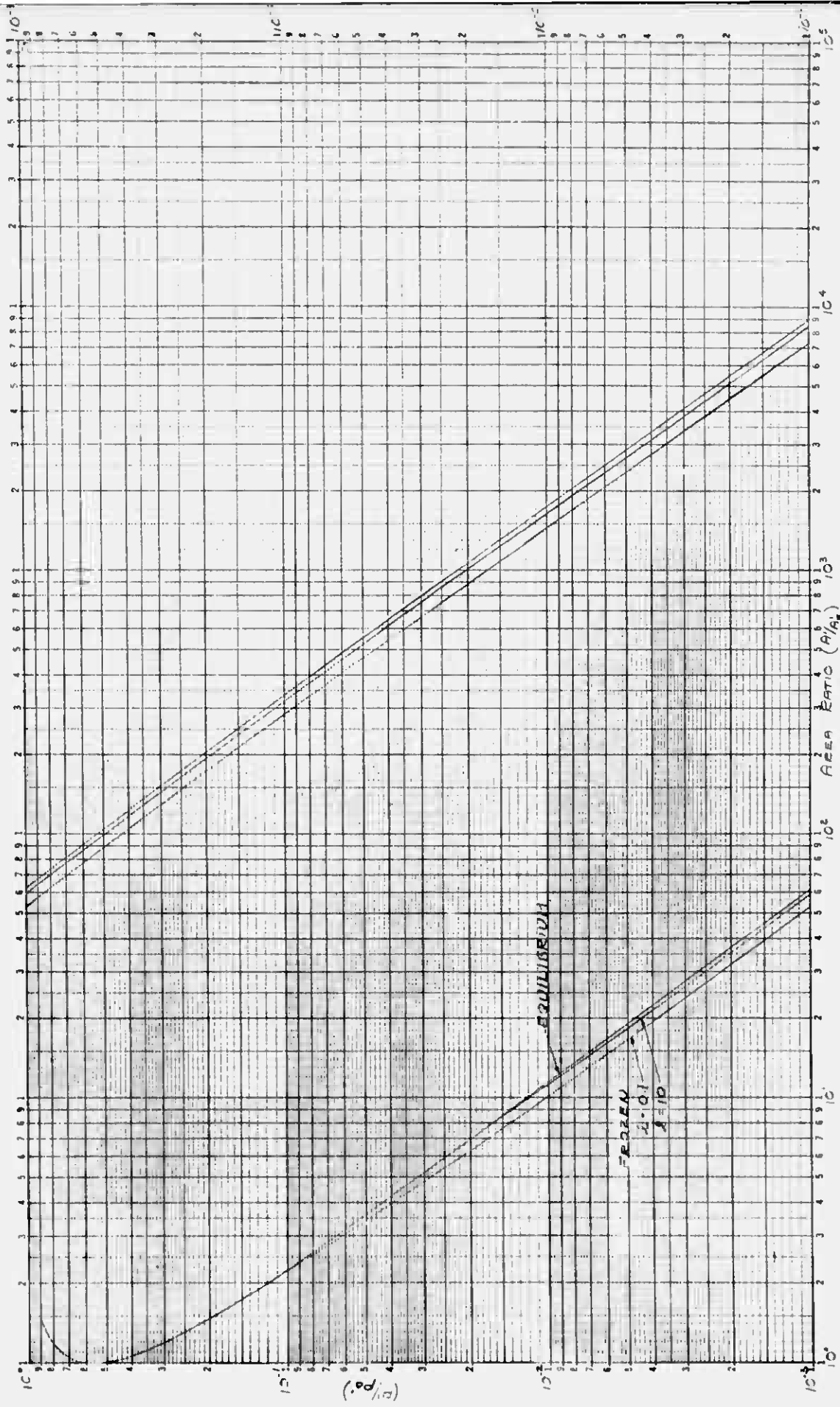


FIGURE NO. 276 HYPERBOLIC NOZZLE $T_0 = 5000^\circ K$ $P_0 = 300 \text{ ATM}$

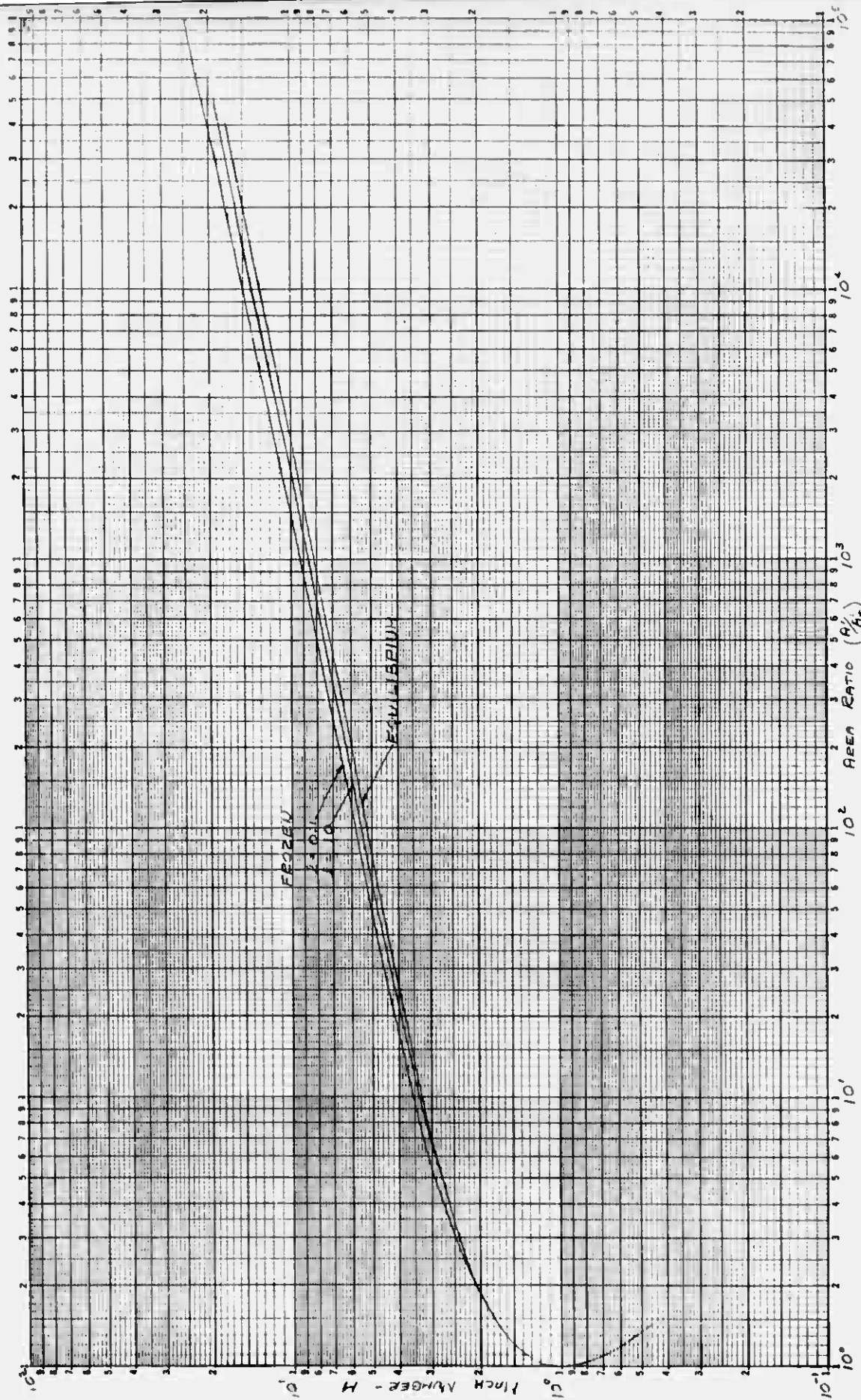


FIGURE NO. 27c HYPERBOLIC NOZZLE $T_0 = 5000^\circ K$ $P_0 = 300 \text{ ATM}$

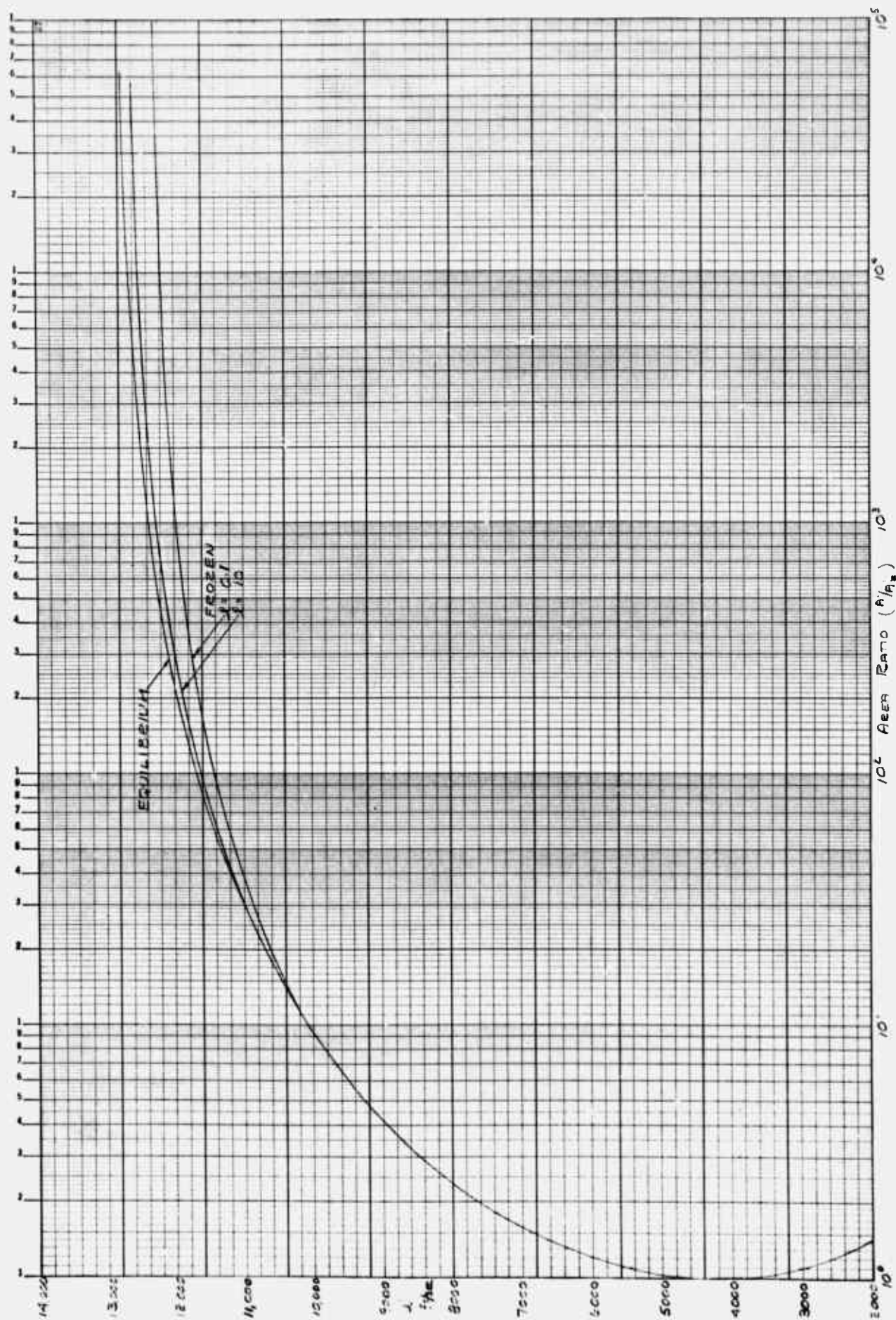


FIGURE NO. 27d HYPERBOLIC NOZZLE $T'_0 = 5000^\circ K$ $P'_0 = 300 \text{ ATM}$

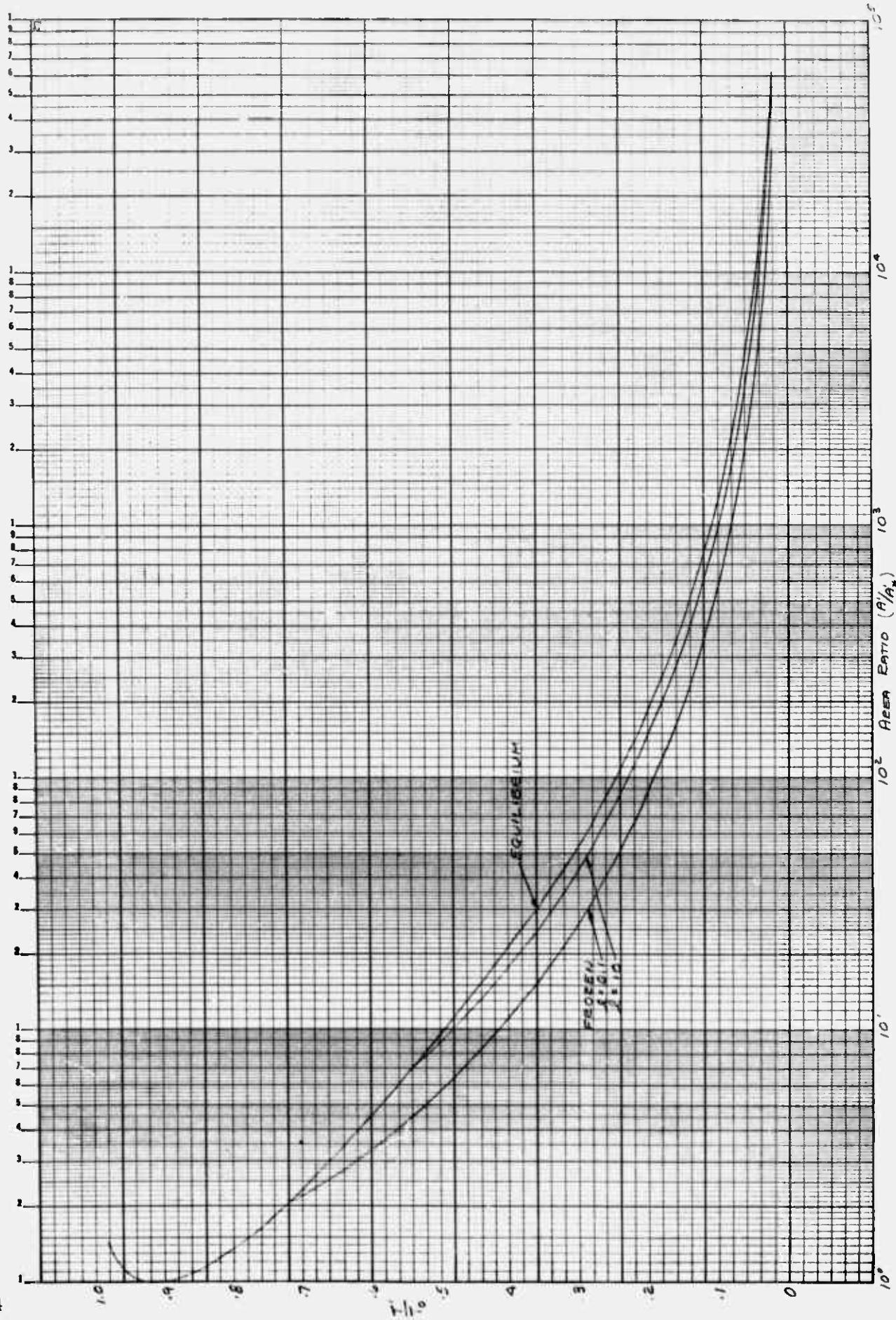


FIGURE NO. 27e HYPERBOLIC NOZZLE $T_c = 5000^\circ K$ $P'_0 = 300 \text{ ATM}$

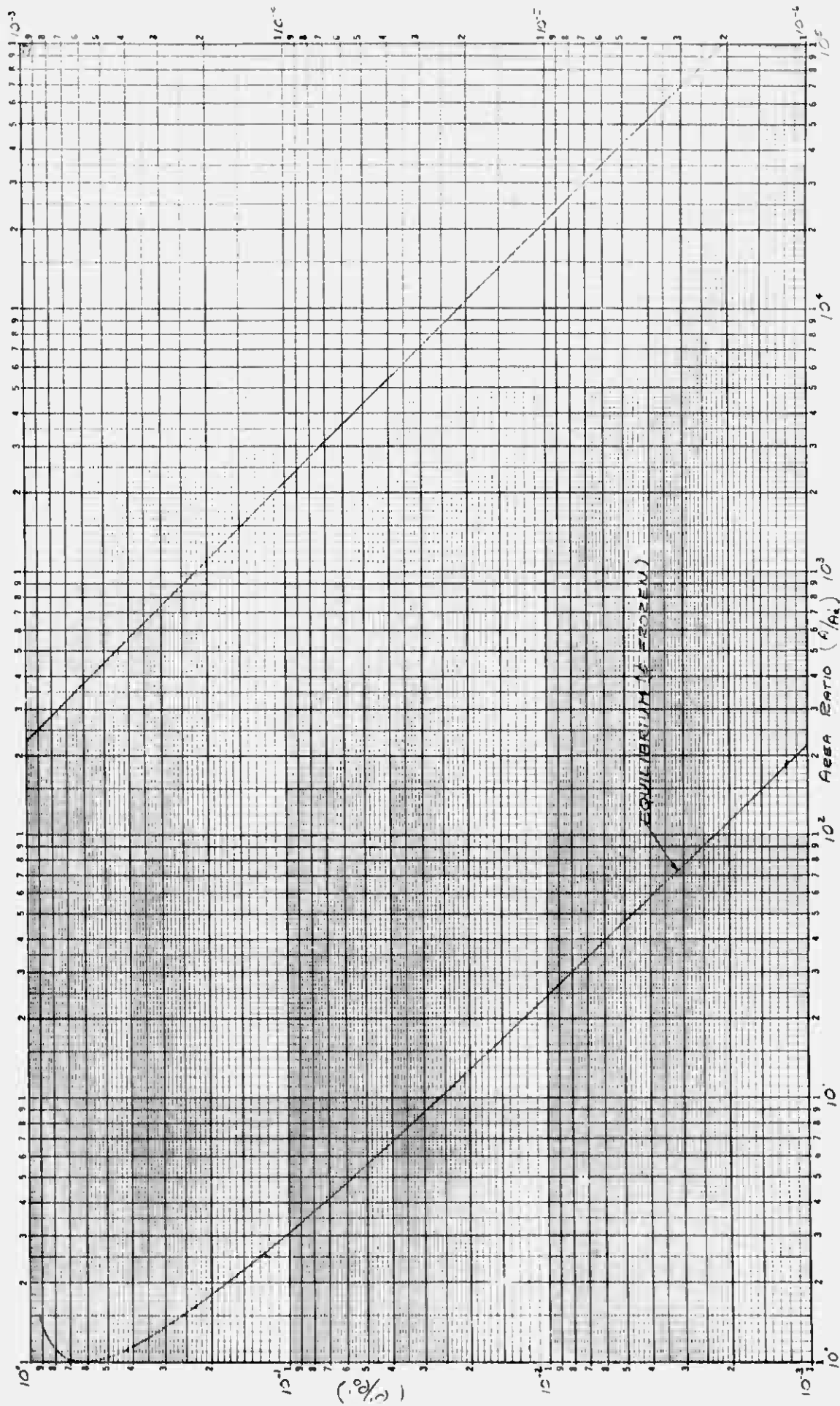


FIGURE NO. 28a HYPERBOLIC NOZZLE $T'_0 = 5000^\circ K$ $P'_0 = 1000 \text{ ATM}$

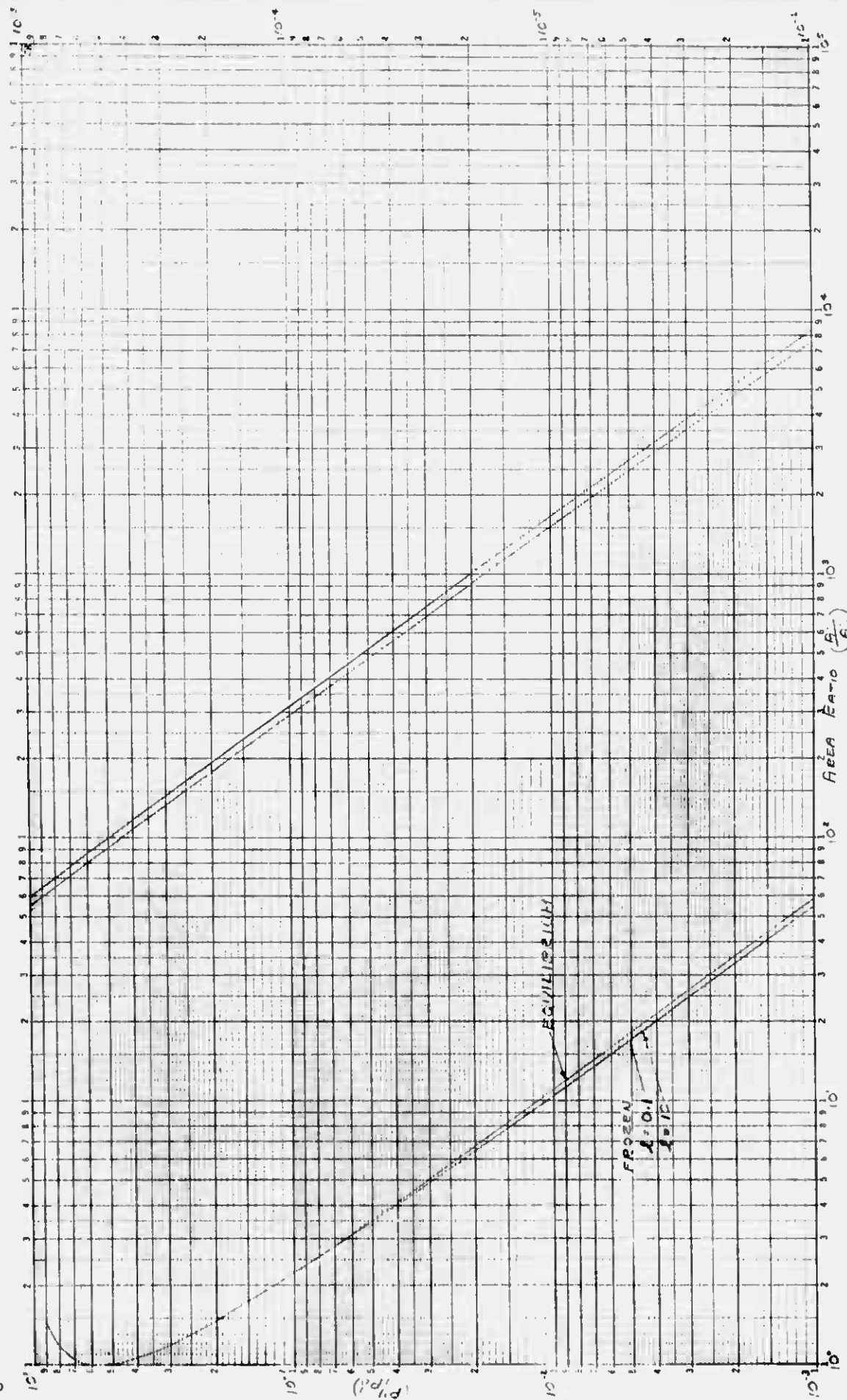


FIGURE NO. 2A6 HYPERBOLIC NOZZLE $T'_0 = 5000^\circ K$ $P'_0 = 1000$ mm Hg

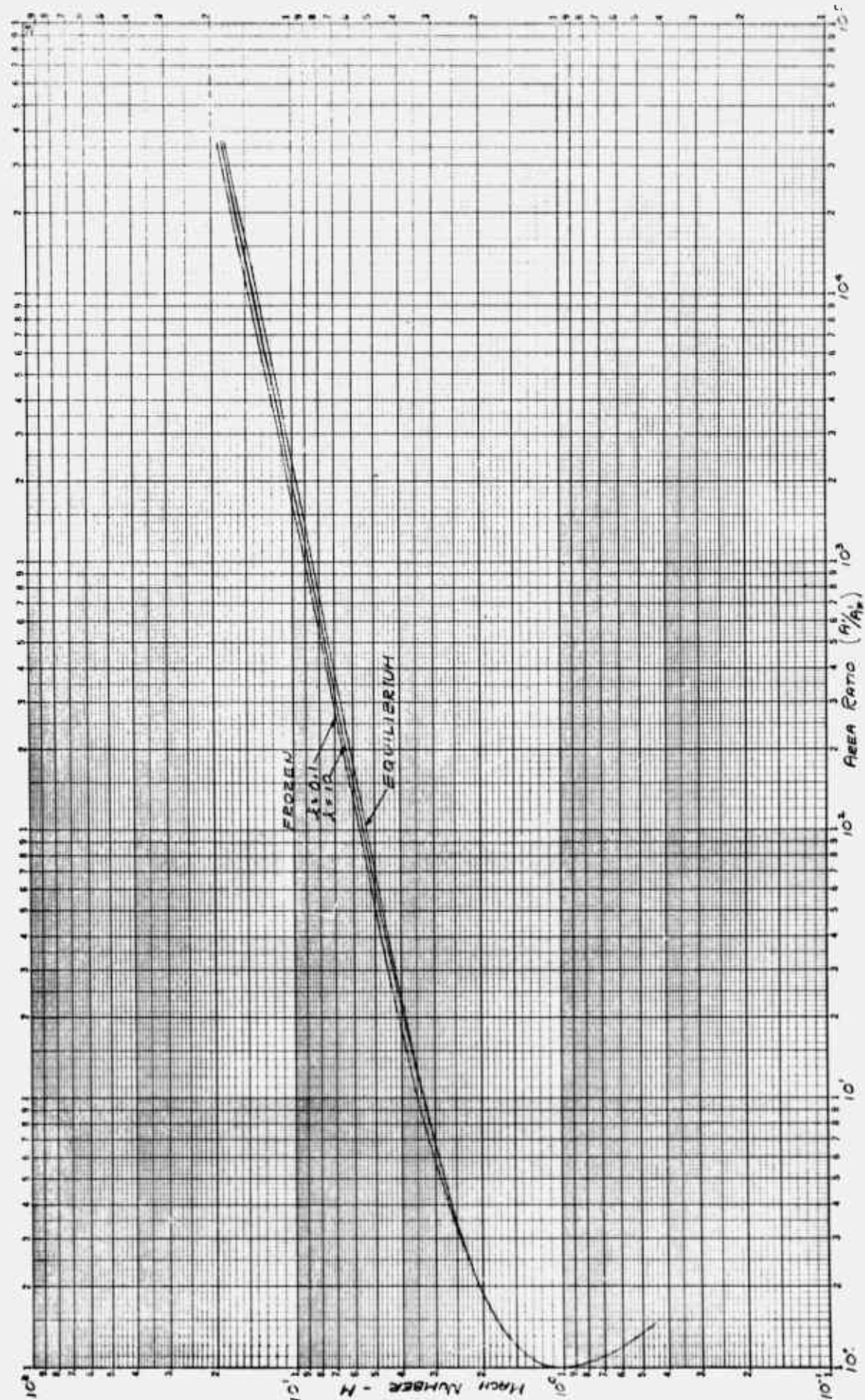


FIGURE NO. 28c HYPERBOLIC NOZZLE $T_0 = 5000^\circ K$ $P_0 = 1000 \text{ ATM}$

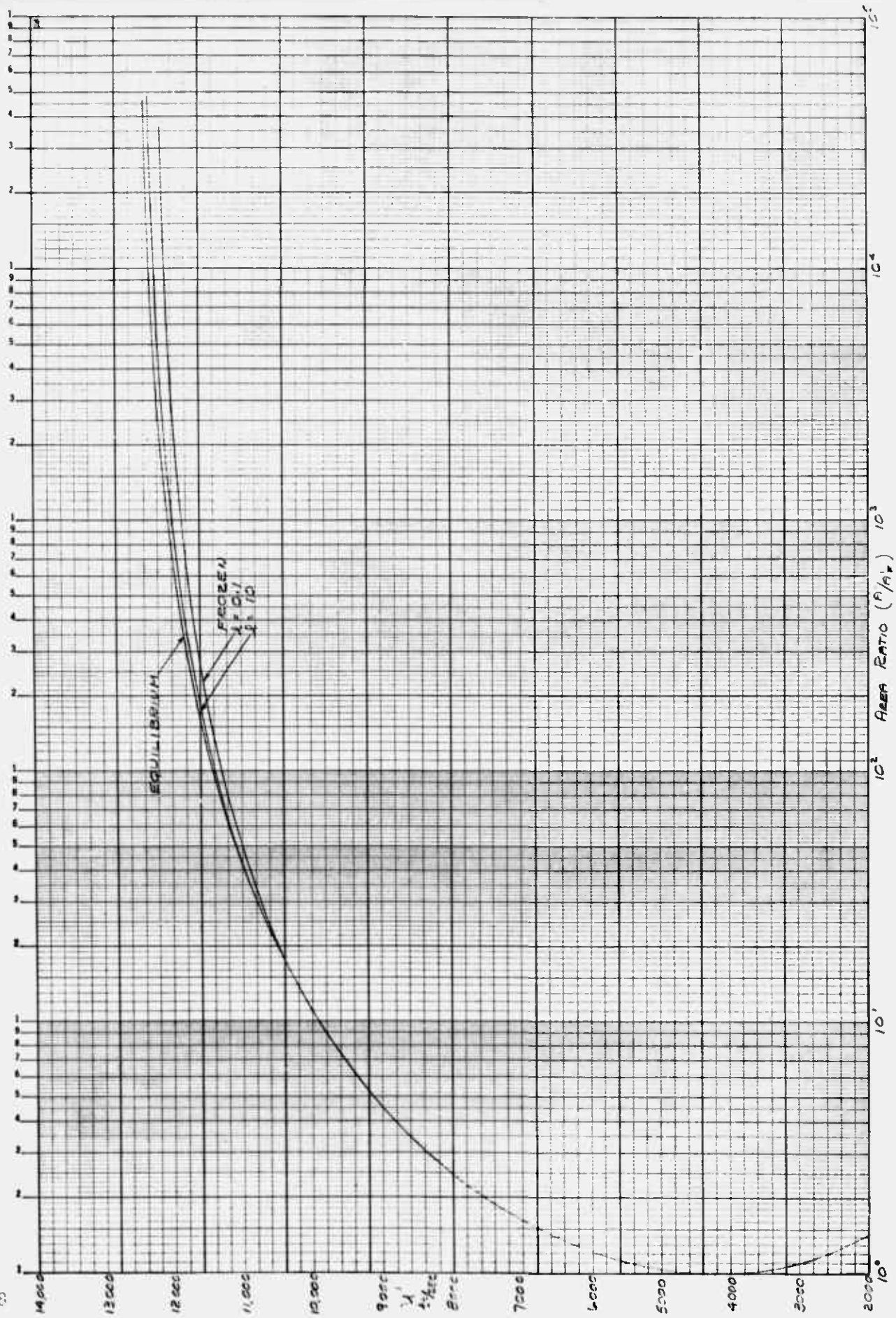


FIGURE NO. 28d HYPERBOLIC NOZZLE $T_0 = 5000^\circ\text{K}$ $P_0 = 1000\text{ ATM}$

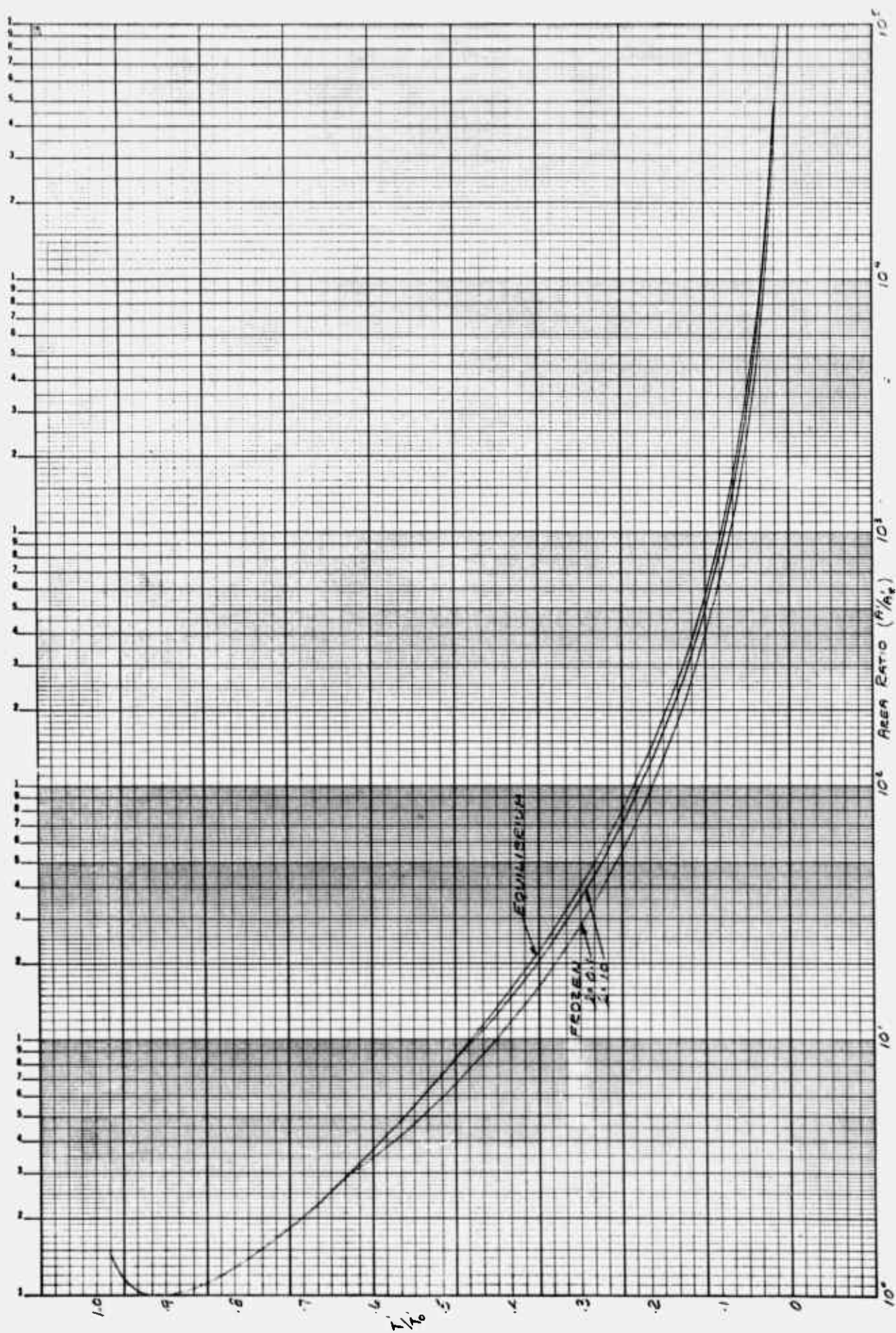


FIGURE NO. 28c HYPERBOLIC NOZZLE $T_1 = 5000^\circ\text{K}$ $P_1 = 1000\text{ ATM}$

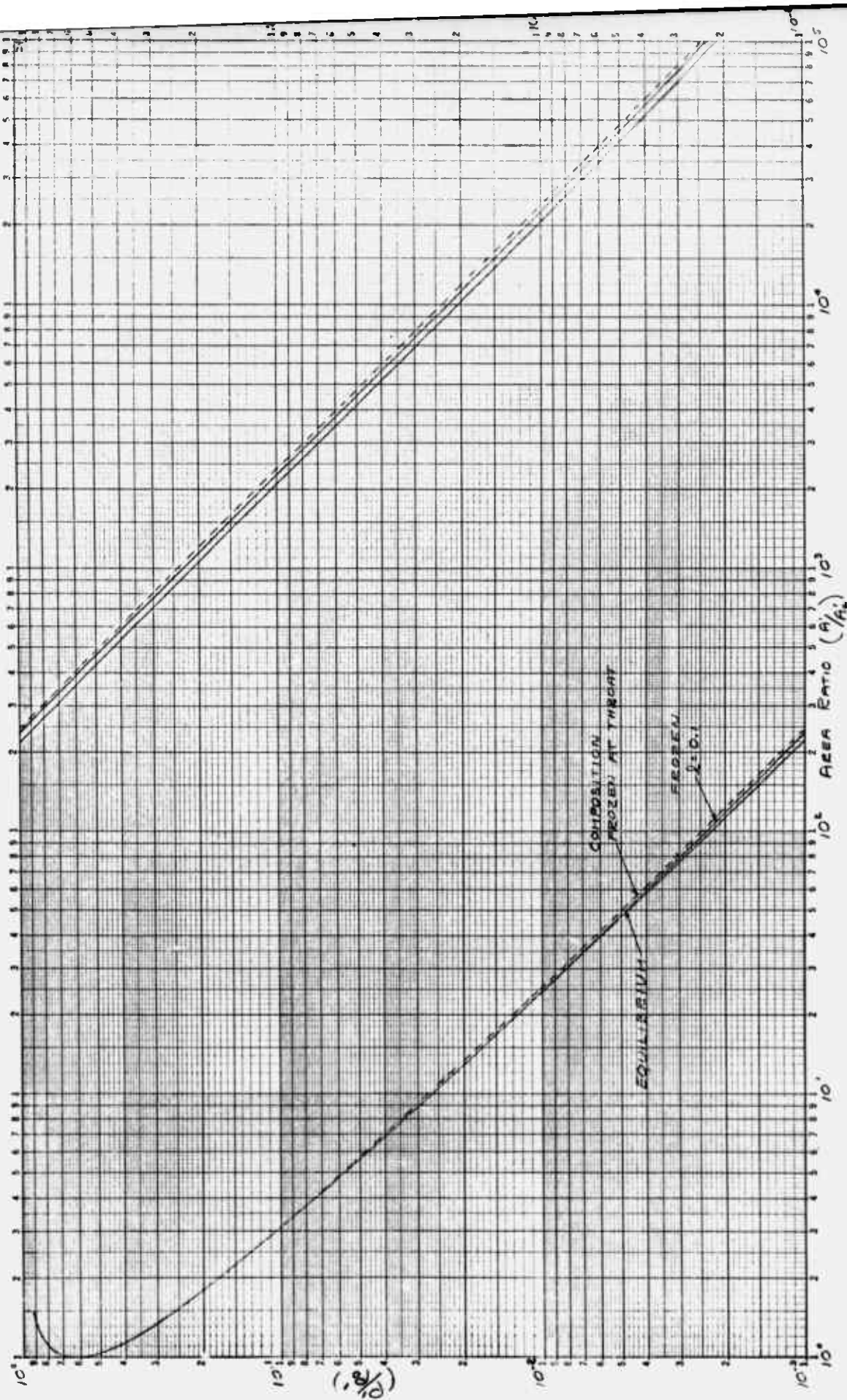


FIGURE NO. 29a HYPERBOLIC NOZZLE $T'_0 = 6000^\circ K$ $P'_0 = 100 \text{ ATM}$

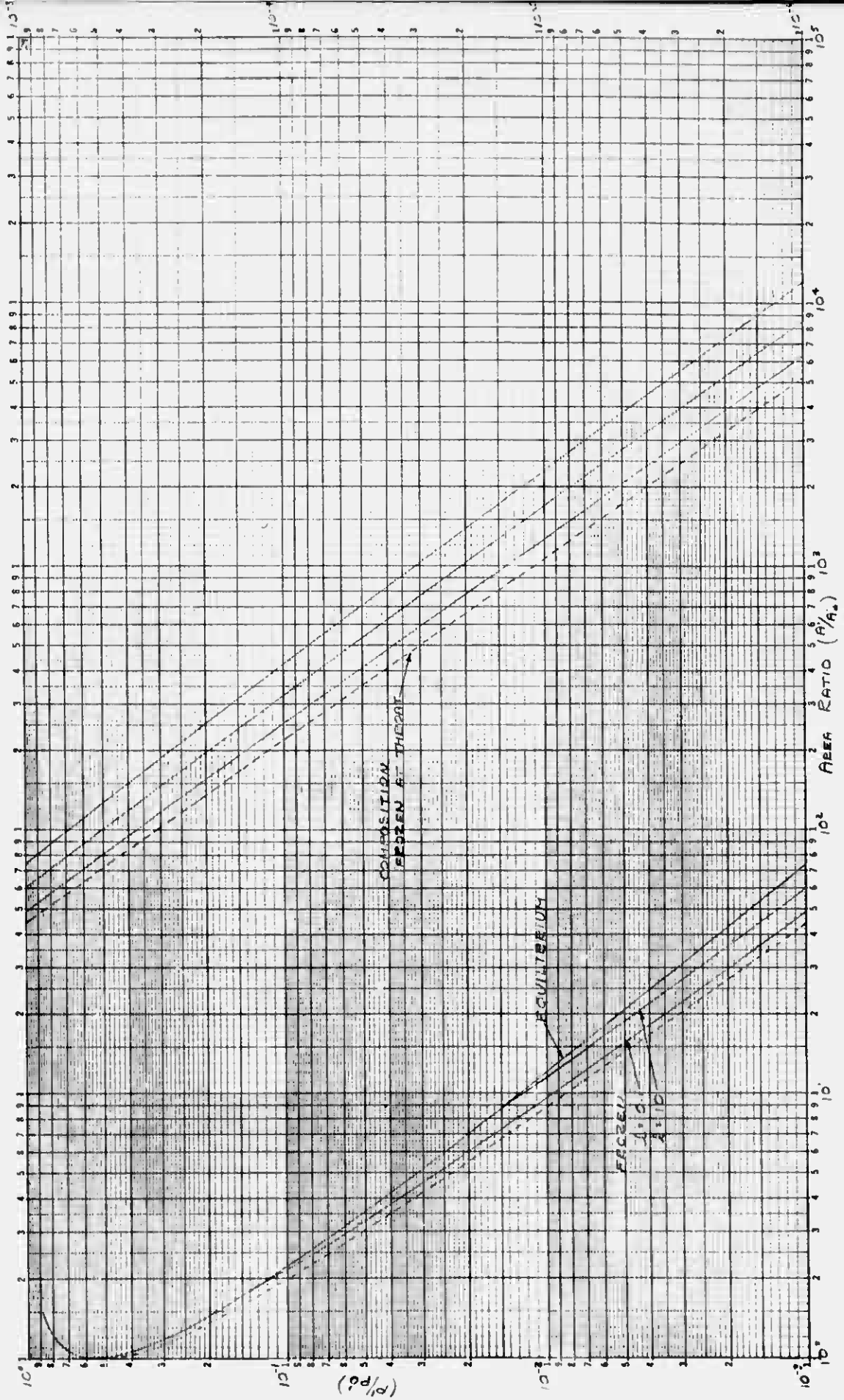


FIGURE NO. 296 HYPERBOLIC NOZZLE $T_c' = 6000^\circ K$ $P_c' = 100 \text{ ATM}$

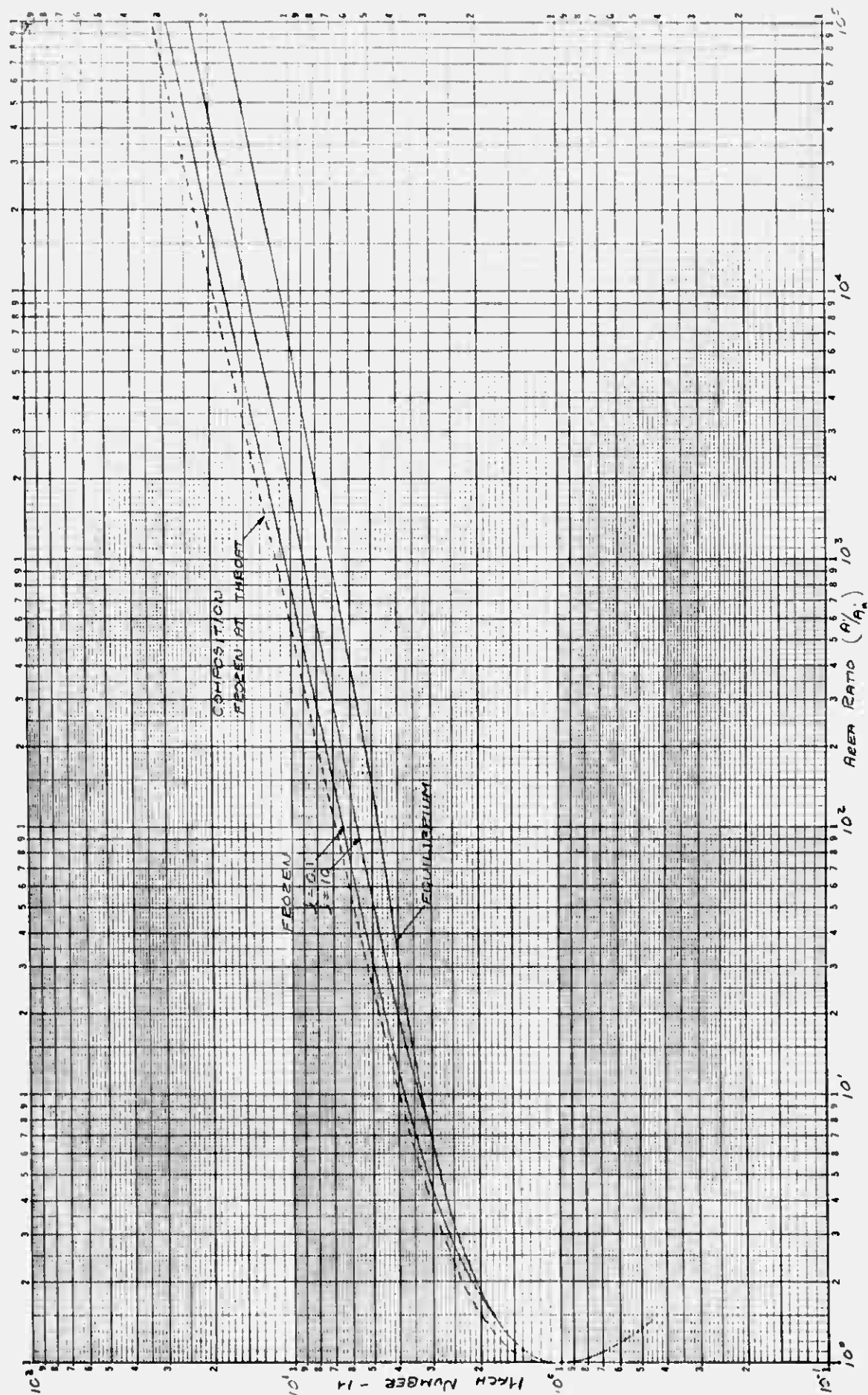


FIGURE NO. 29c HYPERBOLIC NOZZLE $T_0 = 6000^\circ\text{K}$ $P_0 = 100\text{ ATM}$

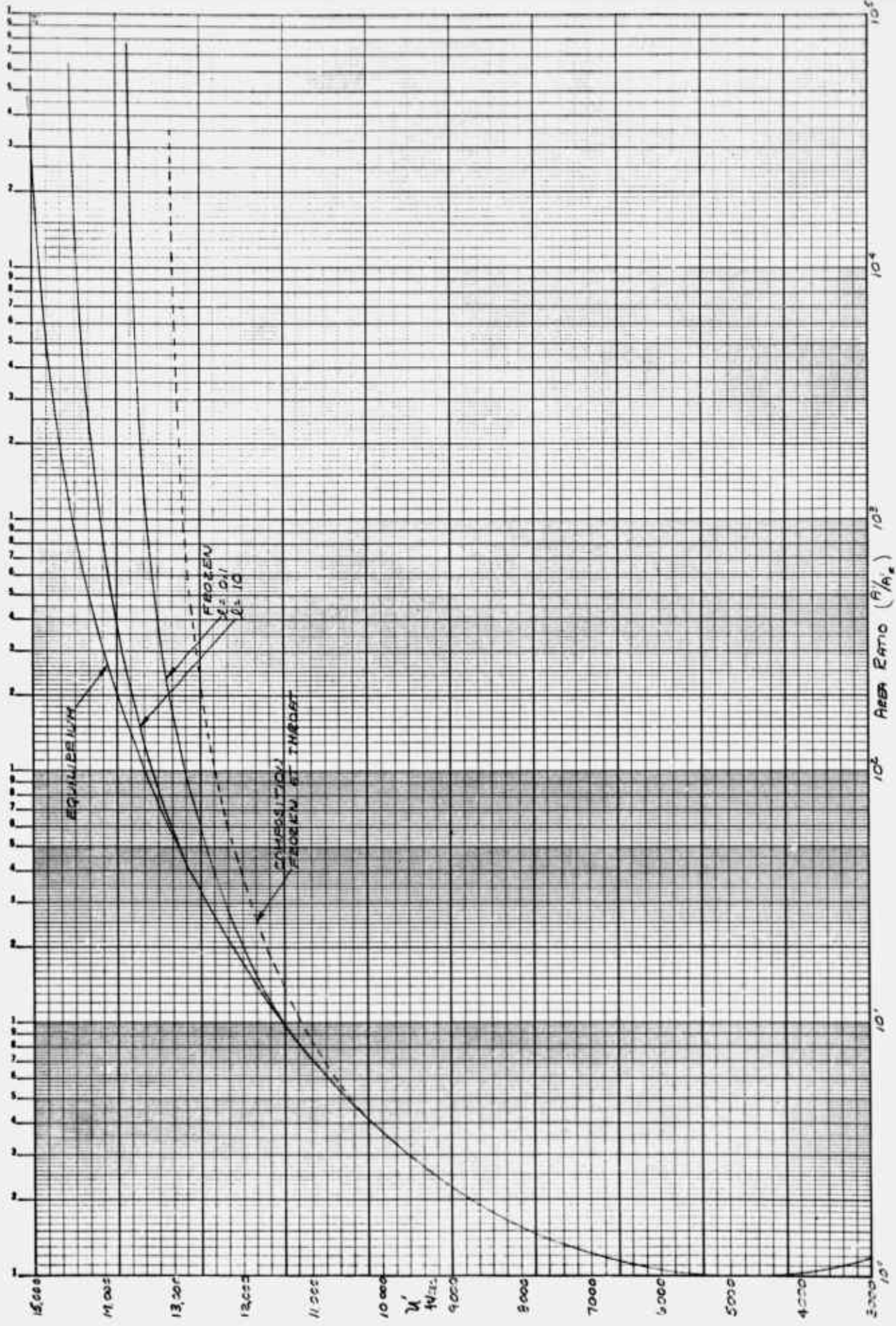


FIGURE NO 29D HYPERBOLIC NOZZLE $T'_0 = 6000^\circ K$ $P'_0 = 100 \text{ ATM}$

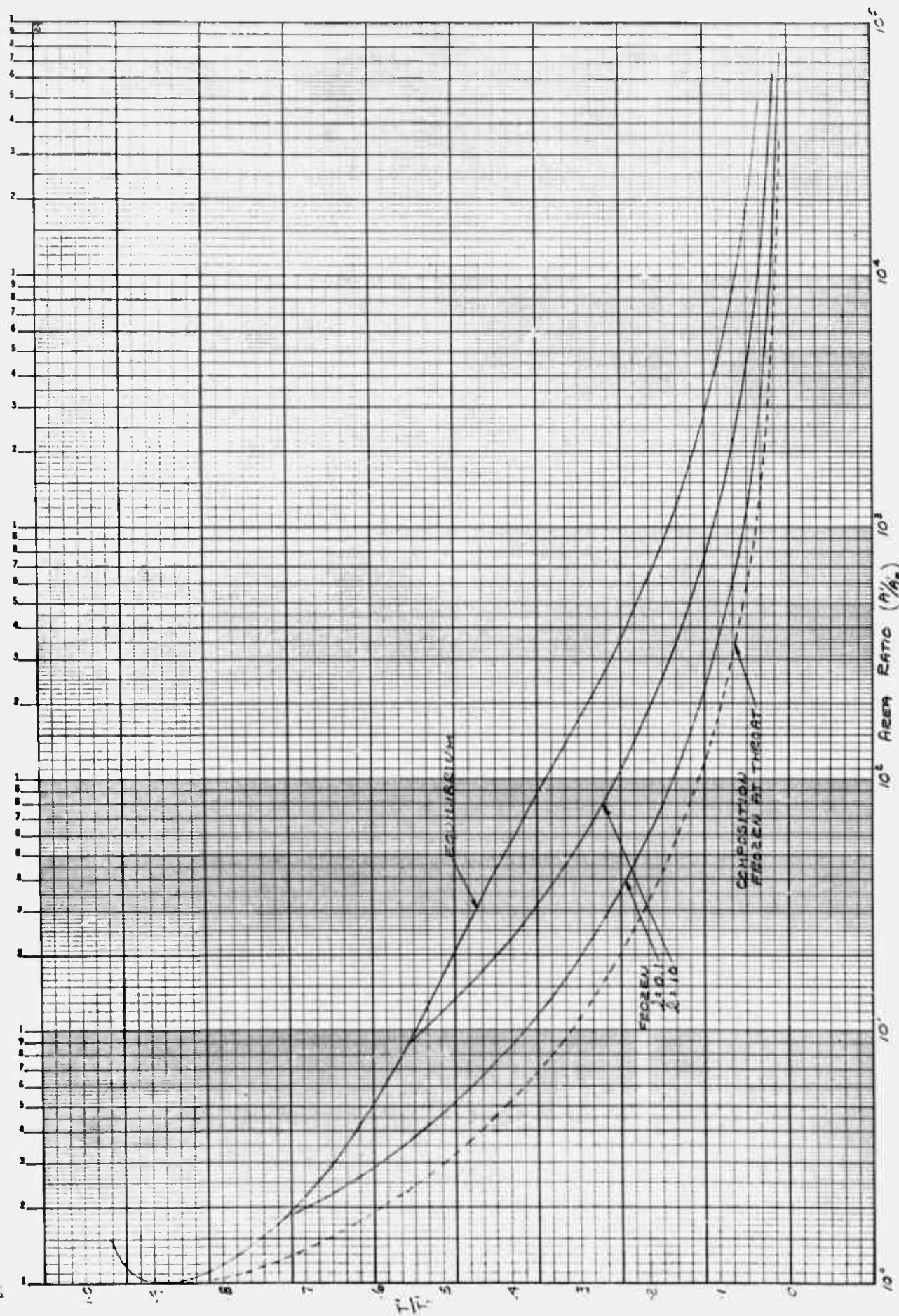


FIGURE NO 292 HYPERSONIC NOZZLE $T_0 = 6000^\circ K$ $P_0 = 100 \text{ ATM}$

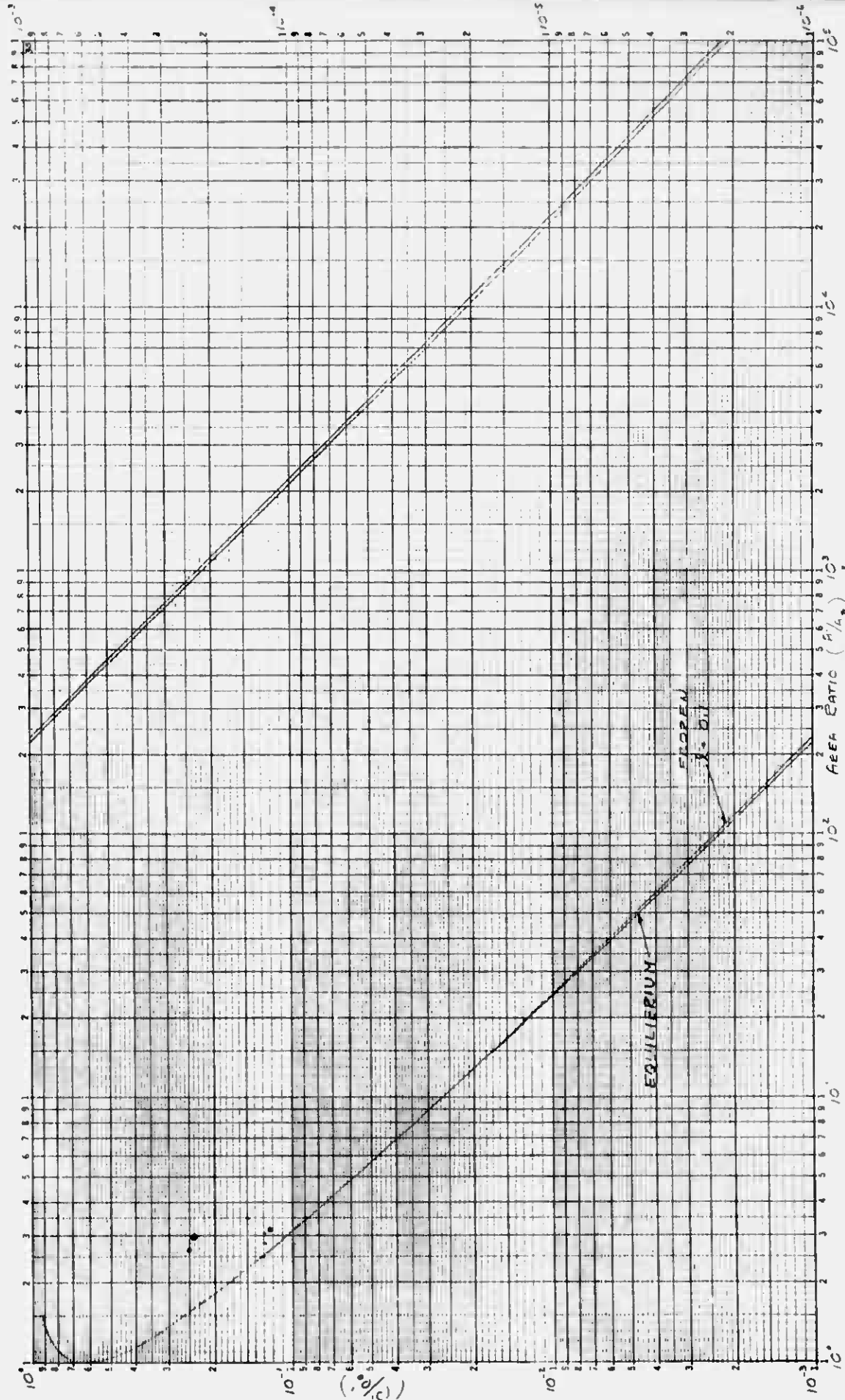
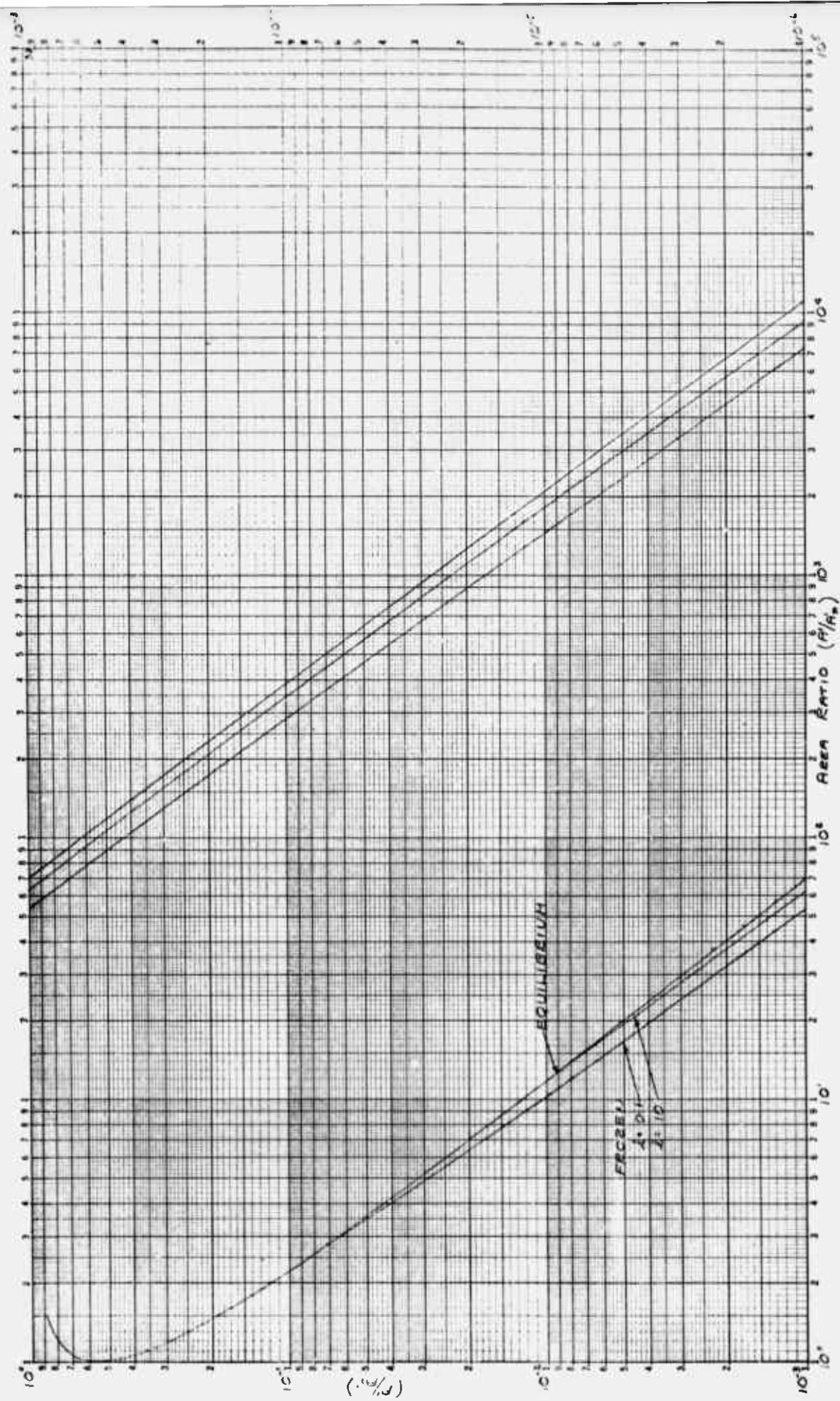


FIGURE NO. 30a HYPERBOLIC NOZZLE $T_0 = 6000^\circ R$ $P_0 = 3000 \text{ PSI}$

FIGURE NO. 30% HYPERBOLIC NOZZLE $T_0 = 6000^\circ K$ $P_0 = 300 \text{ ATM}$

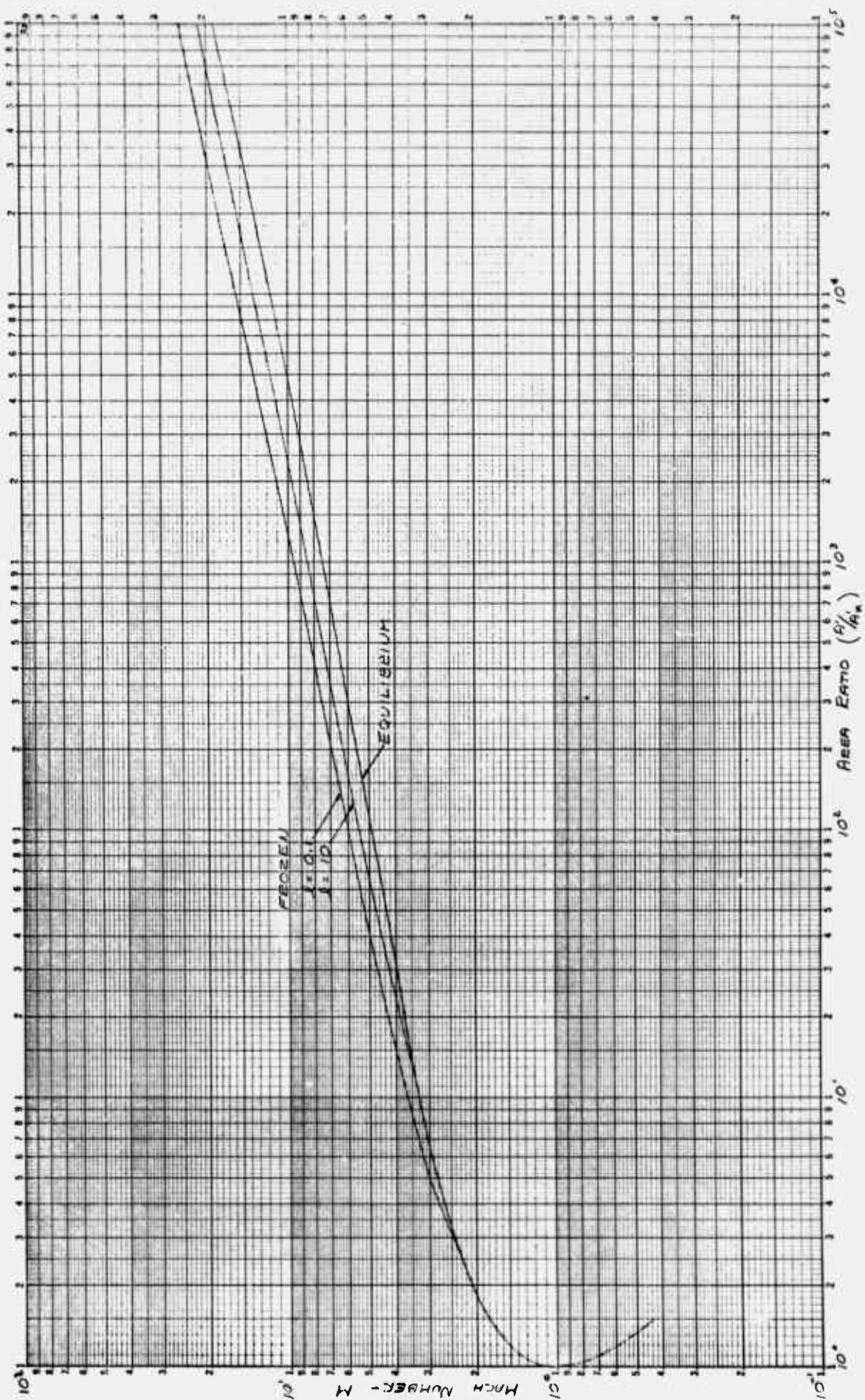


FIGURE NO. 30c HYPERBOLIC NOZZLE $T_i = 6000^\circ K$ $P_i = 300 \text{ ATM}$

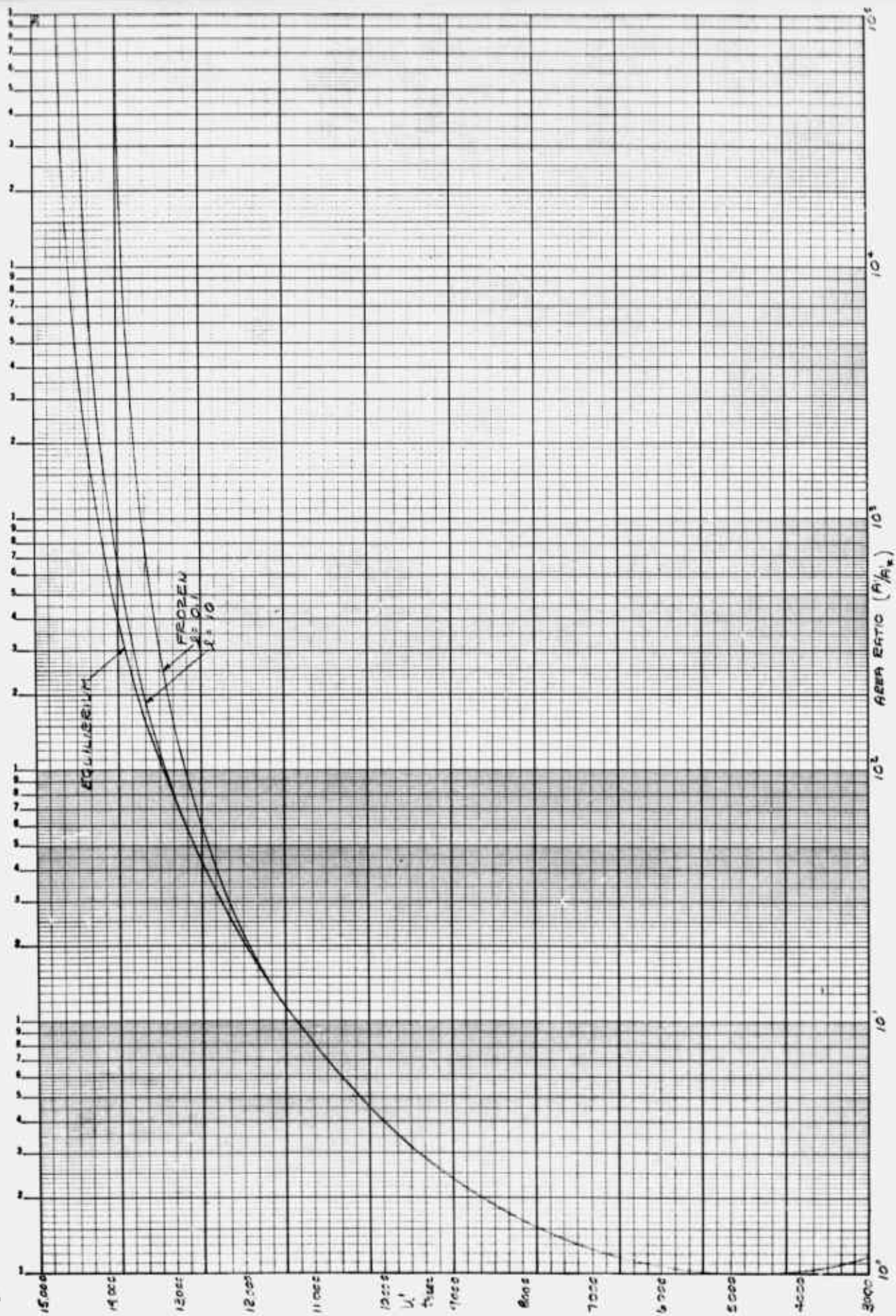


FIGURE NO. 301 HYPERSONIC NOZZLE $T'_0 = 6000^\circ\text{K}$ $P'_0 = 300\text{ ATM}$

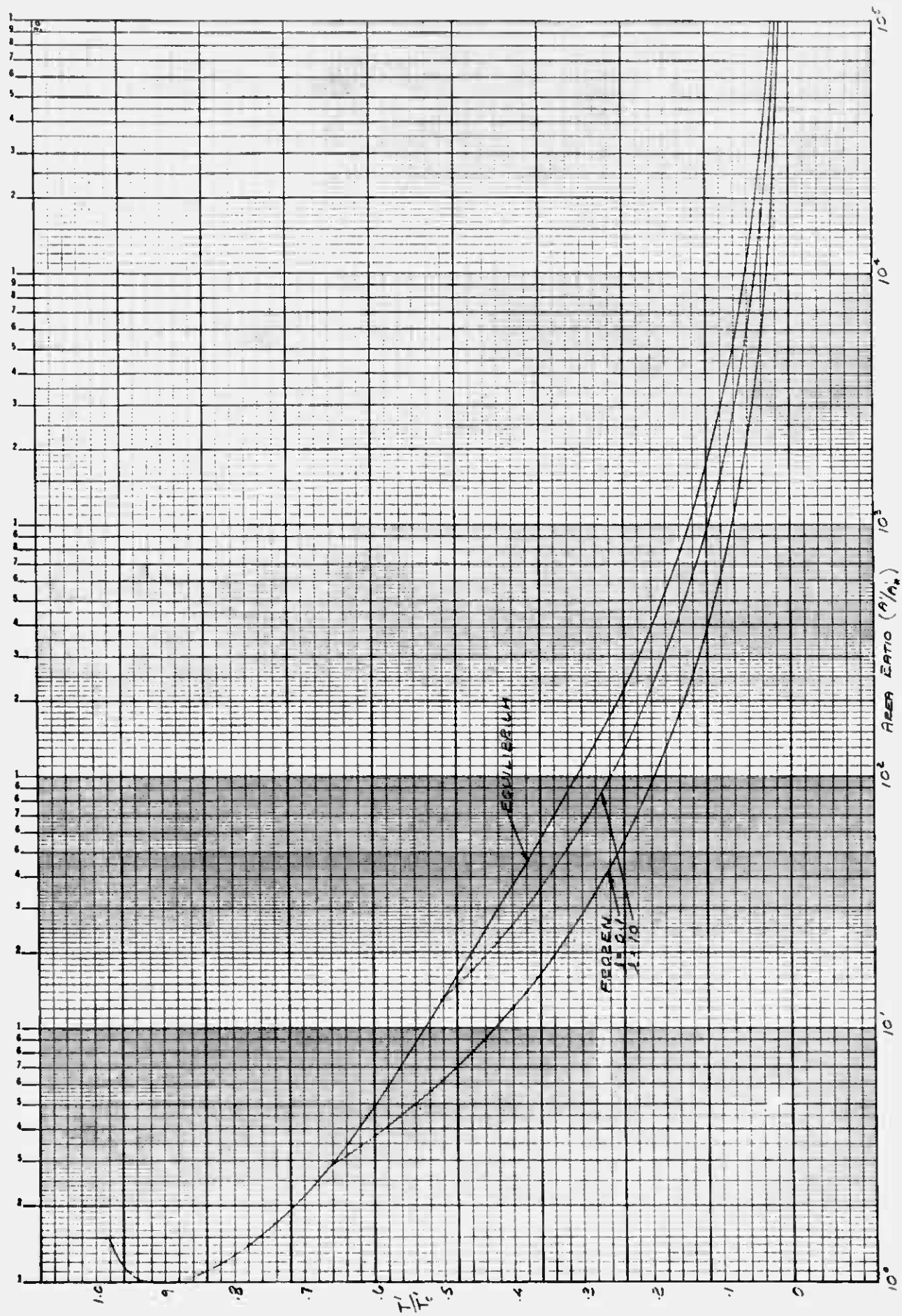


FIGURE NO 30c. HYPERBOLIC NOZZLE $T'_0 = 6000^\circ K$ $P'_0 = 300 \text{ ATM}$

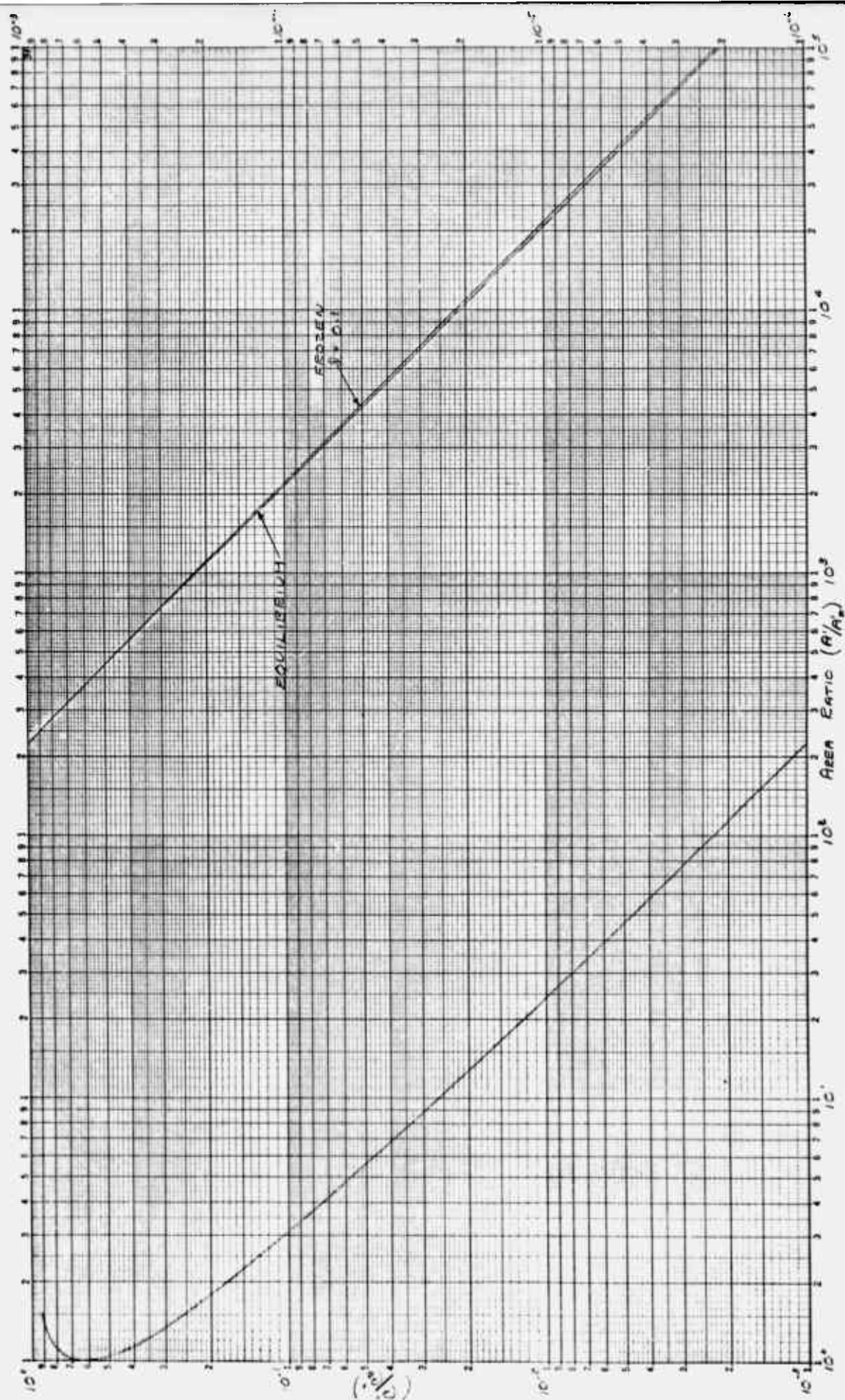


FIGURE NO. 31a HYPERBOLIC NOZZLE $T'_0 = 6000^\circ K$ $P'_0 = 1000 \text{ ATM}$

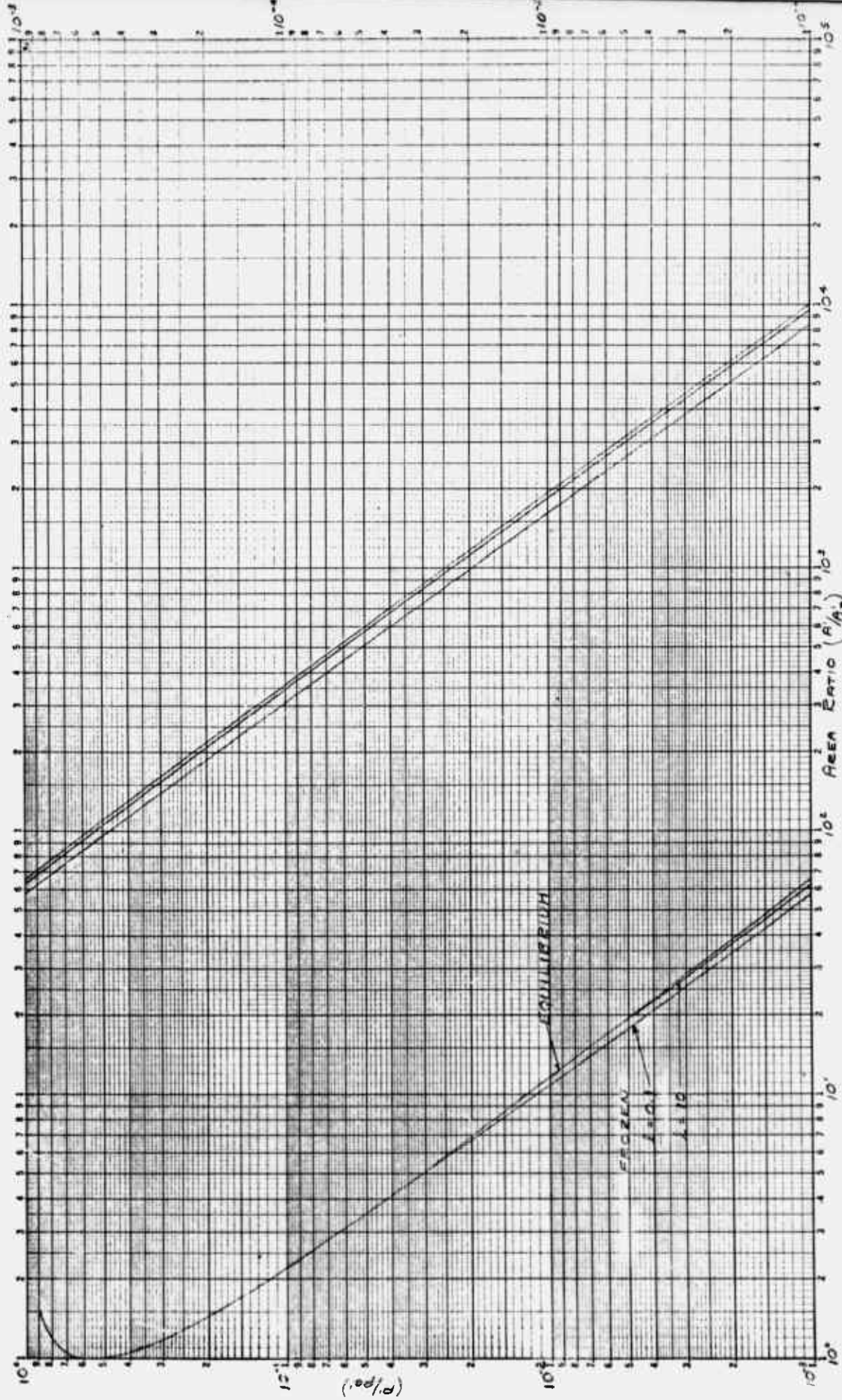


FIGURE NO. 316 HYPERBOLIC NOZZLE $T'_0 = 6000^\circ K$ $P'_0 = 1000 \text{ ATM}$

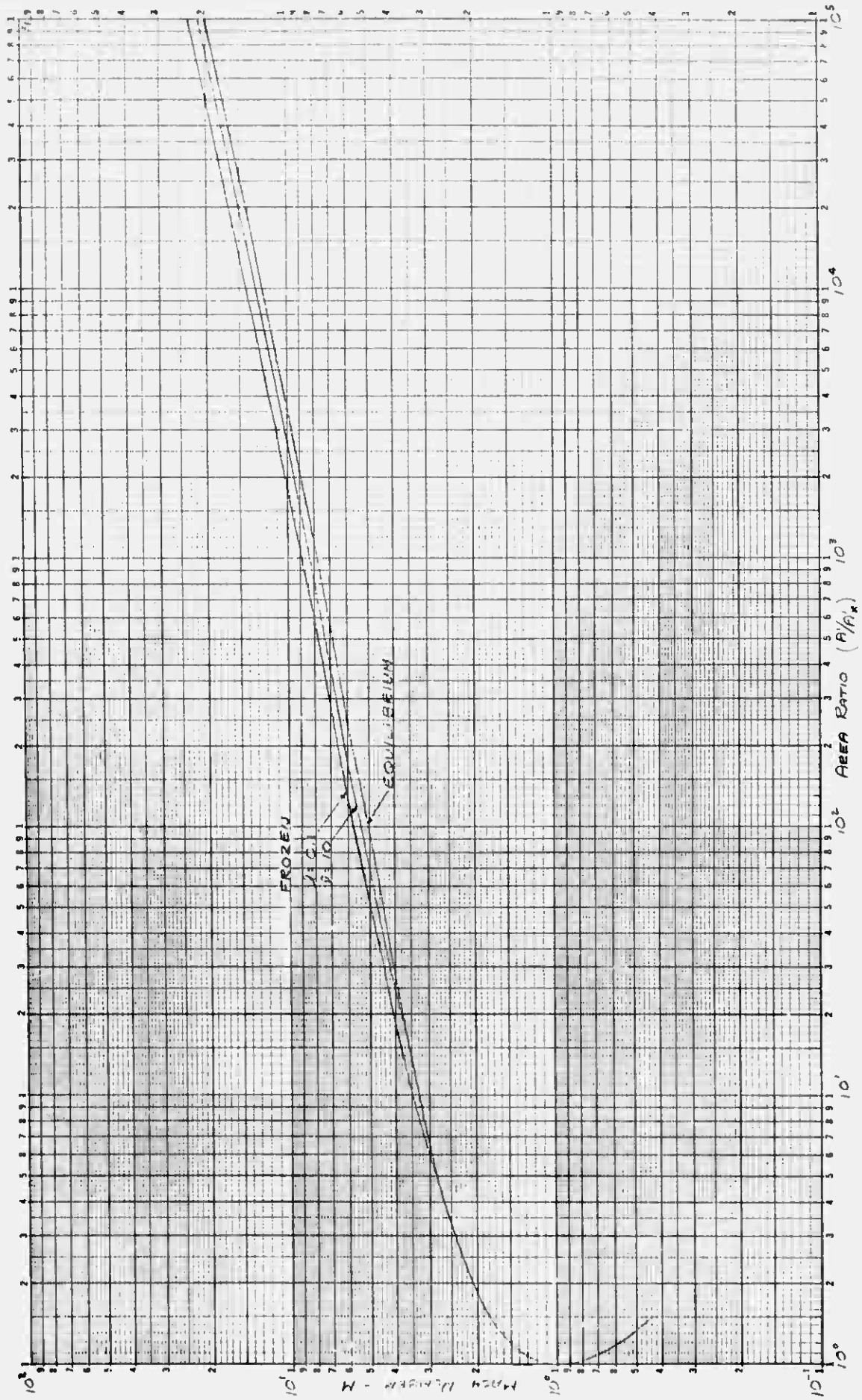


FIGURE NO. 31C HYPERBOLIC NOZZLE $T_0 = 6000^\circ K$ $P_0 = 1000 \text{ ATM}$

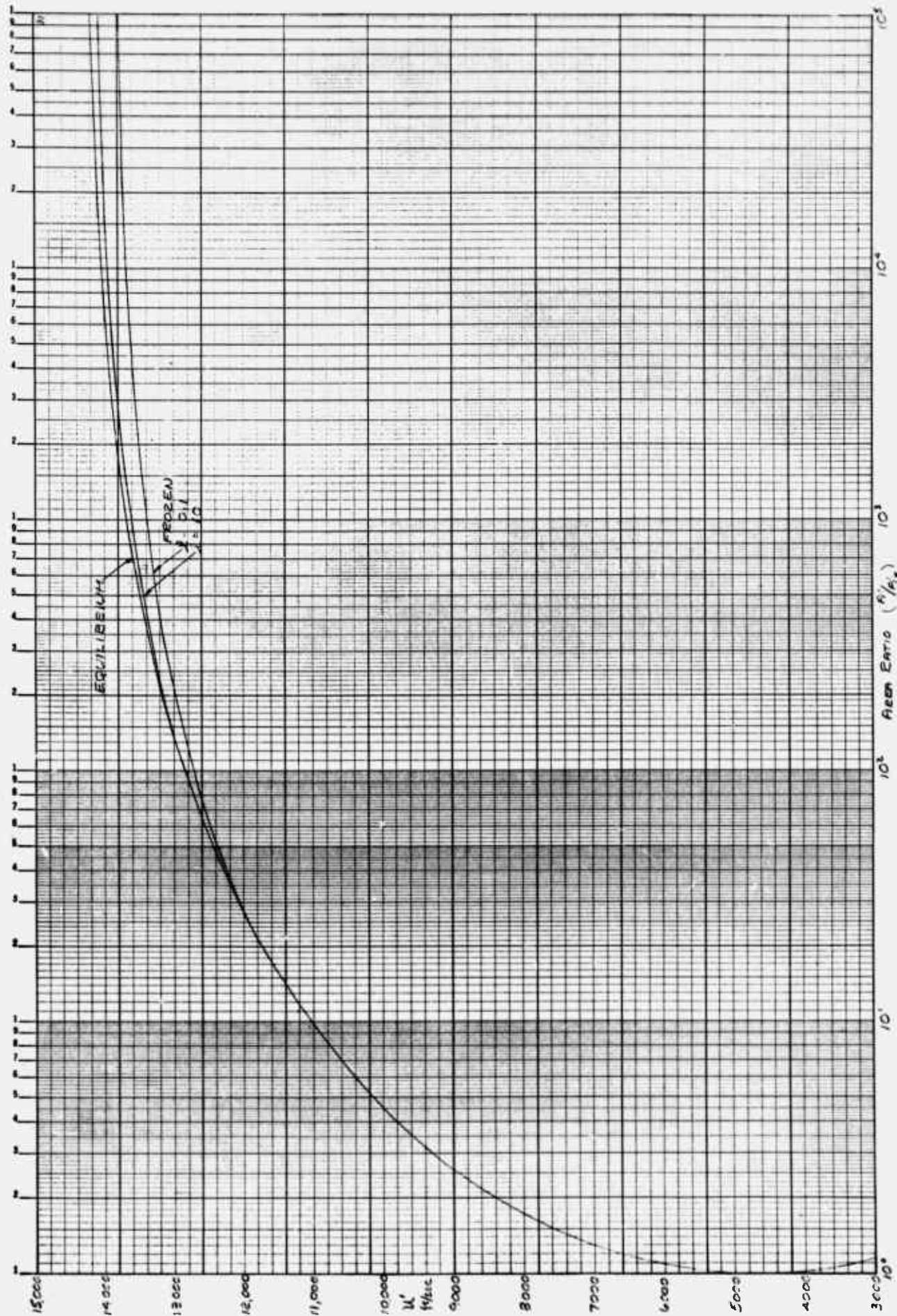
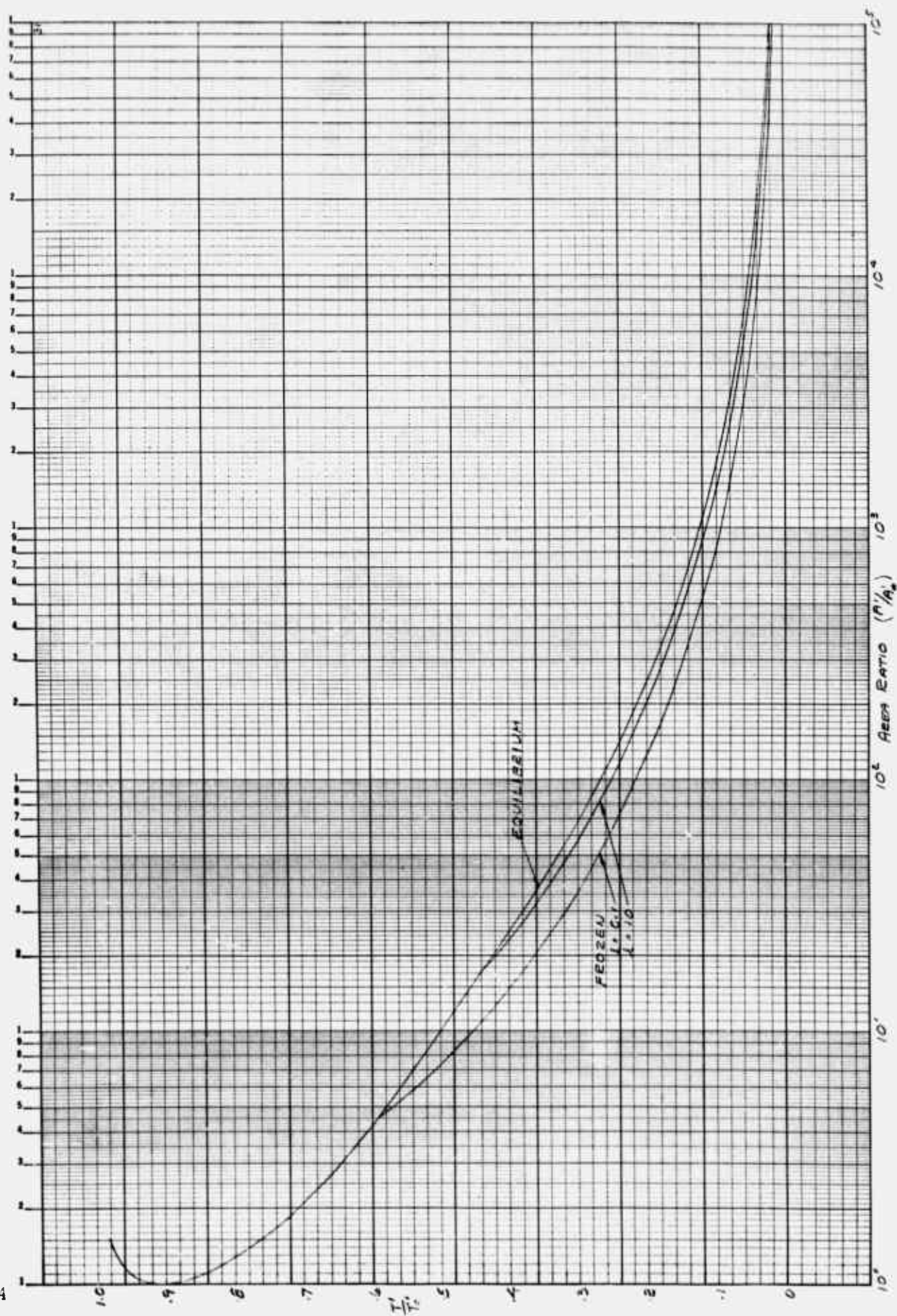


FIGURE NO. 31d HYPERBOLIC NOZZLE $T_0 = 6000^\circ K$ $P_0 = 1000$ ATM

FIGURE NO. 31c HYPERBOLIC NOZZLE $T_0 = 6000^\circ\text{K}$ $P_0 = 1000\text{ ATM}$